

Galaxy Formation

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The Hebrew University of Jerusalem
& Chaire Internationale de Recherche Blaise Pascal, Paris

**Une série de cours (en anglais) pour les étudiants
en thèse (et M2) et les chercheurs**

les mercredis, de 17h00 à 19h00 à l'amphithéâtre de l'IAP
98bis Bd Arago - Paris 14^{ème} - M° St Jacques ou Denfert-Rochereau

20 octobre

1. The standard cosmology
2. Linear growth of fluctuations by gravitational instability

A special lecture series on **Galaxy Formation**

by Avishai Dekel (Chaire Internationale Blaise Pascal)

for graduate students and researchers; IAP/OP Wednesdays 17:00-19:00

- | | |
|-------------|--|
| Octobre 20 | 1. the standard cosmology
2. linear growth of fluctuations by gravitational instability |
| Novembre 17 | 3. statistics of density fluctuations: the CDM scenario
4. nonlinear growth: spherical model, filamentary structure |
| Decembre 8 | 5. numerical simulations of structure formation
6. hierarchical clustering: Press-Schechter formalism, biasing |
| Decembre 15 | 7. dark-matter halos: density profile, cusp/core problem
8. halo substructure: dynamical friction, tidal effects, HOD |
| Janvier 5 | 9. angular momentum problem: tidal torques, disk formation
10. the origin of galaxy scaling relations and their scatter |
| Janvier 12 | 11. semi-analytic modeling: cooling, star formation, mergers
12. feedback processes: supernova, AGN and black holes |
| Fevrier 9 | 13. cold flows versus shock heating
14. origin of bi-modality in galaxies |
| Fevrier 16 | 15. dwarf galaxies and the "fundamental line"
16. dark-dark halos: effect of cosmological photoionization |

Lectures 15-16 (some 8,12)

Dwarf Galaxies & Dark-Dark Halos: Feedback Processes

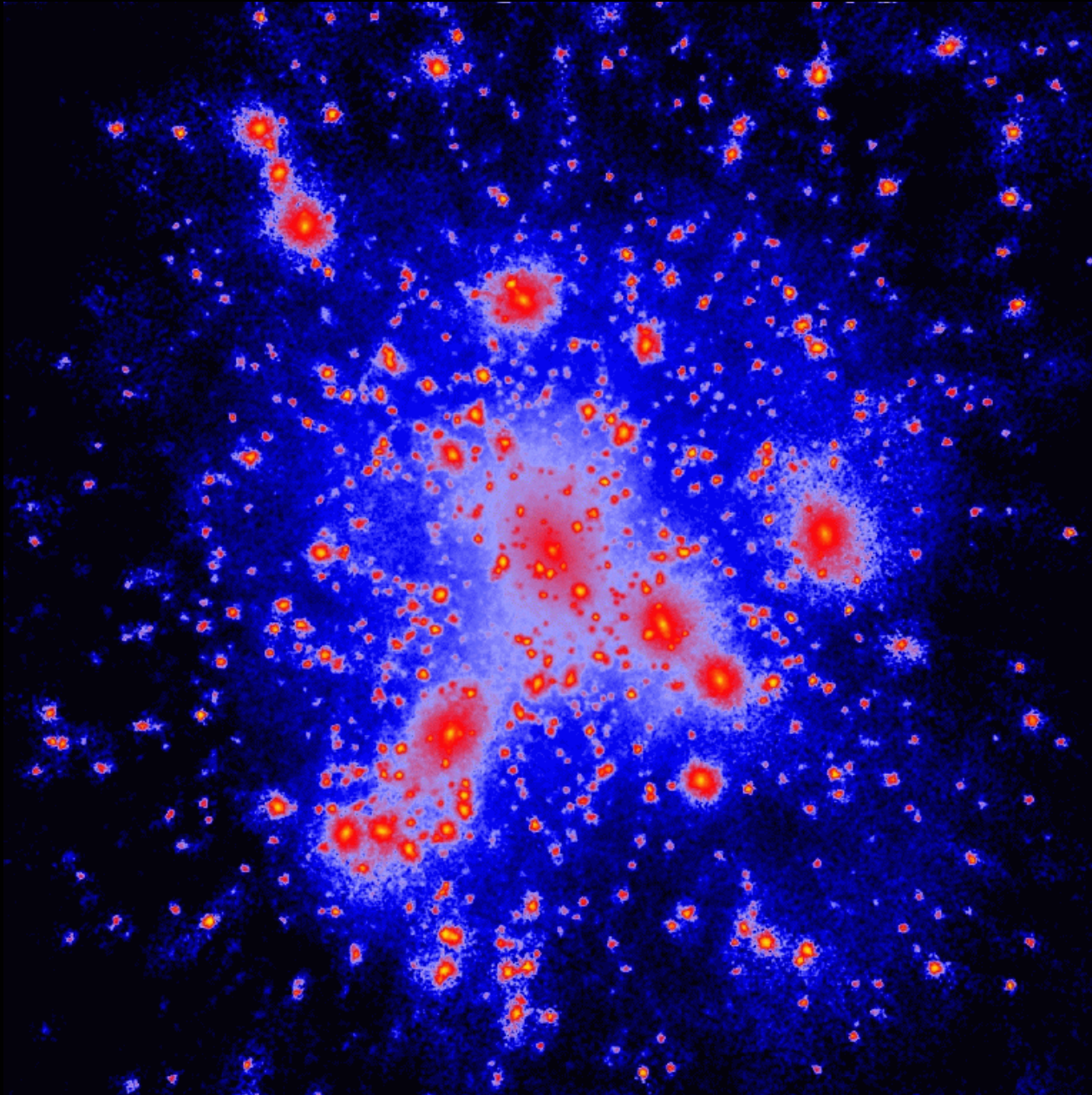
- The “fundamental line”
- Origin of scaling relations:
supernova feedback
- Dark-dark halos (DDH) must exist
- Origin of DDH by photoionization
- Halo substructure: phase-space density

1. Missing Dwarfs & the “Fundamental Line”

Dekel & Woo 2003



Λ CDM model: many dwarf satellites

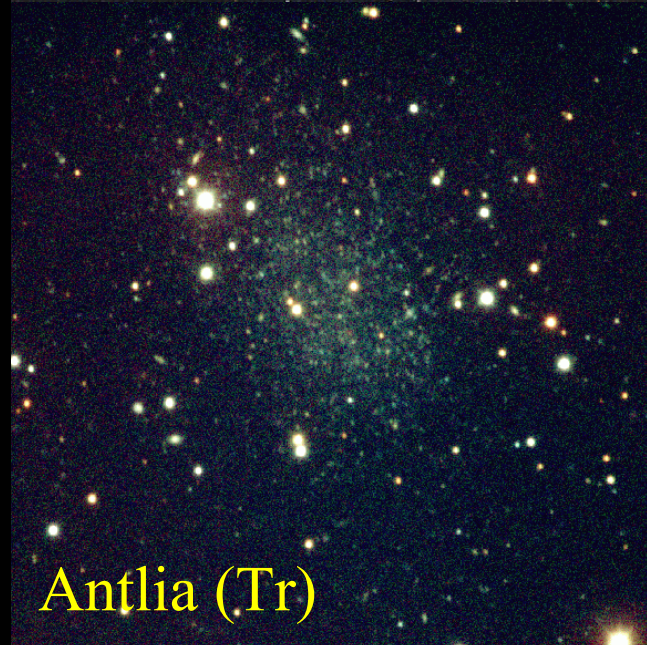
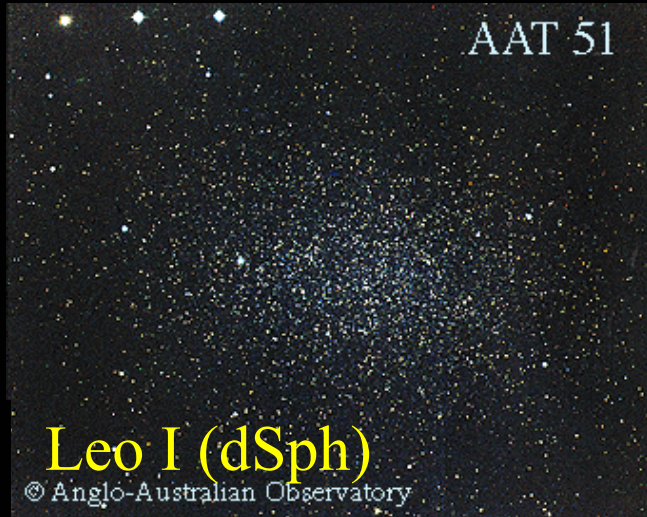


Moore et al

Low-Surface-Brightness & Dwarf Galaxies

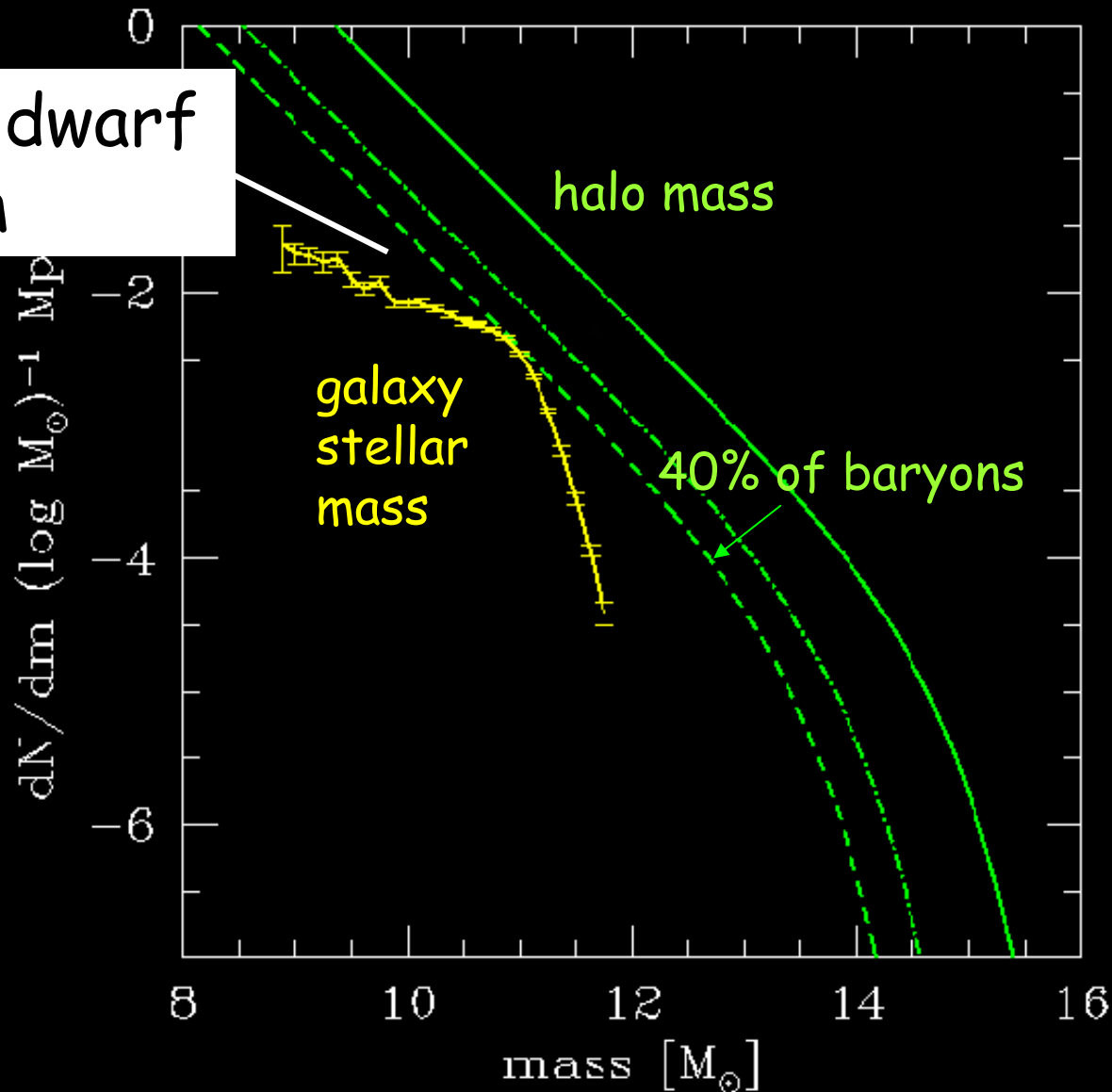


Only a few faint dwarf satellites

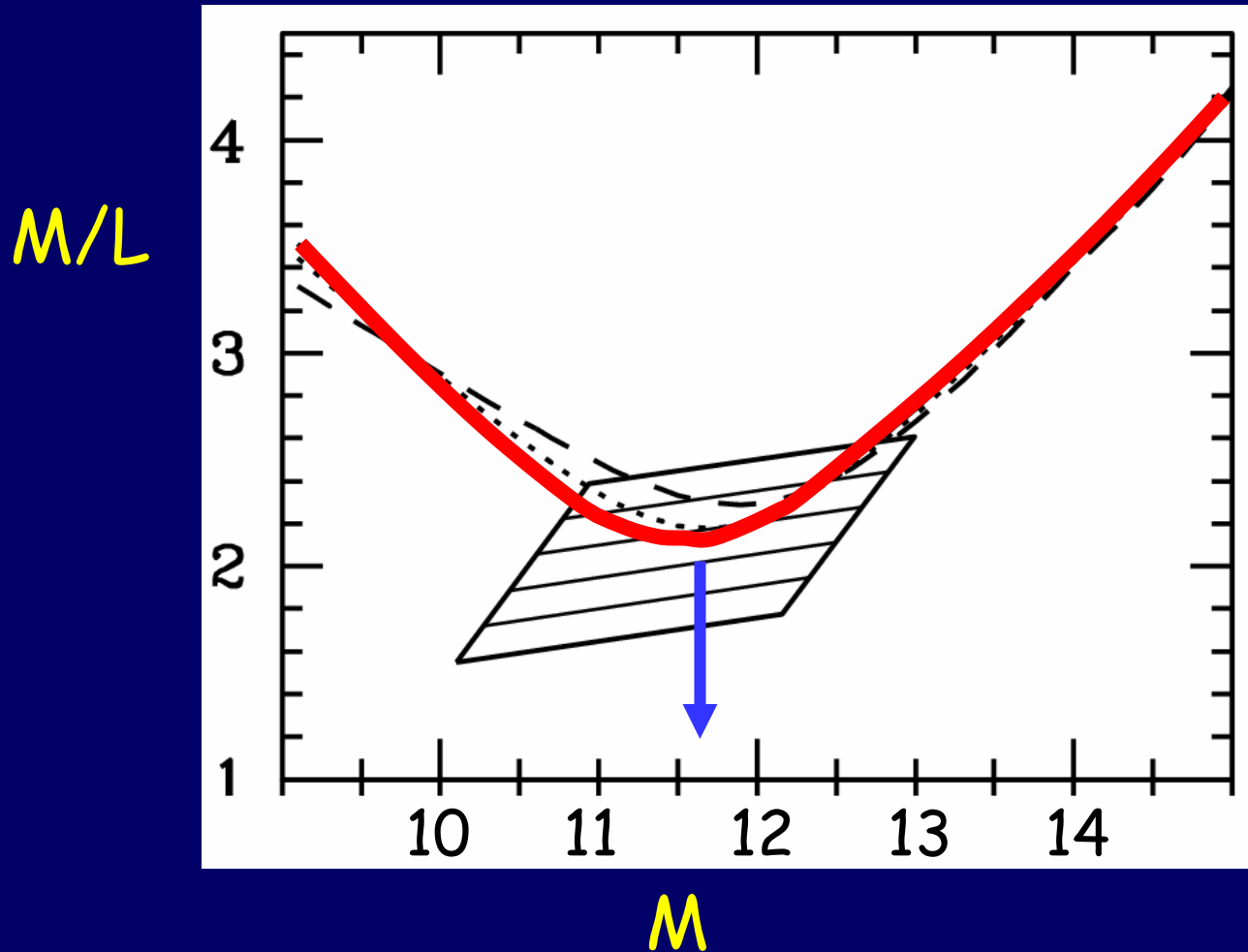


Mass versus Light Distribution

missing dwarf
problem

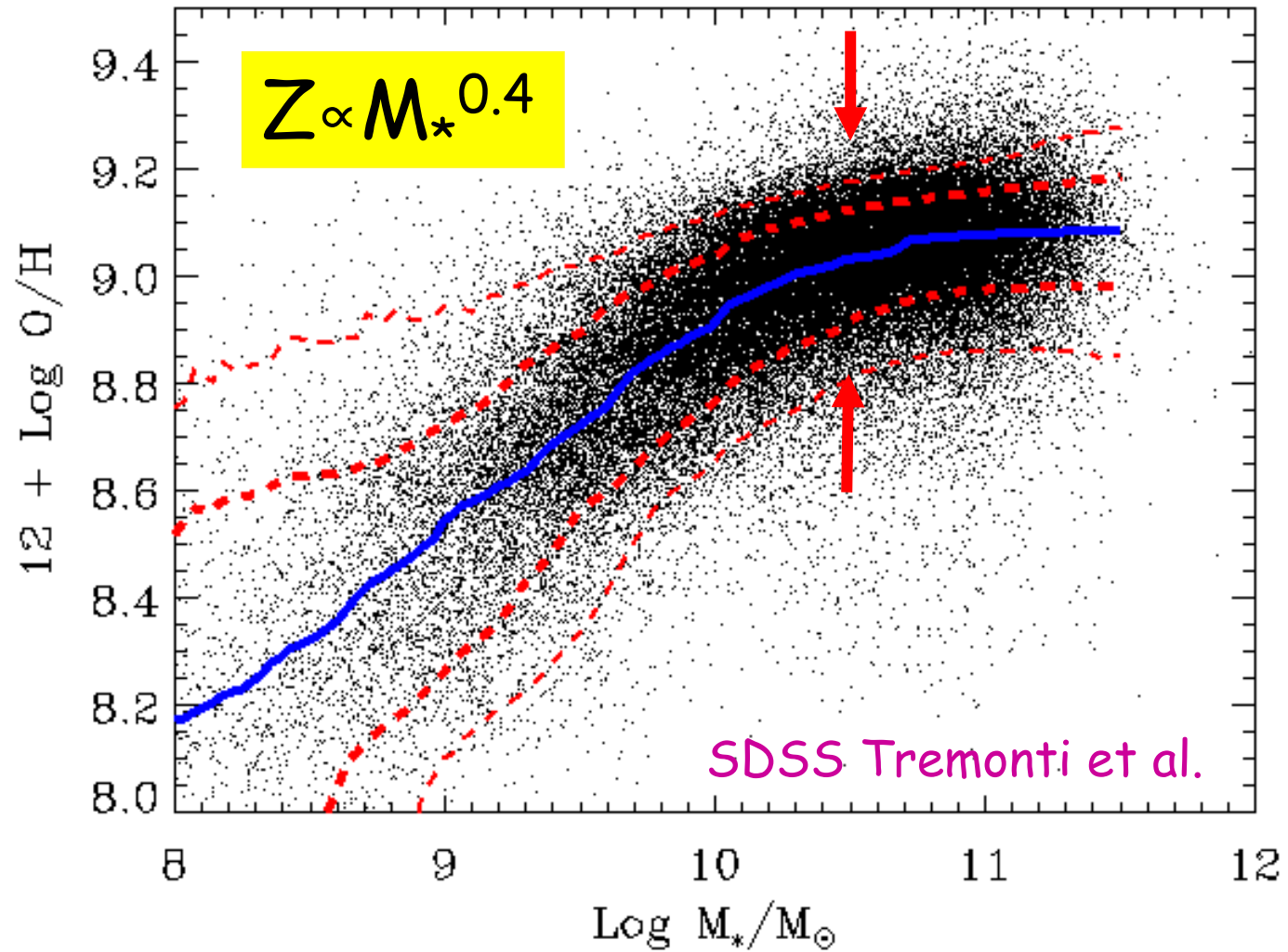


$\langle M/L \rangle$ vs M for halos in 2dF assuming Λ CDM

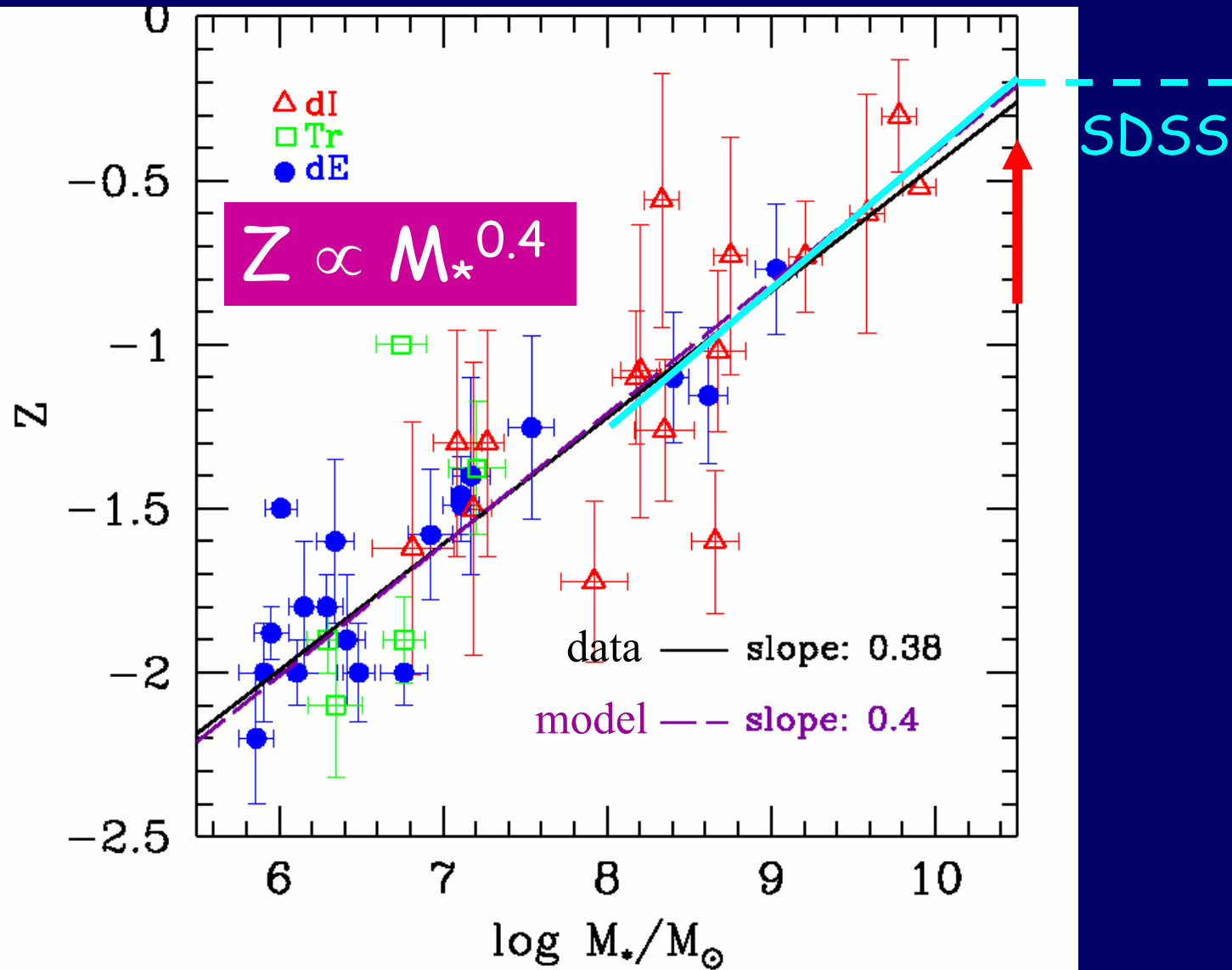


Using conditional luminosity function: Van den Bosch, Mo, Yang 03

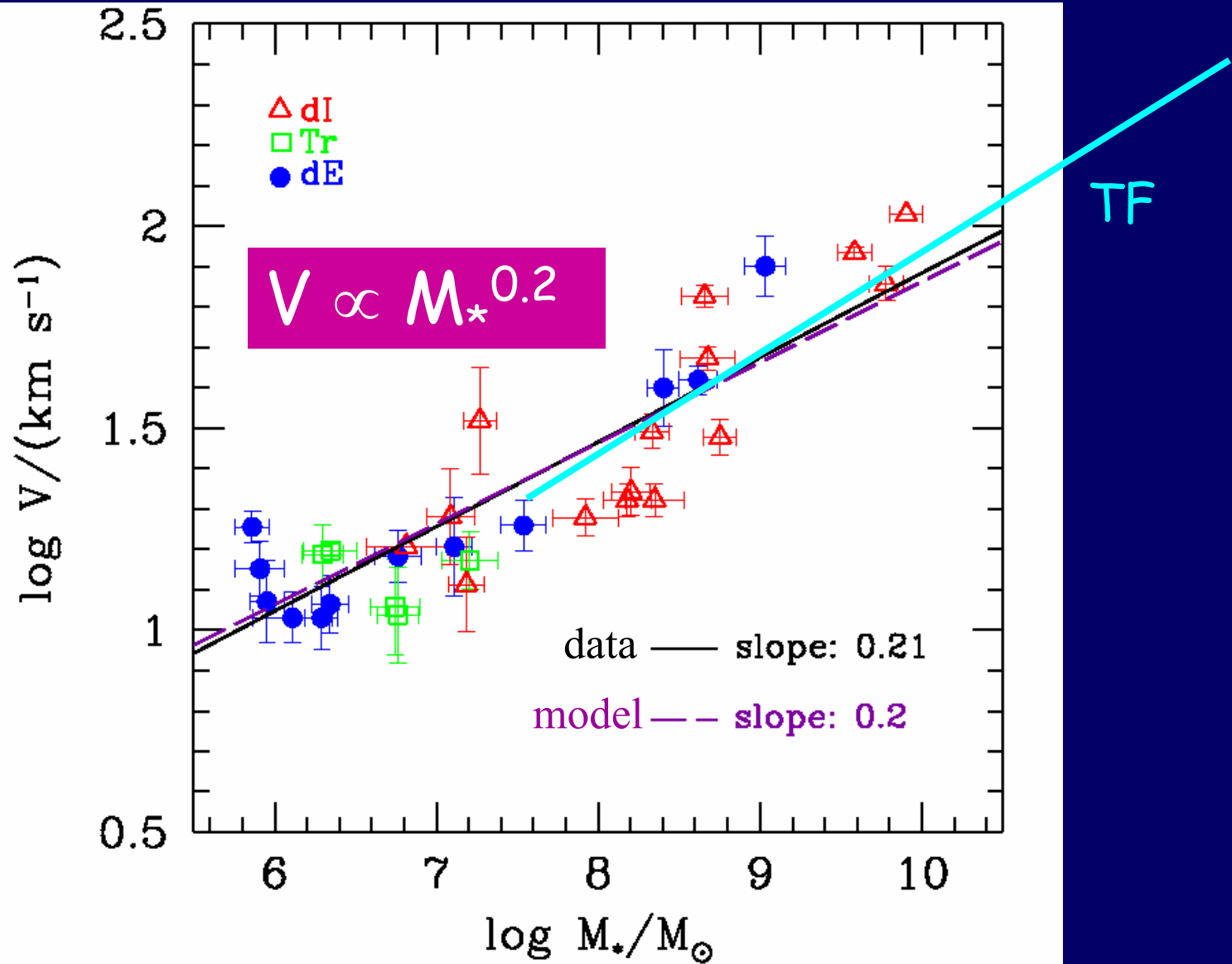
Metallicity



Local Group Dwarfs: Metallicity



LG Dwarfs: Velocity



2. Origin of Scaling Relations:

virial theorem & spherical halo collapse

angular momentum

feedback

Bright Galaxies

- virial halo

$$V^2 \propto \frac{GM}{R}$$

- top hat

$$\frac{M}{R^3} \propto 200 \rho_u$$

$$\rightarrow M \propto V^3 \propto R^3$$

- $M_* \propto M_{\text{gas}} \propto f_{\text{bar}} M$

initial

$$\rightarrow M_* \propto V^3$$

- disk size

$$R_* \approx \lambda R$$

spin

$$\lambda \approx \text{const.}$$

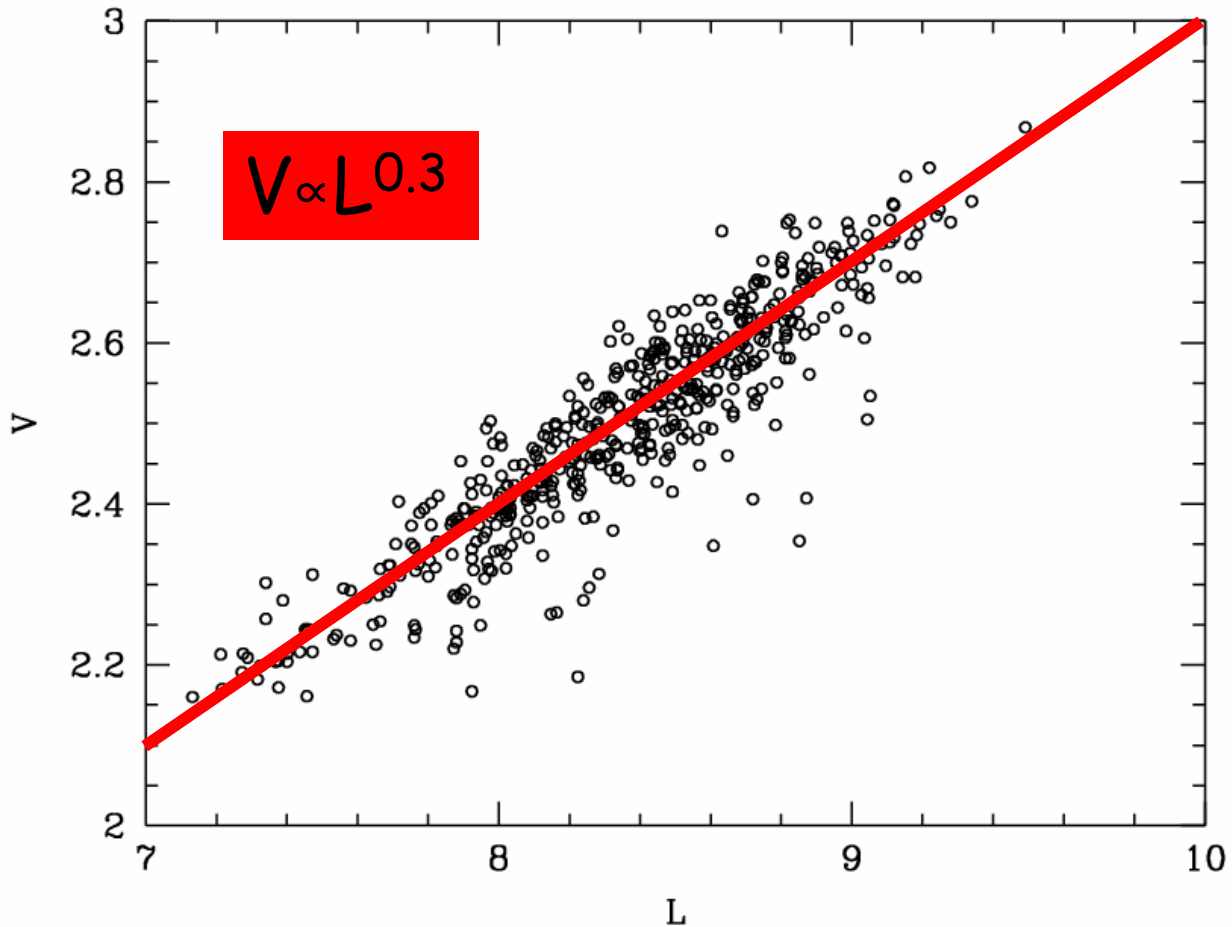
$$\mu_* \propto M_* / R_*^2 \propto \lambda^{-2} M_*^{1/3}$$

$$\rightarrow \mu_* \propto M_*^{1/3}$$

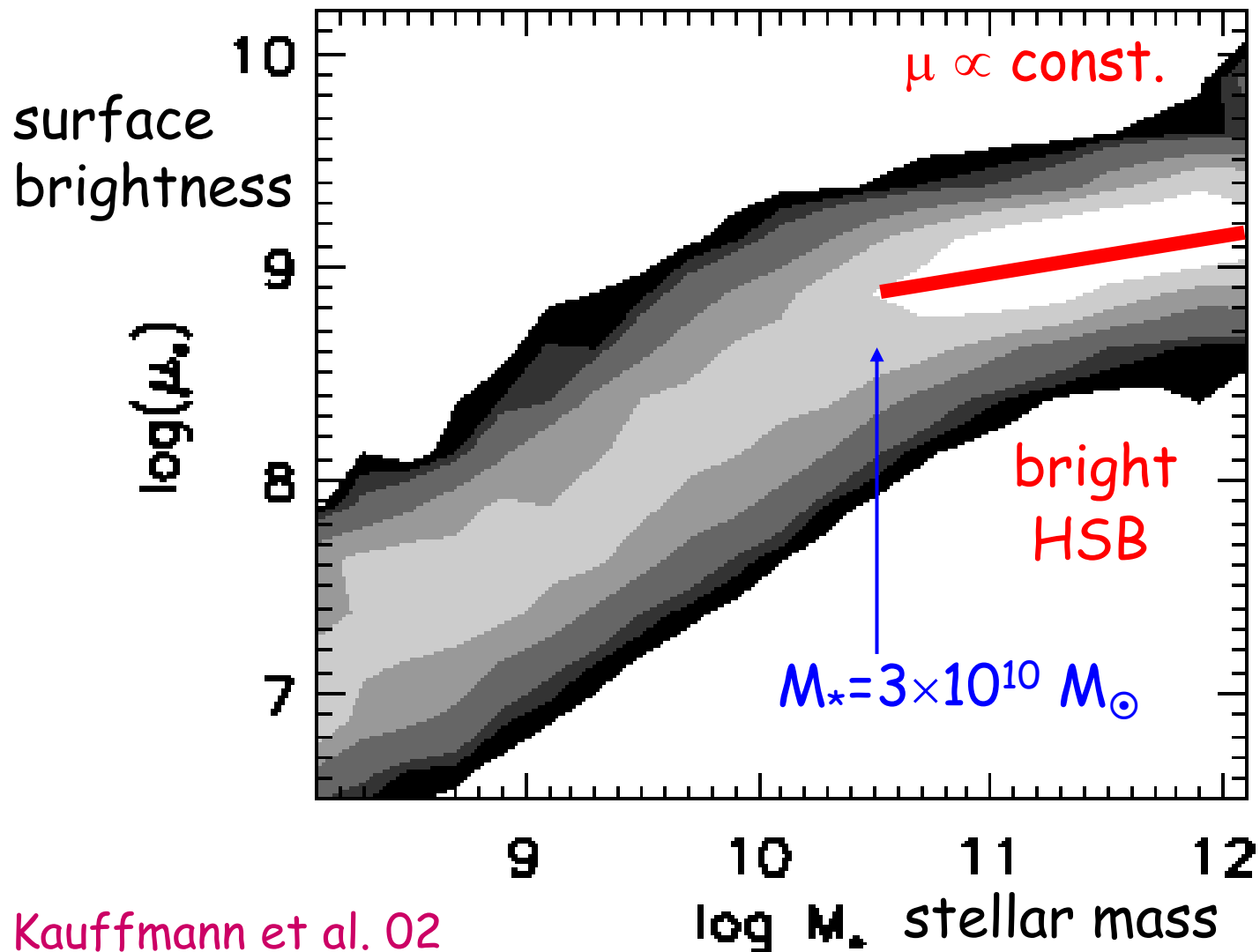
- $Z \propto M_* / M_{\text{gas}}$

$$\rightarrow Z \propto \text{const.}$$

Bright Galaxies: Tully Fisher Relation

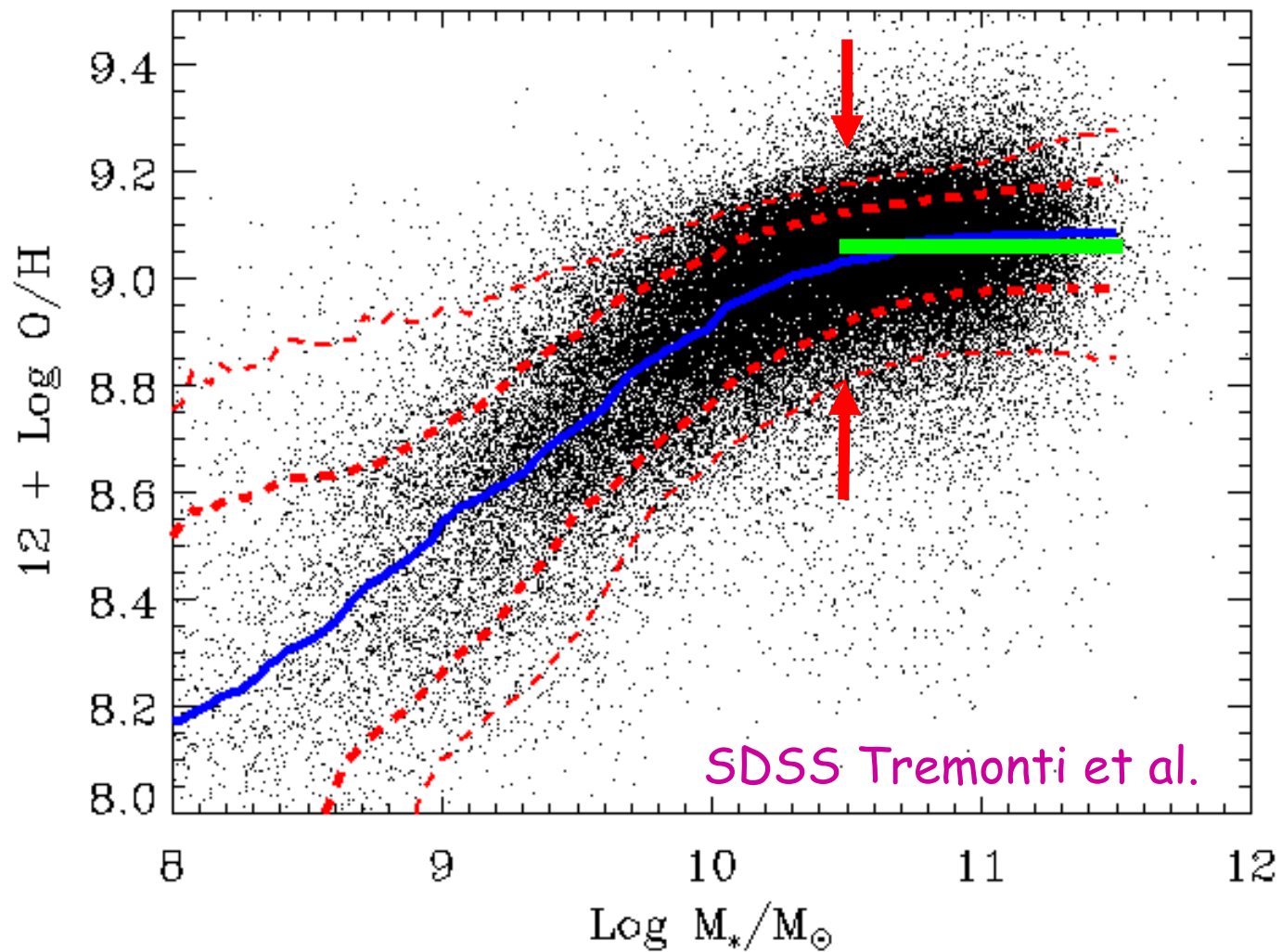


Surface Brightness: SDSS



Kauffmann et al. 02

Metallicity SDSS

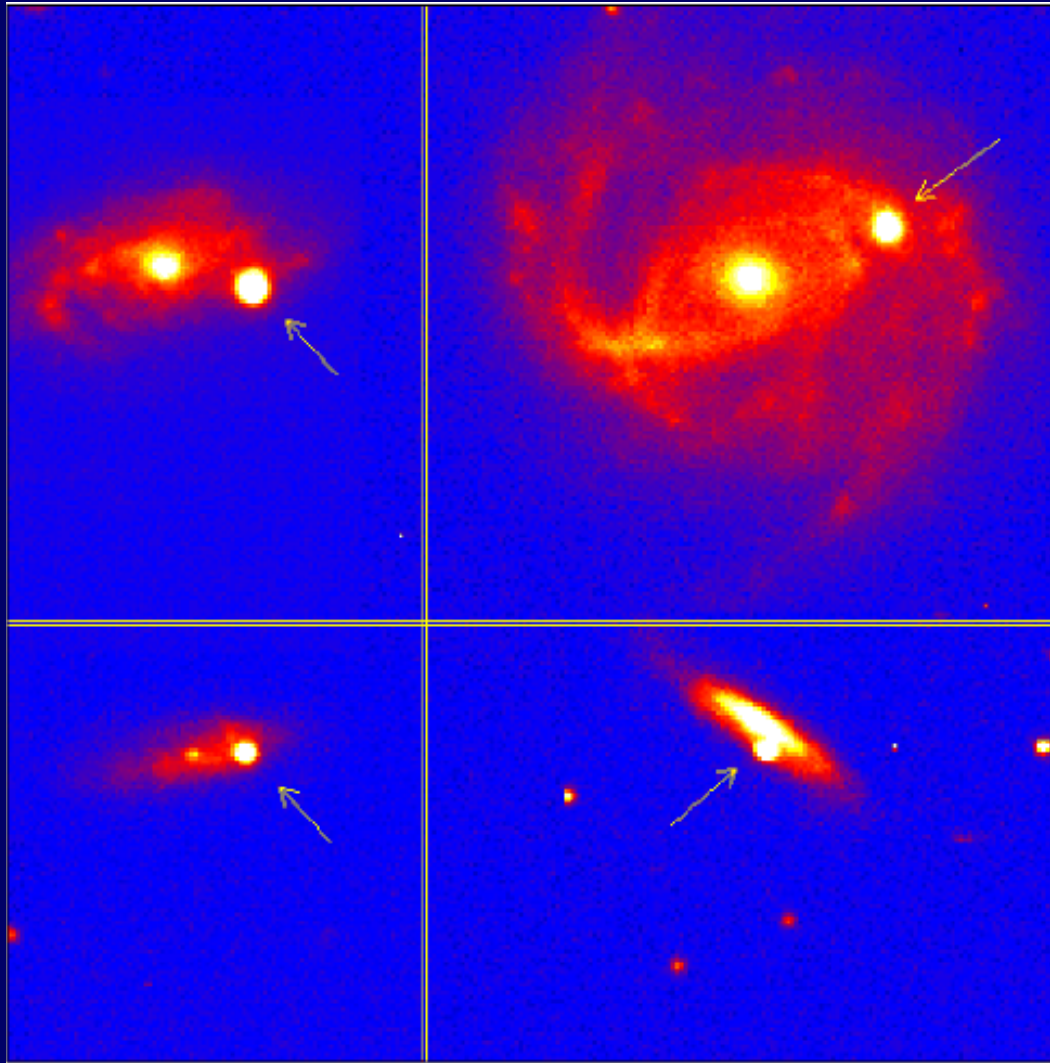


Supernova Feedback

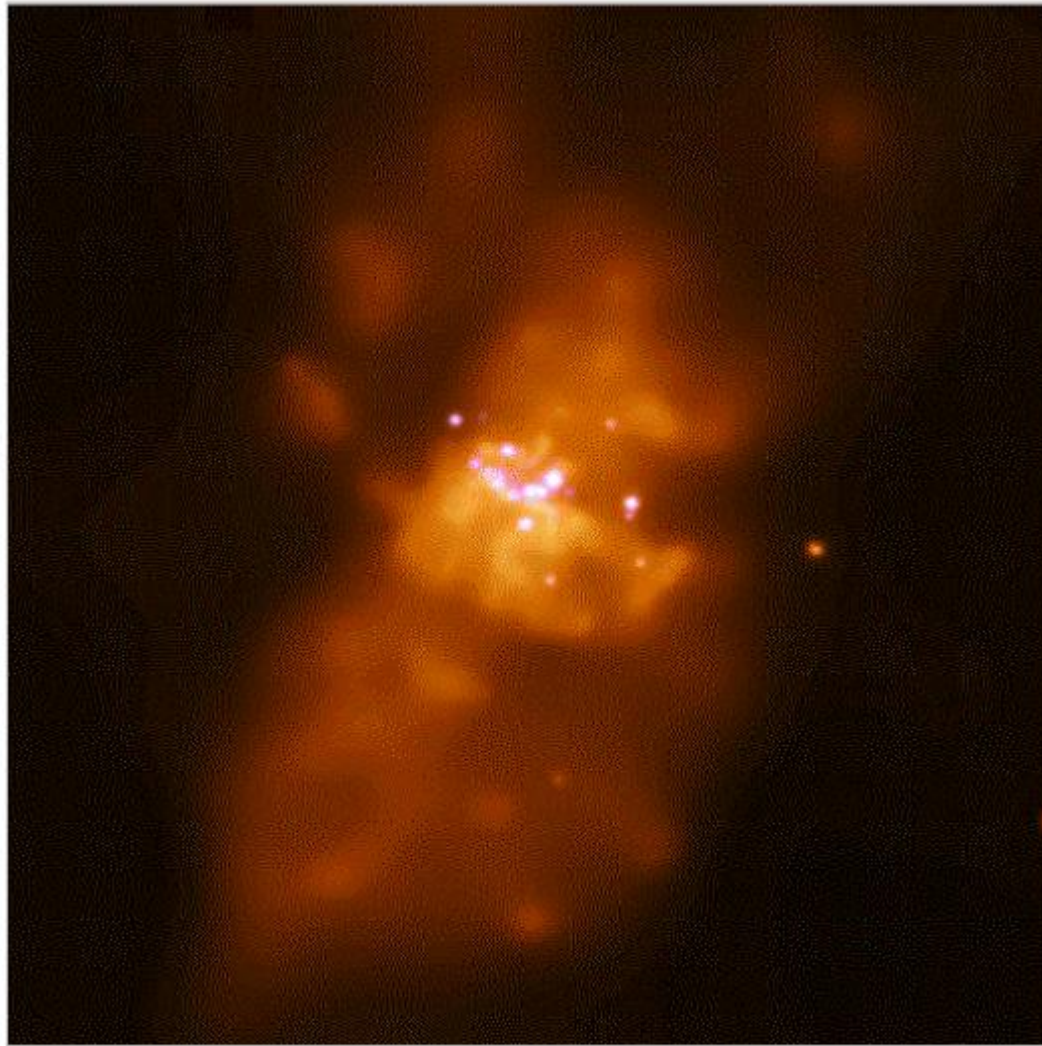
Dekel & Silk 86
Dekel & Woo 03



Much energy in SNe



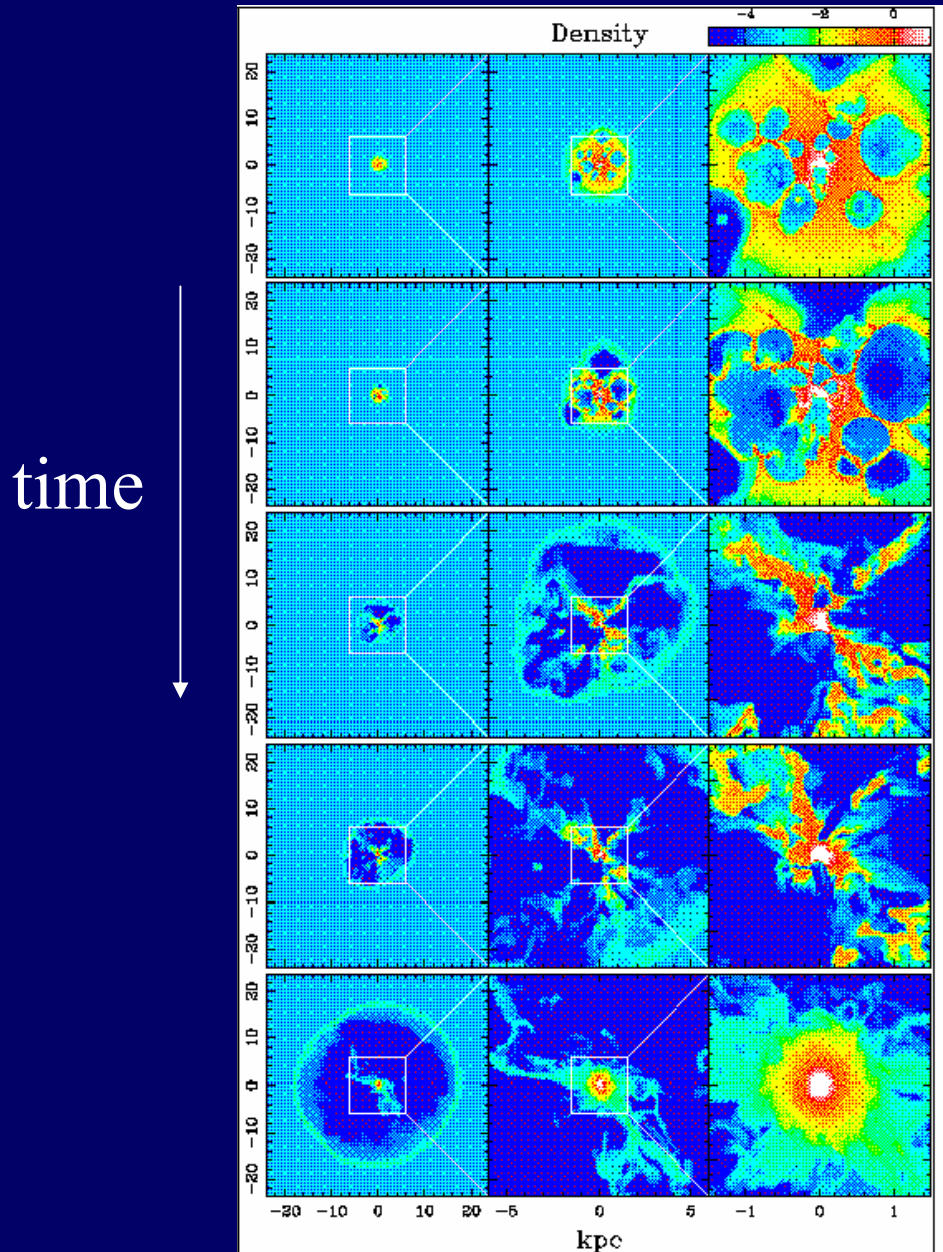
Galactic wind M82



Chandra X-Ray Observatory image of M82

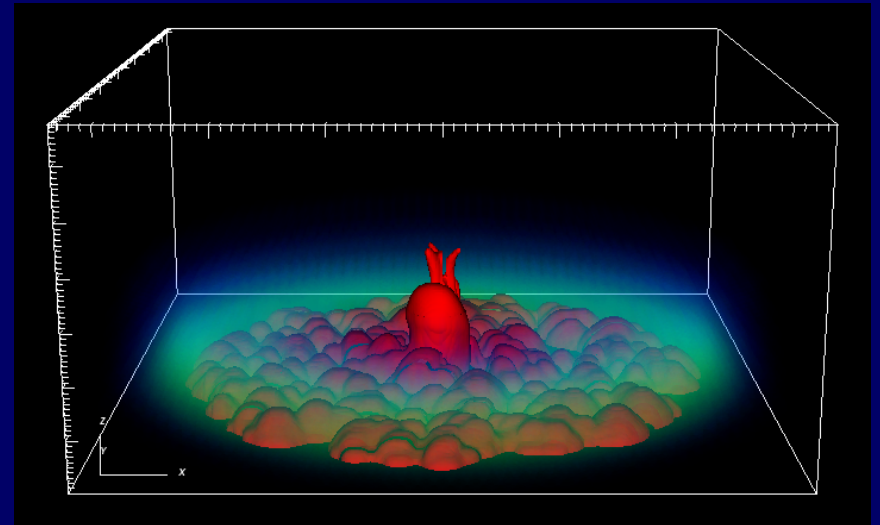
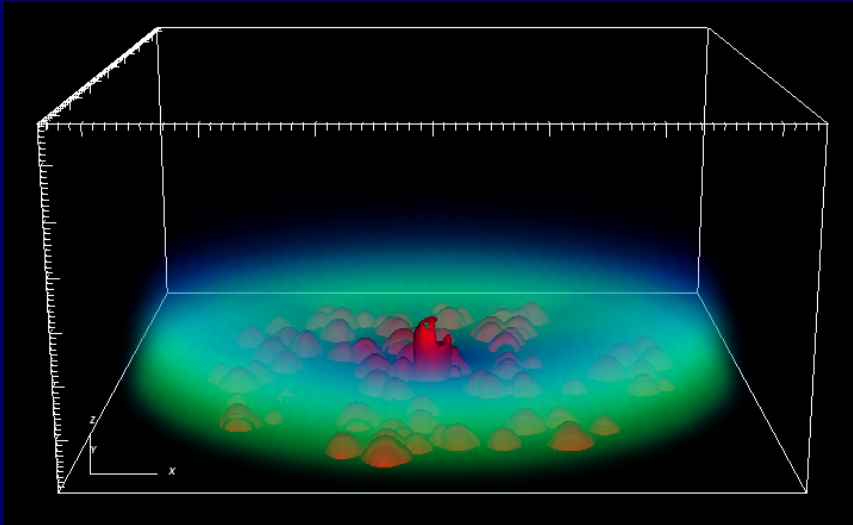
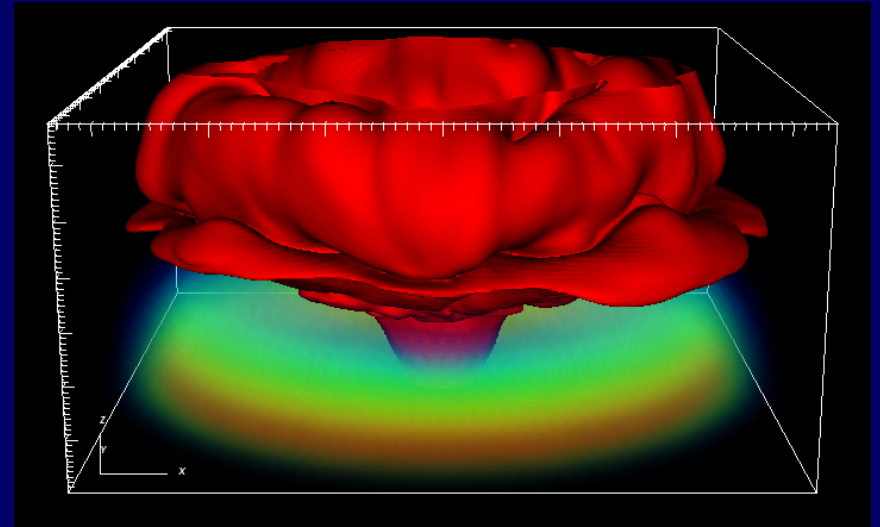
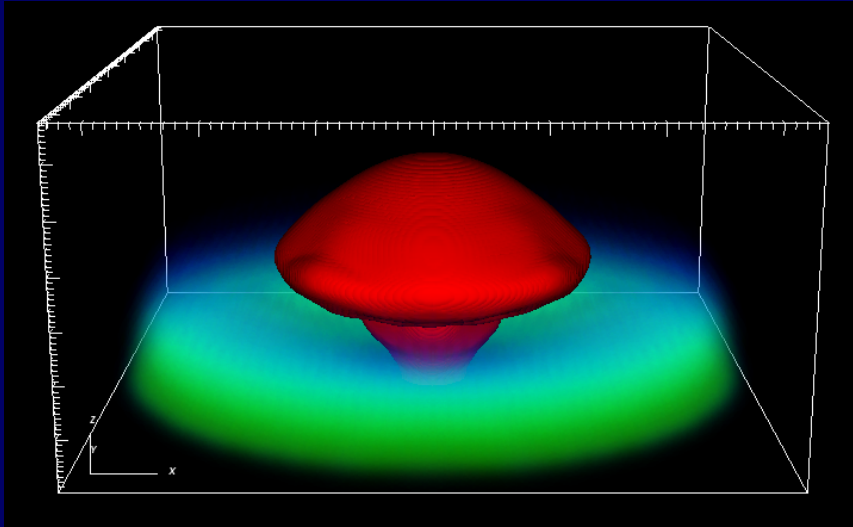


Simulation of supernova blowout



Mori et al.

Supernova Feedback



Fragile, Murray, Lin 04

Supernova Feedback Scale

(Dekel & Silk 86)

Energy fed to the ISM during the “adiabatic” phase:

$$E_{\text{SN}} \approx \nu \varepsilon \dot{M}_* t_{\text{rad}} \propto M_* (t_{\text{rad}} / t_{\text{ff}})$$

$$\dot{M}_* \approx M_* / t_{\text{ff}}$$

$$\approx 0.01$$

for $\Lambda \propto T^{-1}$ at $T \sim 10^5 \text{ K}$

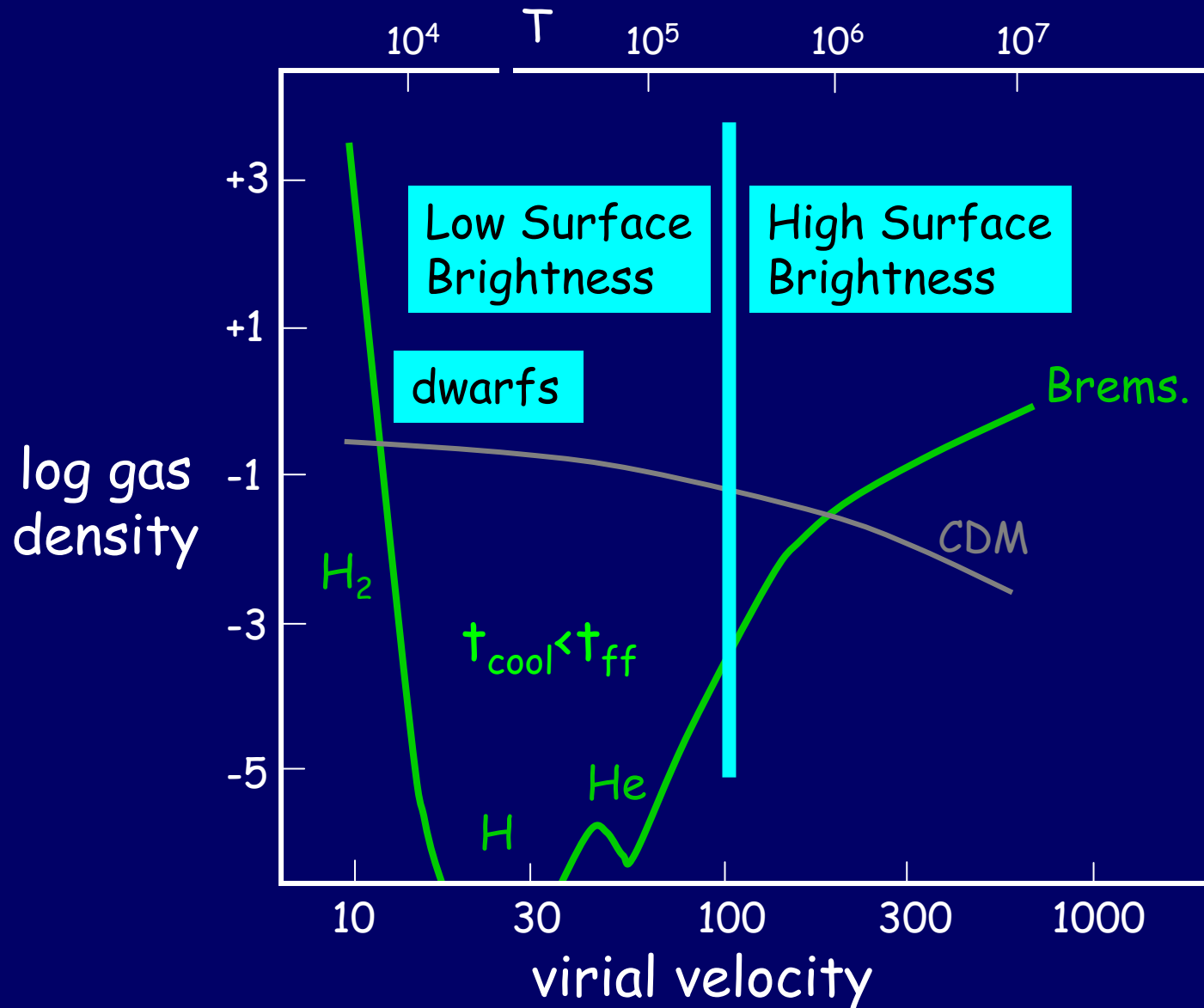
Energy required for blowout:

$$E_{\text{SN}} \approx M_{\text{gas}} V^2$$

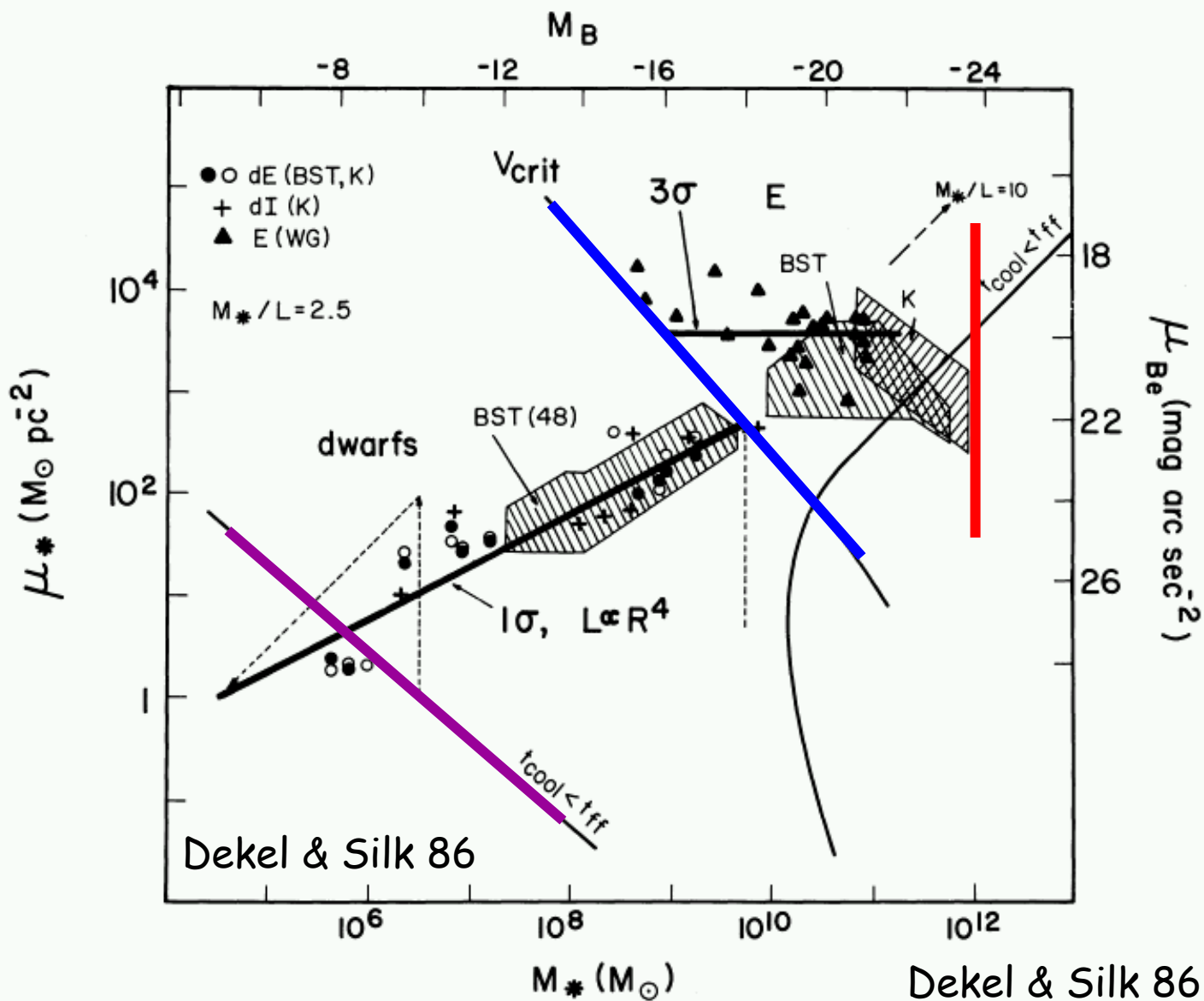
$$\rightarrow V_{\text{crit}} \approx 100 \text{ km/s} \rightarrow M_{*\text{crit}} \approx 3 \times 10^{10} M_{\odot}$$

Supernova Feedback Scale

Dekel & Silk 86



LSB vs HSB



Model: fundamental line of LSB/Dwarfs

(Dekel & Woo 03)

• Energy: $E_{\text{SN}} \propto M_* \propto M_{\text{gas}} V^2$

$$M_* / M_{\text{gas}} \propto V^2$$

• Virial halo: $V^3 \propto M \propto R^3$

$\ll 1$

$$V \propto M_*^{1/5}$$

“Tully Fisher”

$$Z \propto M_* / M_{\text{gas}}$$

$$Z \propto M_*^{2/5}$$

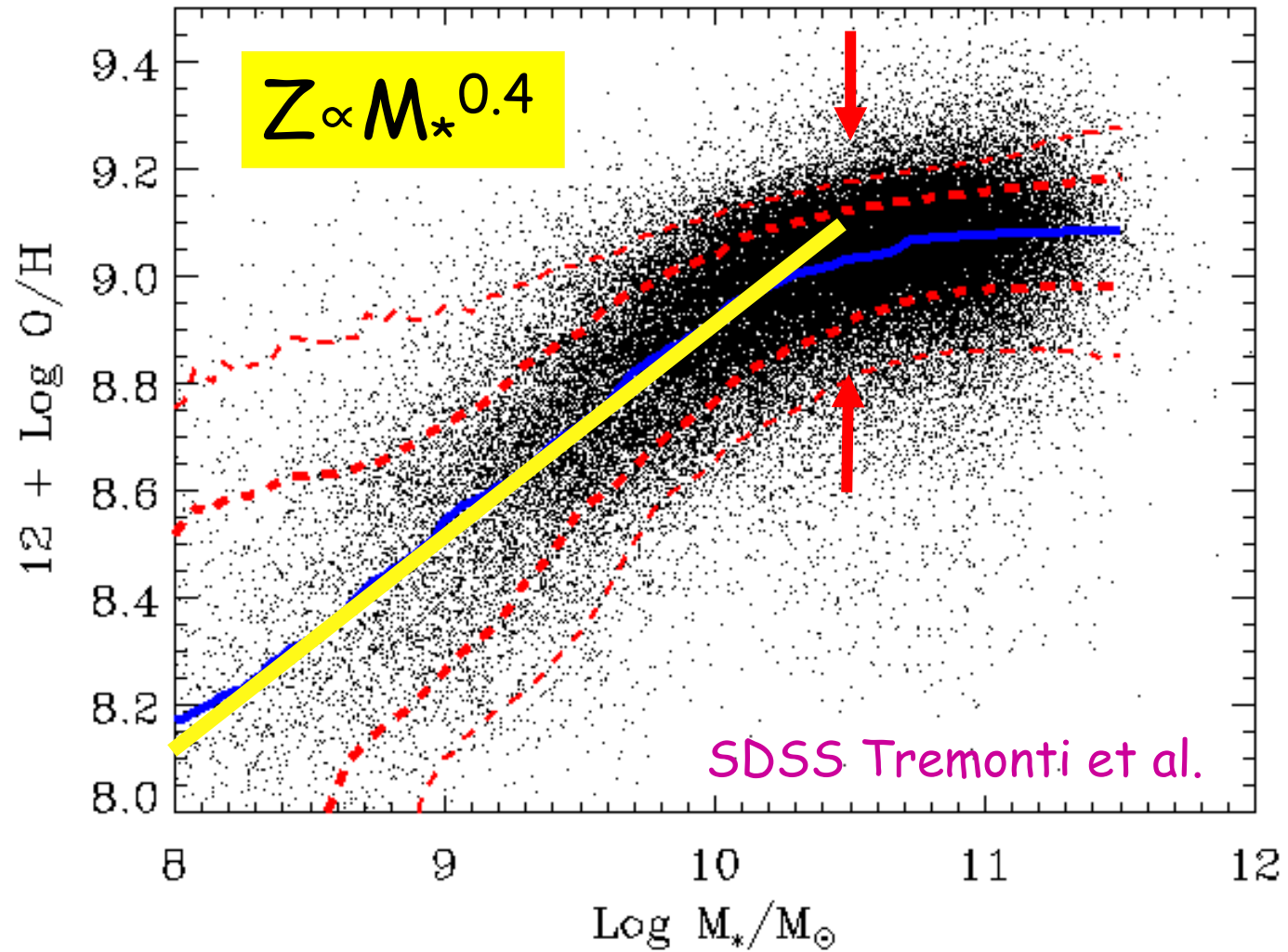
metallicity

$$R_* \approx \lambda R$$

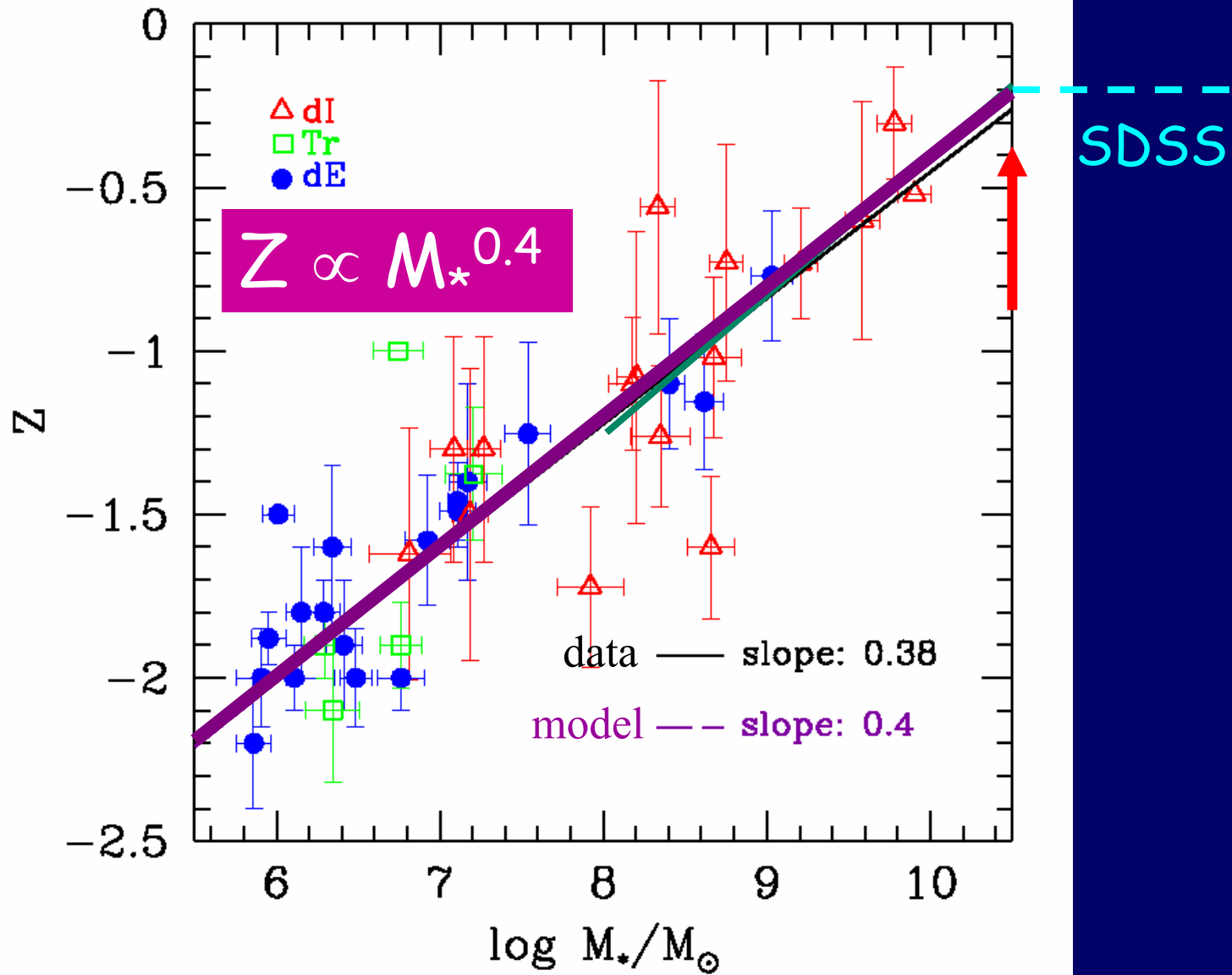
$$\mu \propto \lambda^{-2} M_*^{3/5}$$

surface brightness

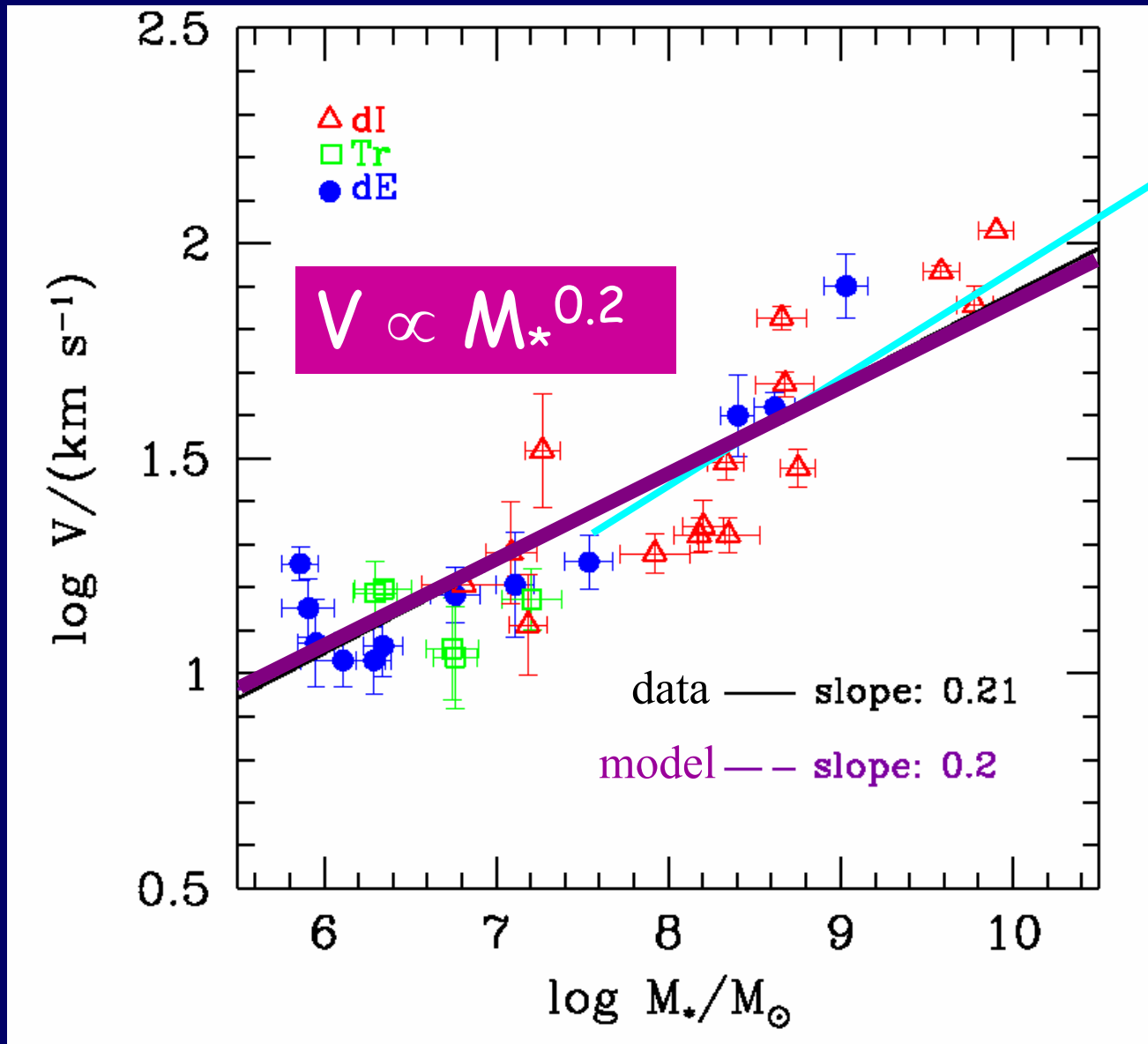
Metallicity



Local Group Dwarfs: Metallicity



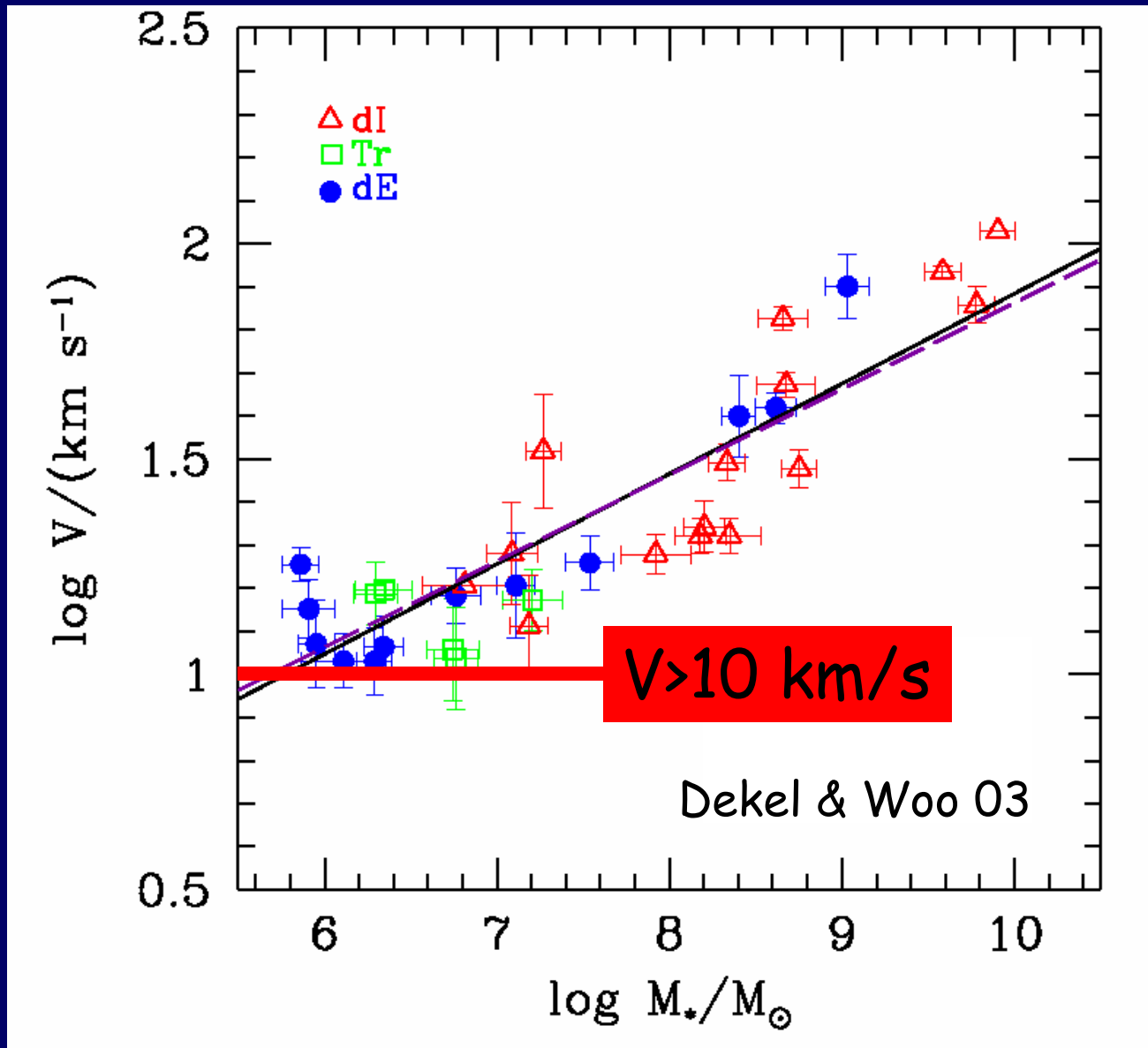
LG Dwarfs: Velocity



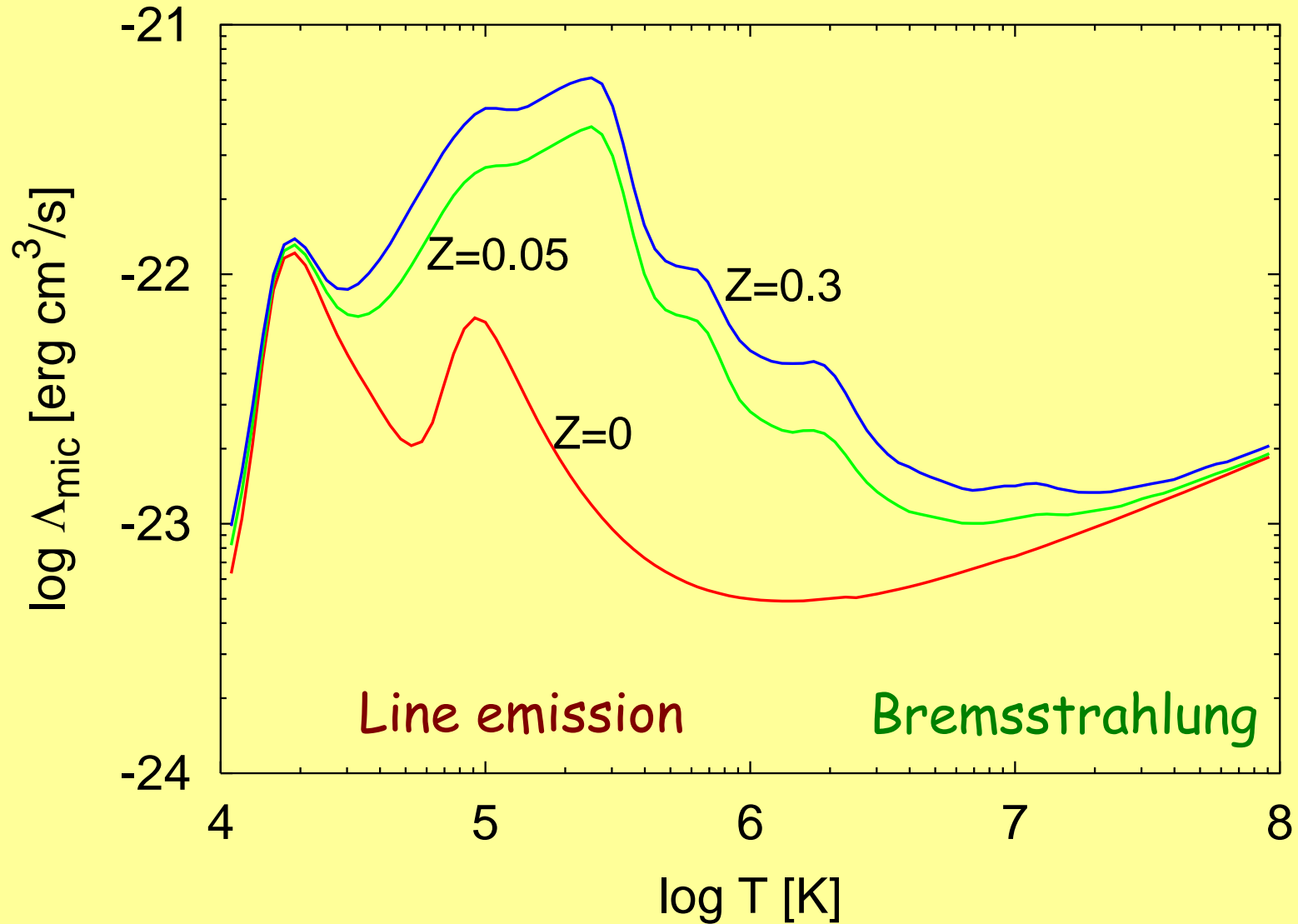
Summary: SN feedback

Could be responsible for the transition scale at $M_* = 3 \times 10^{10}$, and the “fundamental line” of LSB/dwarf galaxies, $M^*/M \propto V^2$.

A lower bound for galaxies



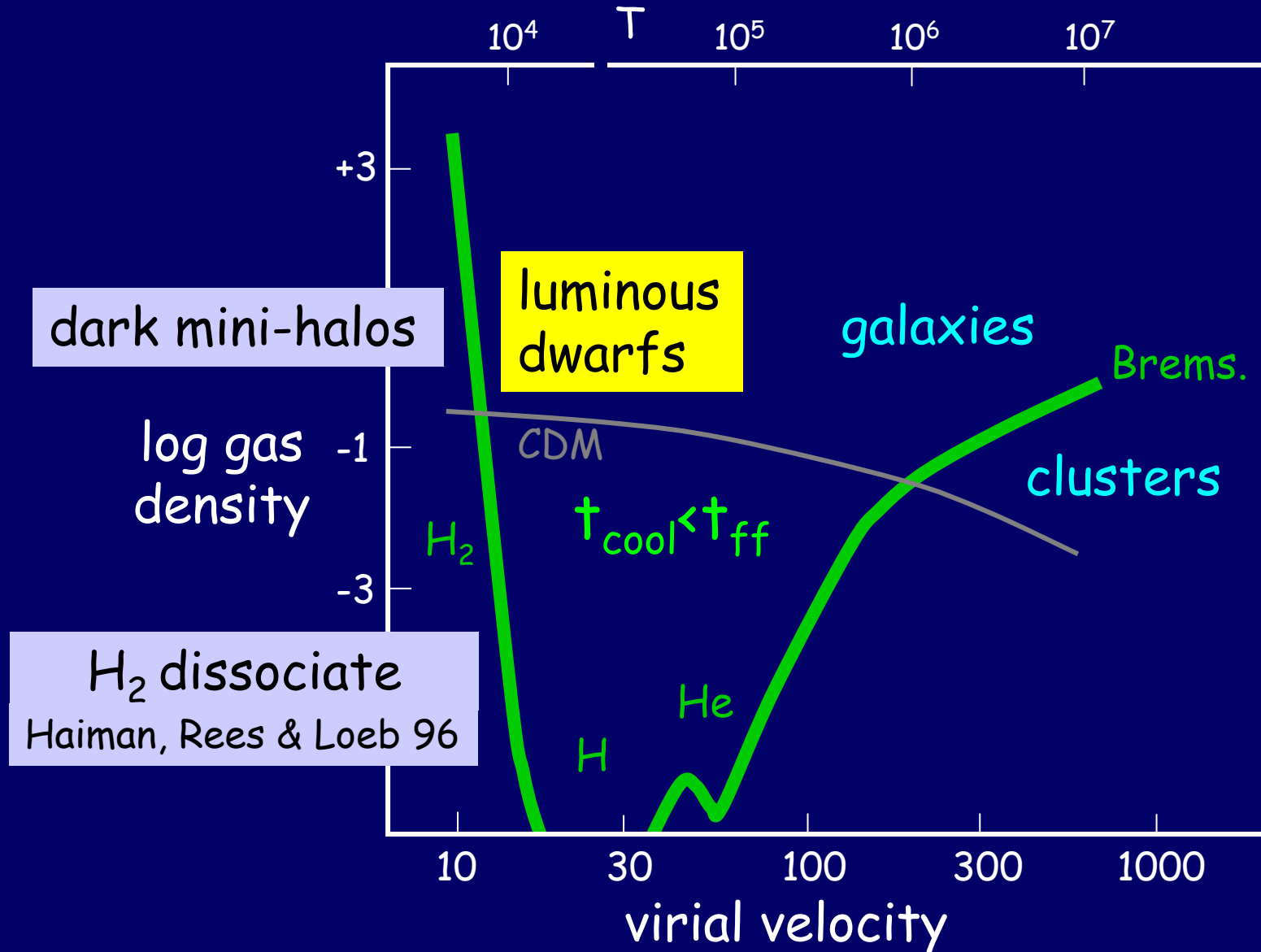
Cooling rate



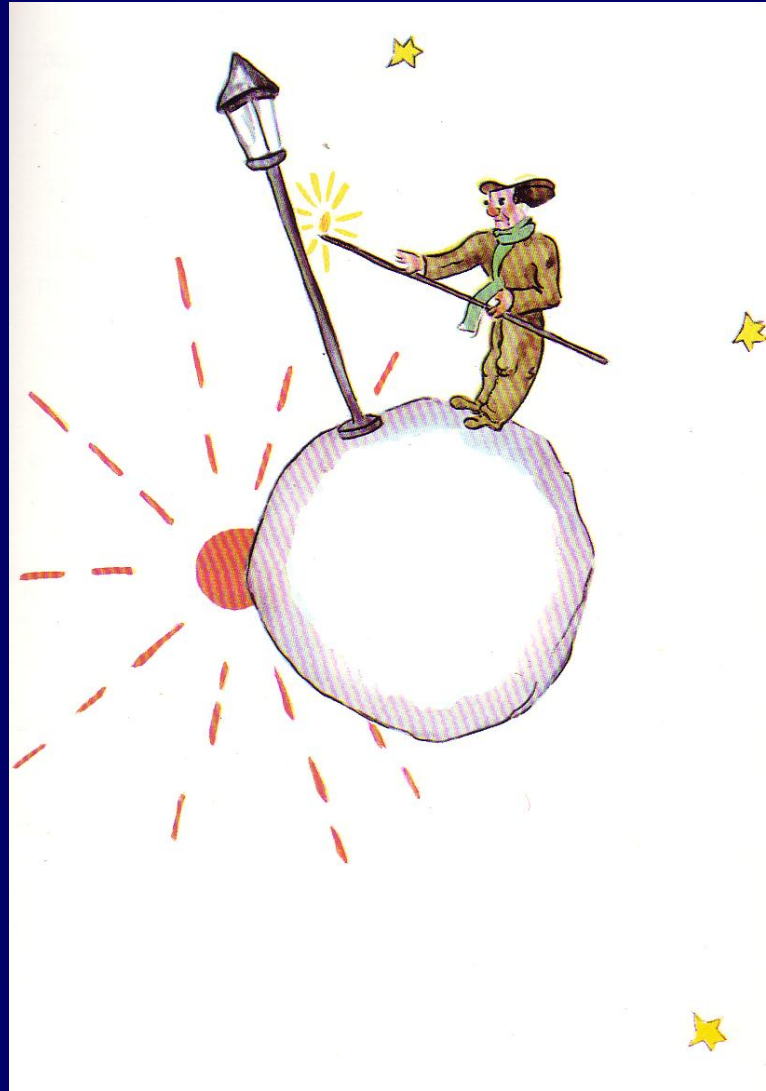
$$q = \frac{N_A^2 \chi^2}{\mu^2} \Lambda(T) \rho \quad [\text{erg g}^{-1} \text{ s}^{-1}] \quad N_A / \mu \text{ molecules per g} \quad \chi e^- \text{ per particle}$$

The Cooling Barrier

Rees & Ostriker 77, Silk 77, White & Rees 78



It isn't that simple to turn on the light



3. Dark-Dark Halos Must Exist



Dark-Dark Halos at $V < 30$ km/s

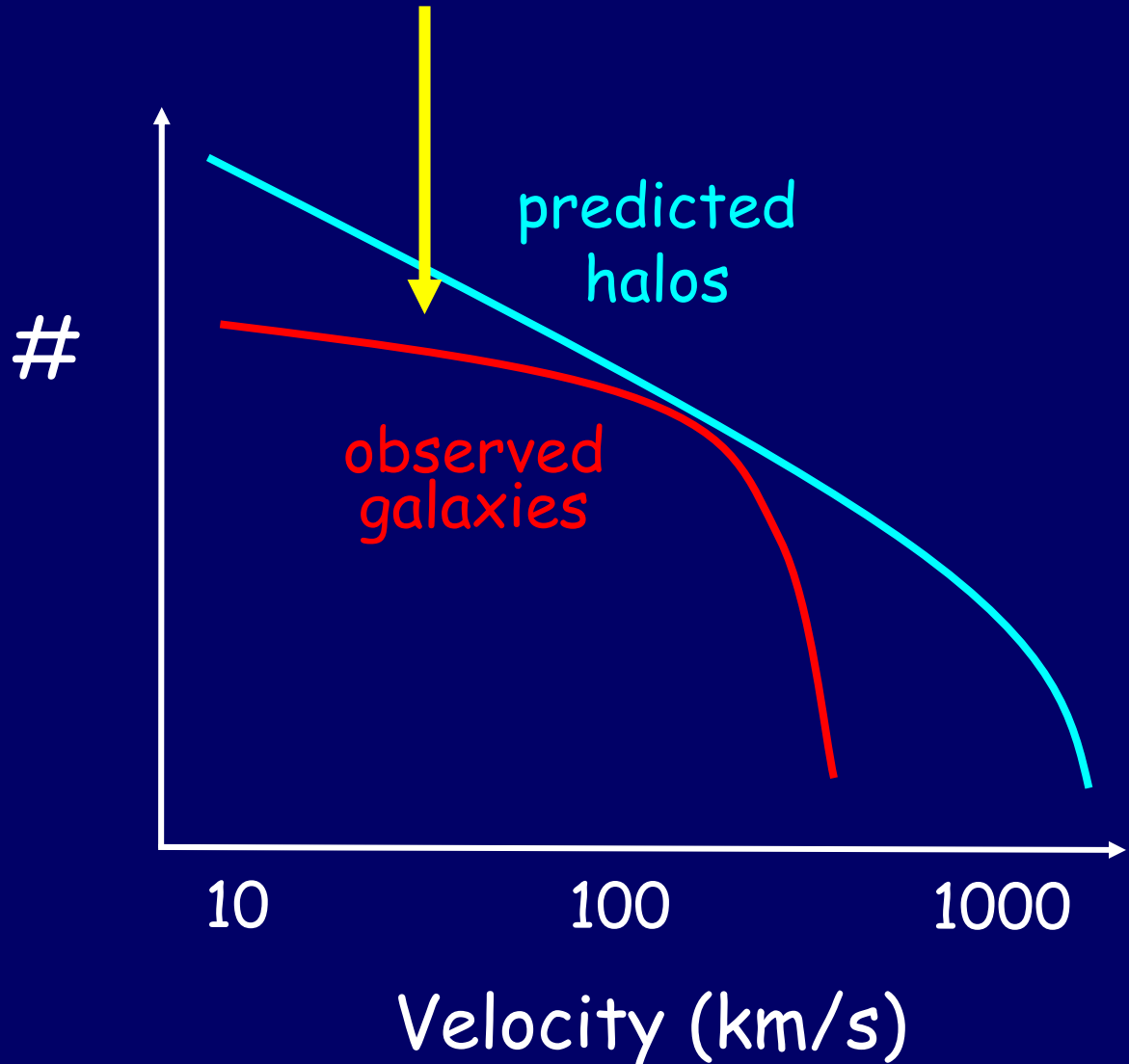
$$\text{TF: } L \sim V^4$$

$$\text{Virial: } M \sim V^3$$

$$\psi(M) \sim M^{-2}$$

$$\phi(L) \sim L^{-1}$$

Cannot be
reconciled!



Dark Dark Halos must exist !

virial, top-hat:

$$M \propto V^3$$

Tully Fisher:

$$L \propto M^x \propto V^{3x} \quad 3x \approx 5 \quad \rightarrow dL/dM \propto M^{x-1}$$

luminosity function:

$$\varphi(L) \propto L^{-\alpha} \quad \alpha \approx 1.2$$

mass function:

$$\psi(M) \propto M^{-\beta} \quad \beta \approx 1.8$$

$$f_L(M) \psi(M) dM = \varphi(L) dL \rightarrow dL/dM \propto M^{\alpha x - \beta + \gamma}$$

$$\Rightarrow (\alpha - 1)x = \beta - 1 - \gamma$$

$$\rightarrow \gamma \approx 0.5 - 0.8$$

Cannot reconcile TF with luminosity and mass functions !

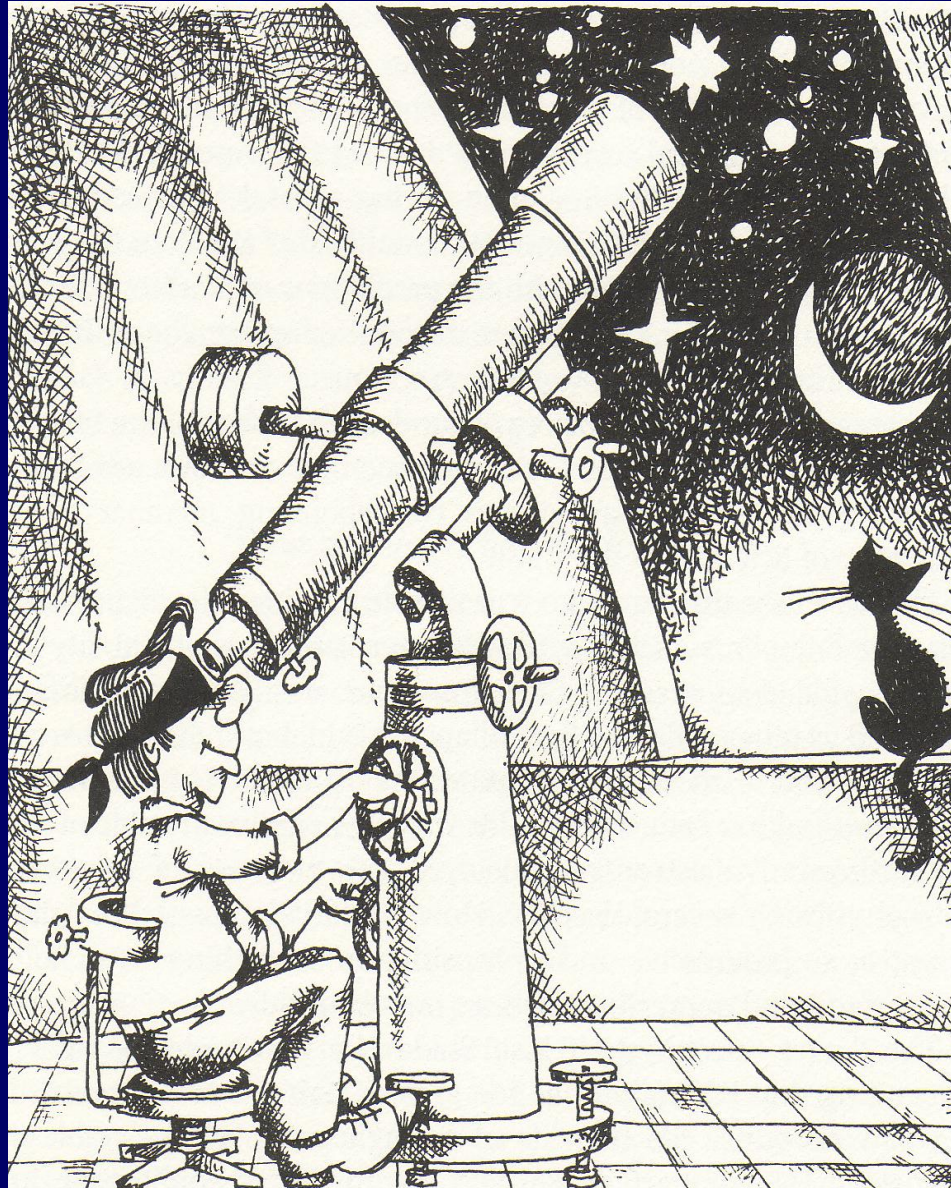
fraction of halos with luminous component:

$$f_L \propto M^\gamma$$

$$\rightarrow f_L \propto M^{0.5-0.8} \propto V^2 \quad L/M \propto V^2$$

completely dark halos SN feedback

Search for DDH



Complete removal of gas from proto-halos?

By **SN** outflow? unlikely



By **ram pressure** due to outflow from a nearby galaxy (Scannapieco, Ferrara & Broadhurst 00) ?

By **radiative** feedback?

4. Evaporation by Thermal Winds

Shaviv & Dekel 2003



Radiative Feedback

Reionization of H by UV flux from stars and AGN
by $z_{\text{ion}} \sim 10 \rightarrow$ heating gas to $T \approx (1-2) \times 10^4 \text{K}$

Jeans scale – **no** infall into halos of $V < 30 \text{ km/s}$

Efstathiou 92; Thoul & Weinberg 96; Gnedin & Ostriker 97; Gnedin 00

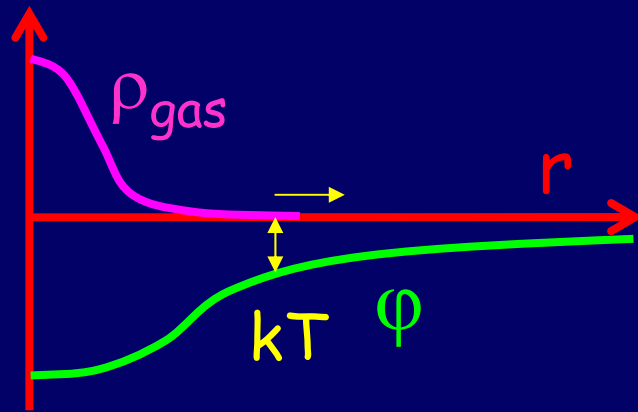
But complete gas removal?

Evaporation from halos of ~~$V < 10 \text{ km/s}$~~ Barkana & Loeb 99

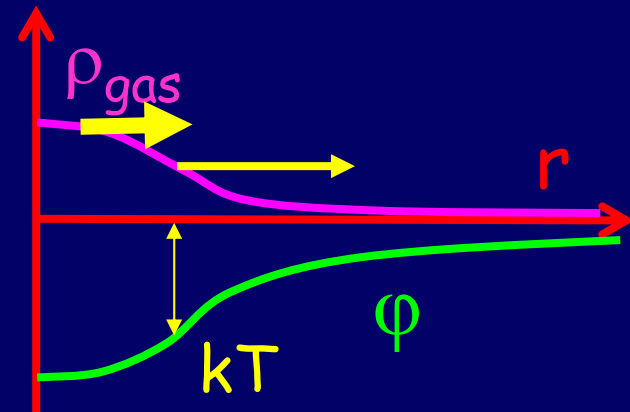
$V < 30 \text{ km/s}$ Shaviv & Dekel 03

May eliminate luminous dwarfs in small halos, $10 < V < 30$

Evaporation of hot gas



cold gas



hot gas

Mass loss from top of potential well $t_{\text{evap}} \approx t_{\text{dyn}} e^{\phi/kT}$

It is continuously replenished and lost

Continuous energy input by the ionizing flux
→ steady wind

Steady Thermal Wind

In stars: Parker 1960. In galaxies: extended potential well

Hydrodynamics:

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \mathbf{v})$$

$$\rho \frac{D\mathbf{v}}{Dt} = -\nabla P + \mathbf{f}_{\text{grav}}$$

$$P = c_s^2 \rho$$

Assume:

spherical, $c_s = \text{const.}$, steady state $\dot{M}(r) = \text{const.} \rightarrow \dot{\rho} = 0 \quad \dot{v} = 0$

→ wind equation:

$$\left(v(r) - \frac{c_s^2}{v(r)} \right) v'(r) = -\phi'(r) + \frac{2c_s^2}{r}$$

→ the sonic radius:

$$\phi'(r_s) = 2c_s^2 / r_s$$

$$\rightarrow r_s \approx GM / c_s^2$$

wind parameter (NFW):

$$\psi \equiv \frac{GM_c / r_c}{c_s^2}$$

$$t_{\text{evap}} / t_{\text{dyn}} \approx 10^{\psi-1}$$

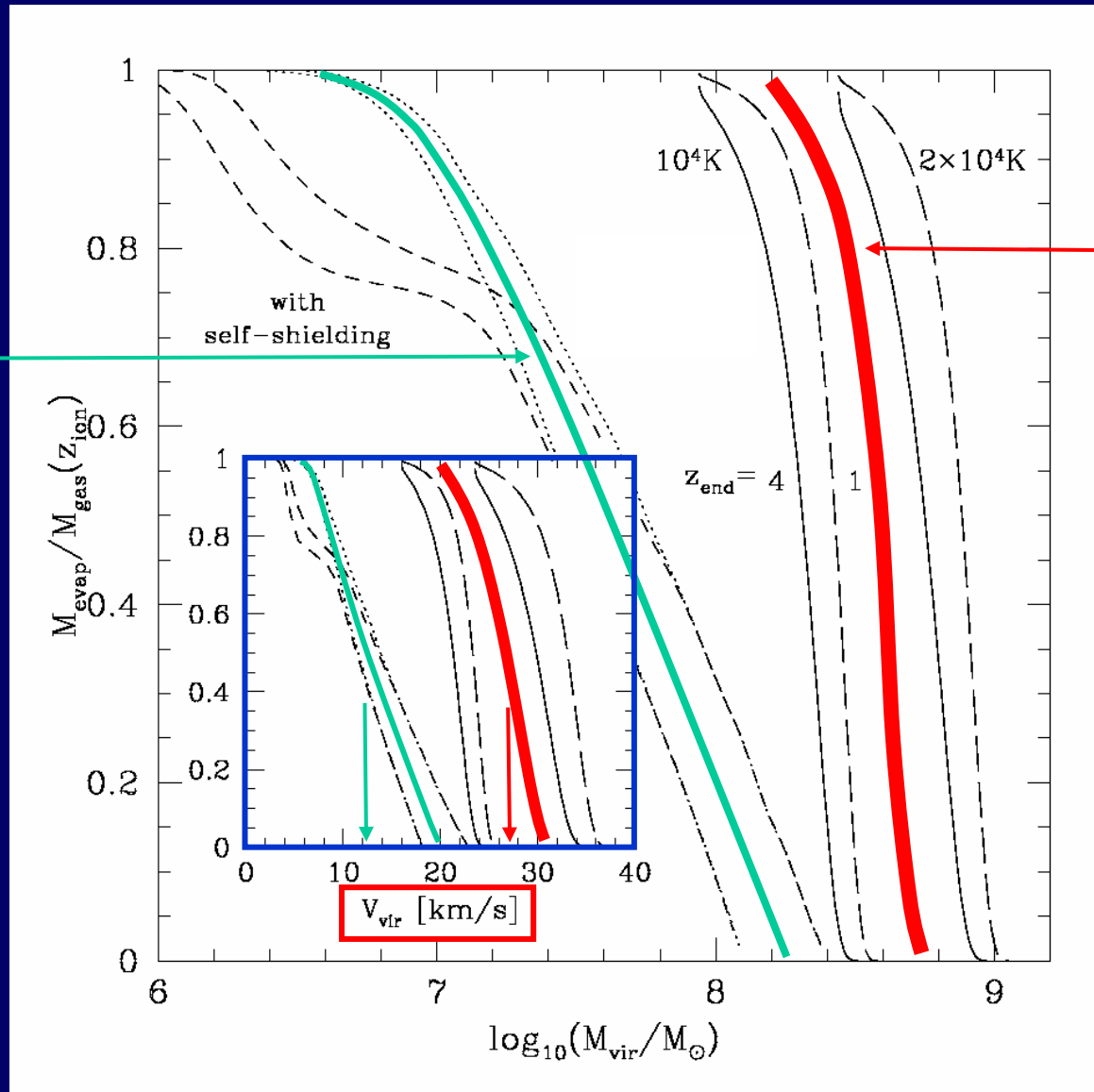
$\psi \gg 1$ tightly bound, no evaporation

$\psi > 1$ bound, but steady wind $\gg t_{\text{dyn}}$

$\psi \leq 1$ rapid evaporation $\sim t_{\text{dyn}}$

Evaporated Mass Fraction

*Barkana &
Loeb 99
instant*



*Shaviv &
Dekel 03
wind*

$z_{\text{ion}} = 8$

$z_{\text{end}} = 2$

Summary Dwarf Halos

Dark-dark halos must exist at $V < 30$ km/s

Half the photo-ionized gas evaporates by steady winds from halos of $V < 30$ km/s.

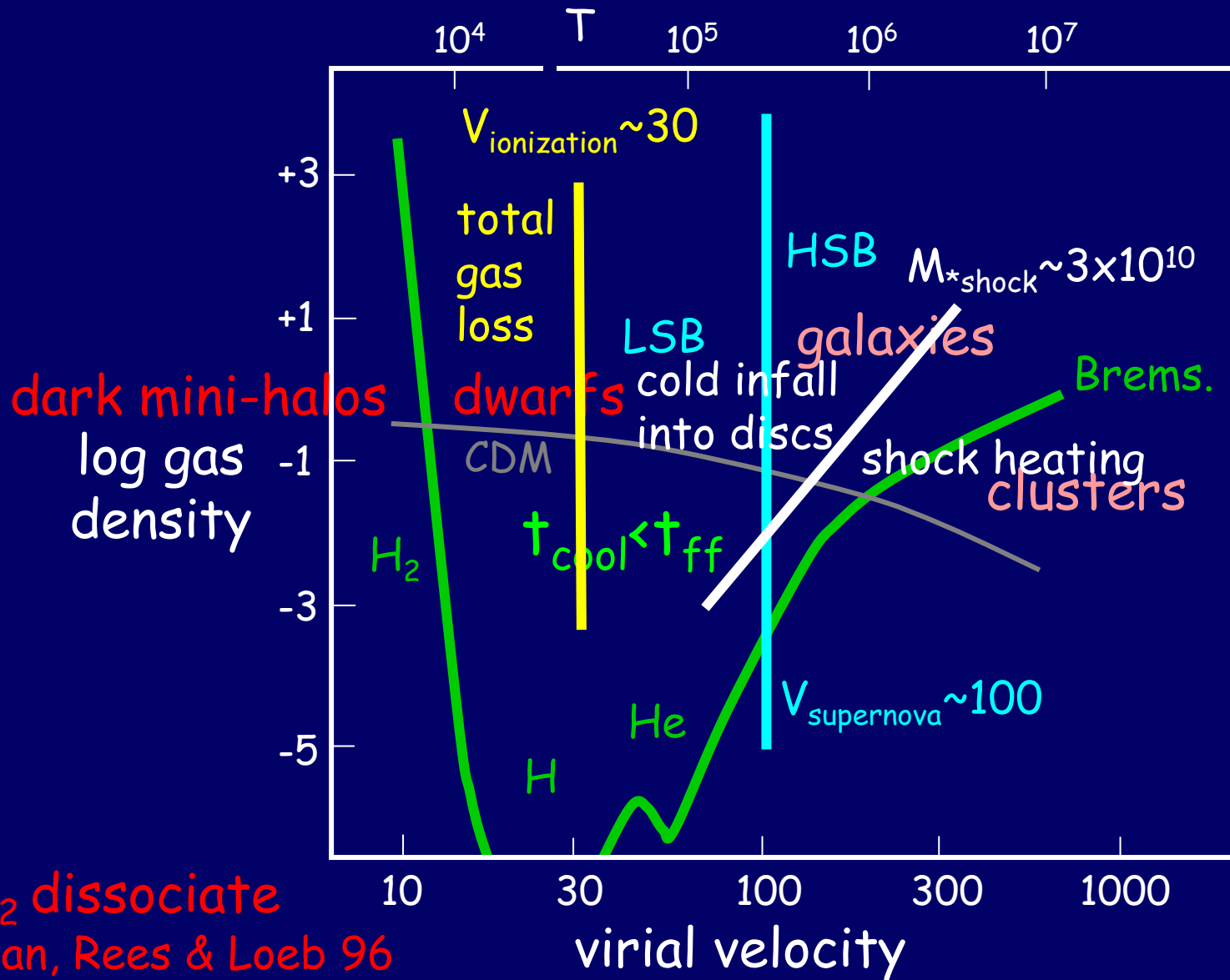
Halos in the range $10 < V < 30$ could be:

- gas-poor dSph / dE
- or totally dark

No galaxies $V < 10$ because of cooling barrier

Cooling vs Free Fall

Rees & Ostriker 77, Silk 77, White & Rees 78

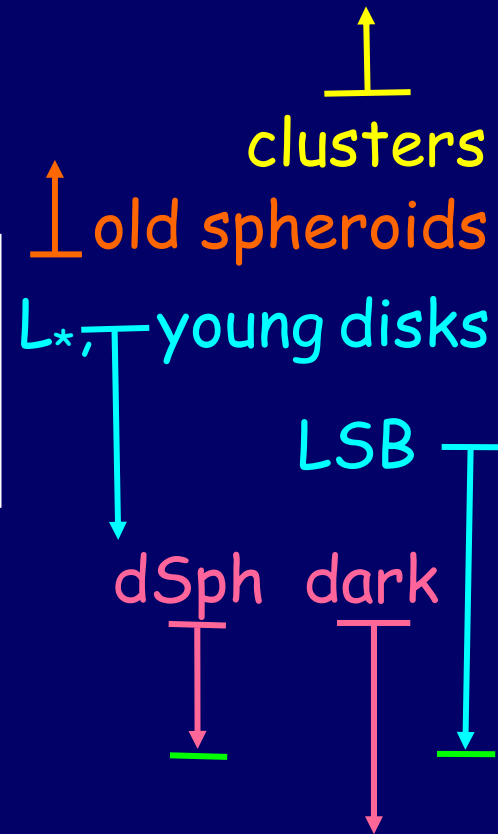


H_2 dissociate

Haiman, Rees & Loeb 96

Summary: Characteristic Scales

	V (km/s)	$M_*(M_\odot)$	$M(M_\odot)$
Cooling (Brems.)	300	2×10^{11}	10^{13}
Shock heating	100	3×10^{10}	6×10^{11}
Supernovae	100	3×10^{10}	6×10^{11}
Photoionization	30	10^8	2×10^{10}
Cooling (H)	10	3×10^5	6×10^8



Phase-Space Density & Halo Substructure

Arad & Dekel, in progress

Phase-Space Density

$$f(\vec{x}, \vec{v})$$

$$\rho(\vec{x}) = \int d\vec{v} f(\vec{x}, \vec{v})$$

Vlasov eq.

$$\partial_t f + \vec{v} \cdot \vec{\nabla}_x f - \vec{\nabla}_x \phi \cdot \nabla_v f = 0$$

Poisson eq.

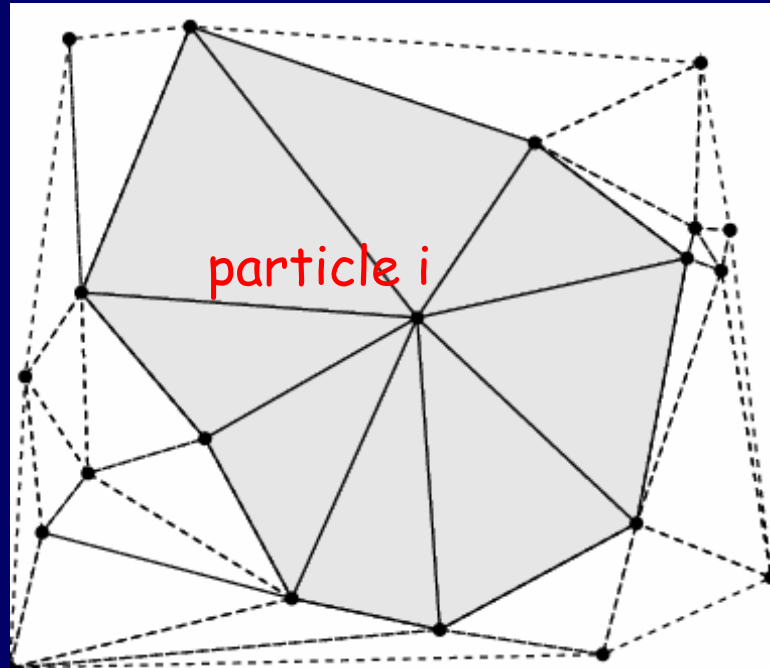
$$\phi(\vec{x}) = -G \int d\vec{x}' d\vec{v} \frac{f(\vec{x}', \vec{v})}{|\vec{x} - \vec{x}'|}$$

Distribution function of f :

$$V(f = f_0) \equiv \int d\vec{x} d\vec{v} \delta_{Dirac}[f(\vec{x}, \vec{v}, t) - f_0]$$

$V(f)df$ = volume of phase space occupied by f in the range $(f, f+df)$

Measuring $f(x,v)$ using an adaptive "grid" Delaunay Tessellation

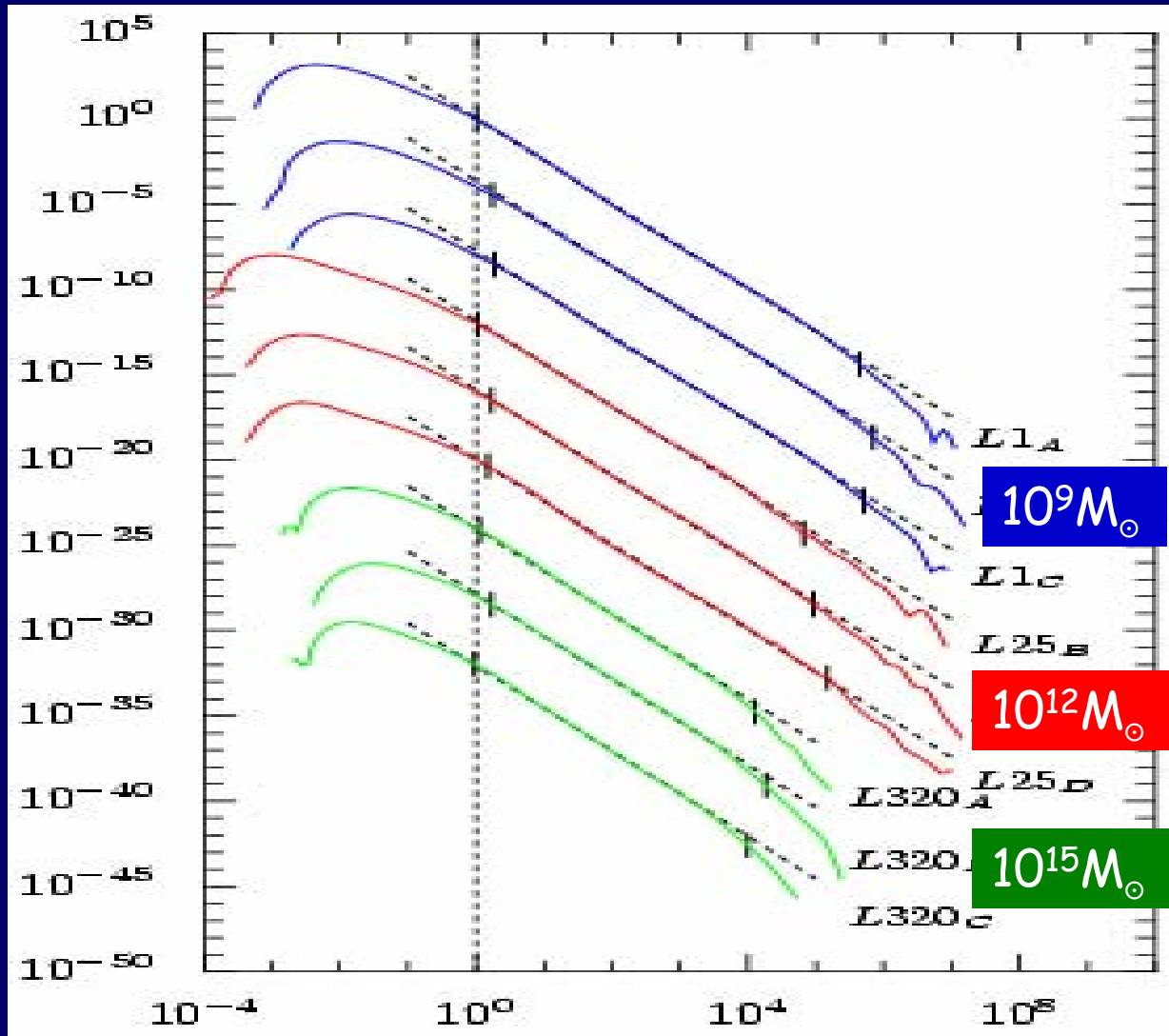


$$f_i = (d + 1) \frac{m}{V_i}$$

Arad, Dekel & Klypin

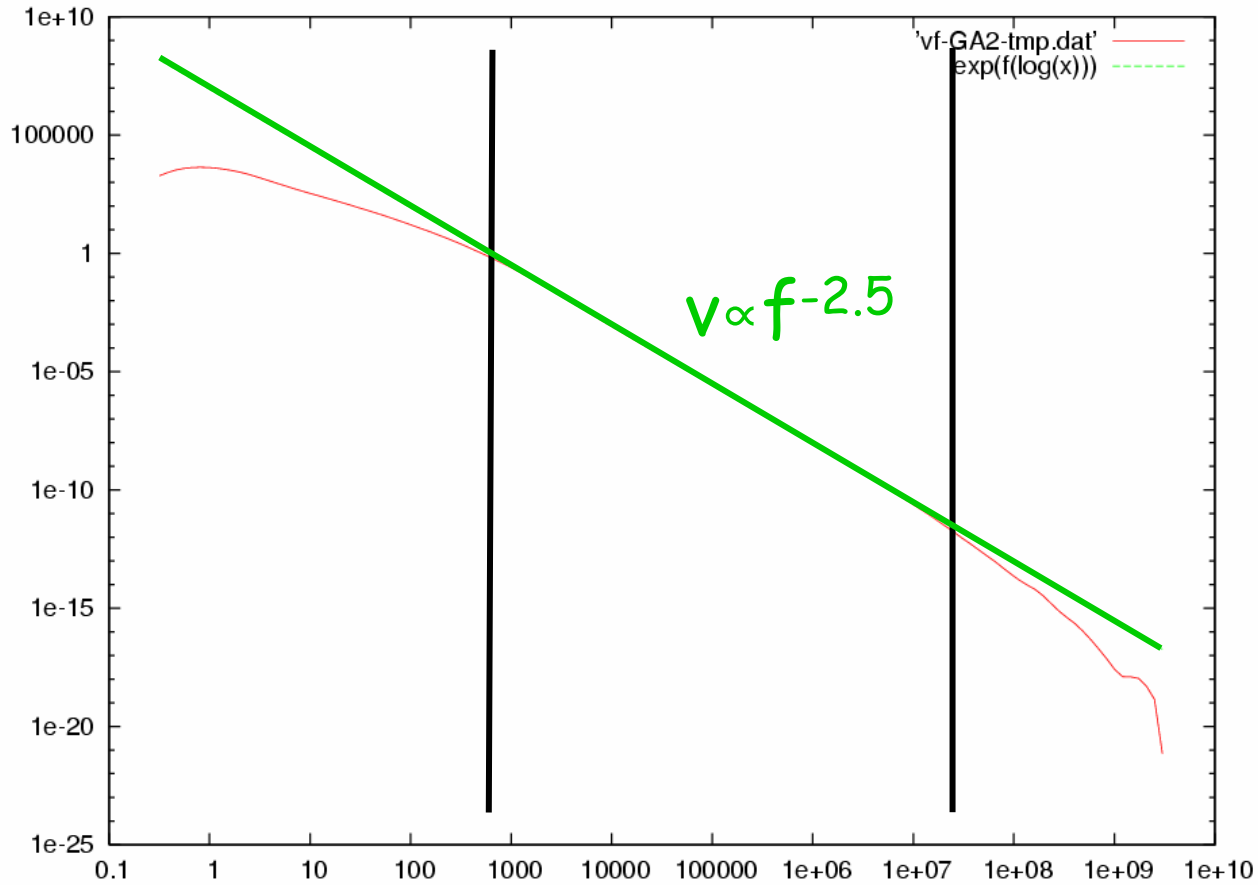
PDF of Phase-Space Density

$V(f)$



PDF of Phase-Space Density

$V(f)$



Arad, Dekel & Klypin

f

$V(f)$ related to $\rho(r)$?

if $f(\vec{x}, \vec{v}) \neq f(E)$ e.g., spherical & isotropic

$$\rho(r) \propto r^{-\alpha}, \quad V(f) \propto f^{-\beta}, \quad \beta = \frac{18 - 4\alpha}{6 - \alpha}$$

$$\alpha = 3 \leftrightarrow \beta = 2$$

$$\alpha = 2 \leftrightarrow \beta = 2.5$$

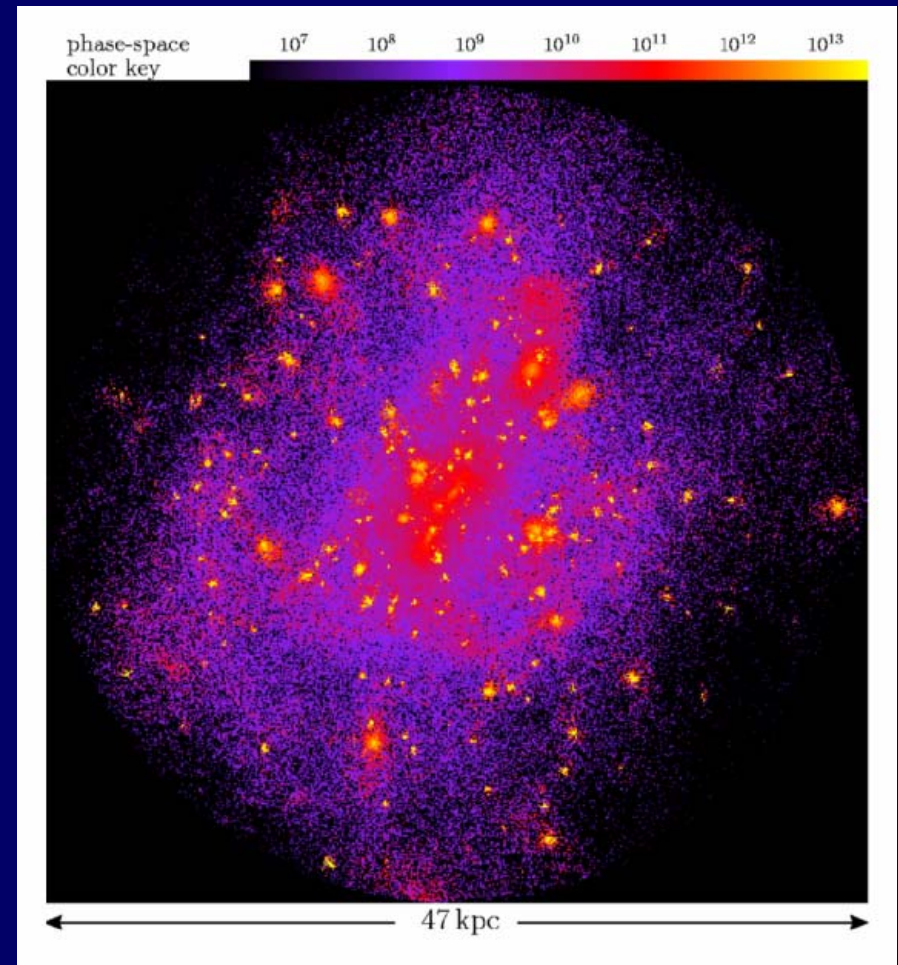
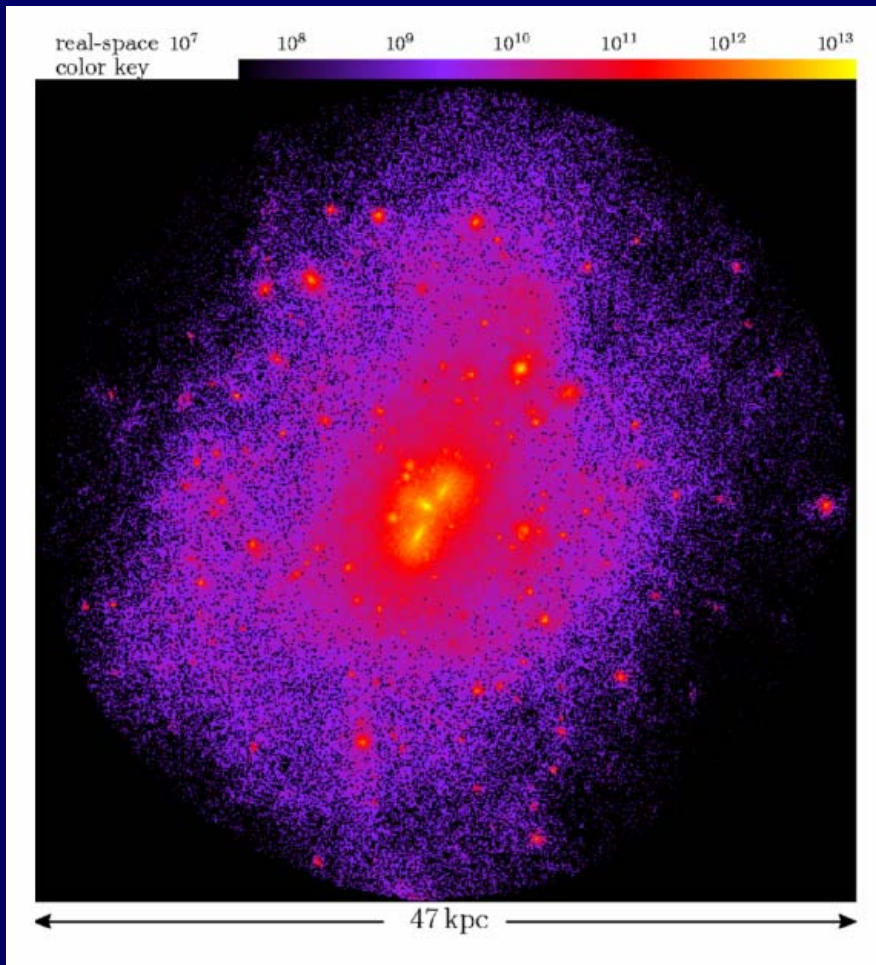
$$\alpha = 1 \leftrightarrow \beta = 2.8$$

$$\alpha = 0 \leftrightarrow \beta = 3$$

Halo Phase-Space Density

Real Density

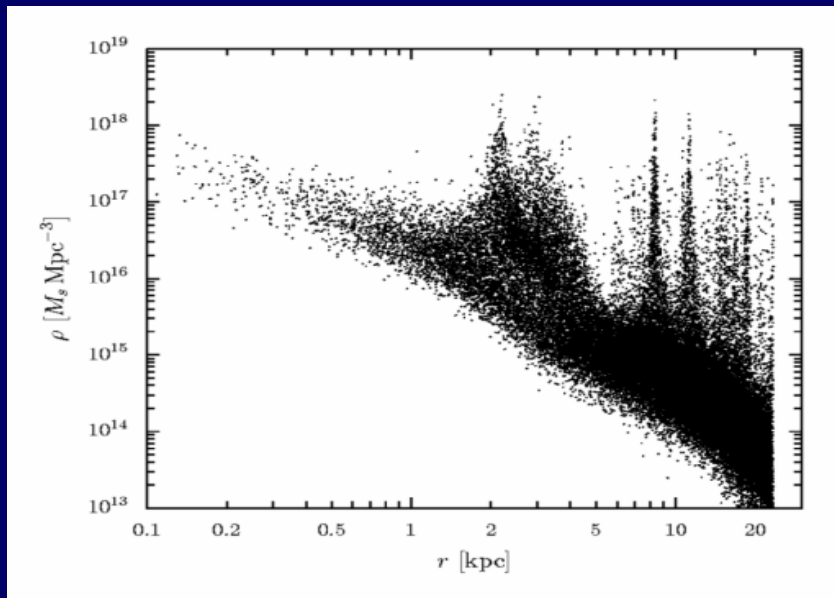
Phase-Space Density



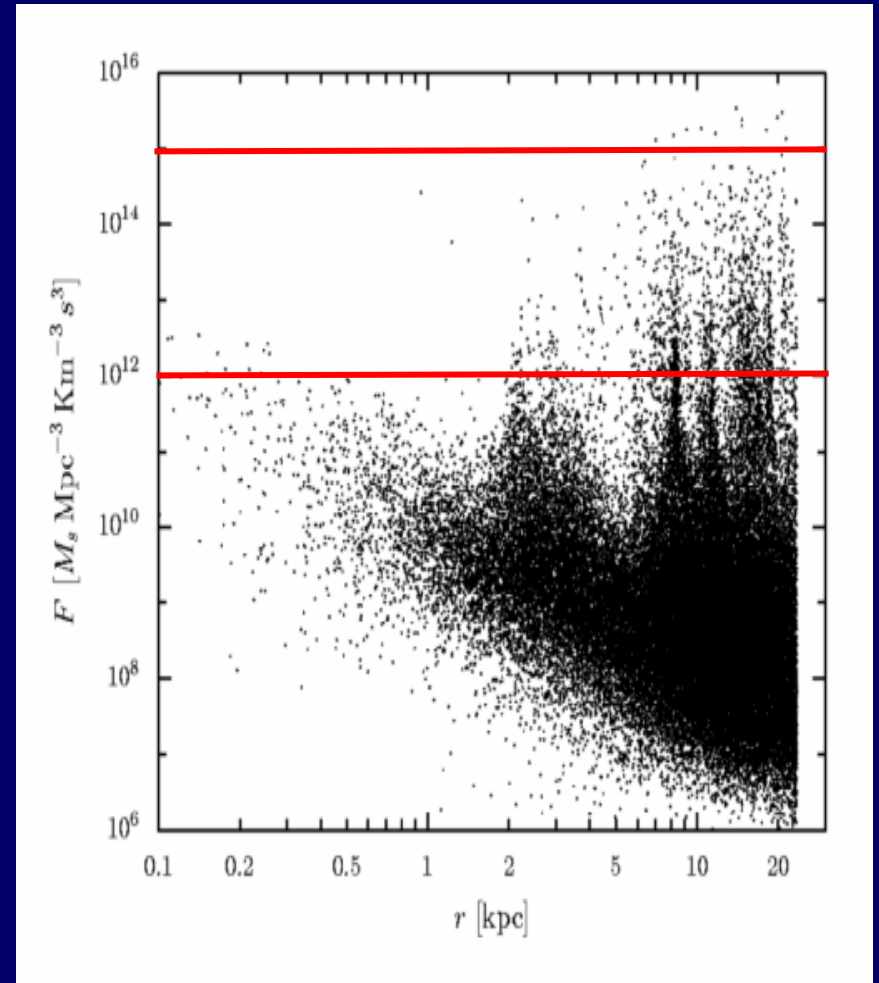
Profiles in Real Space and Phase Space

$f(r)$

$\rho(r)$



radius

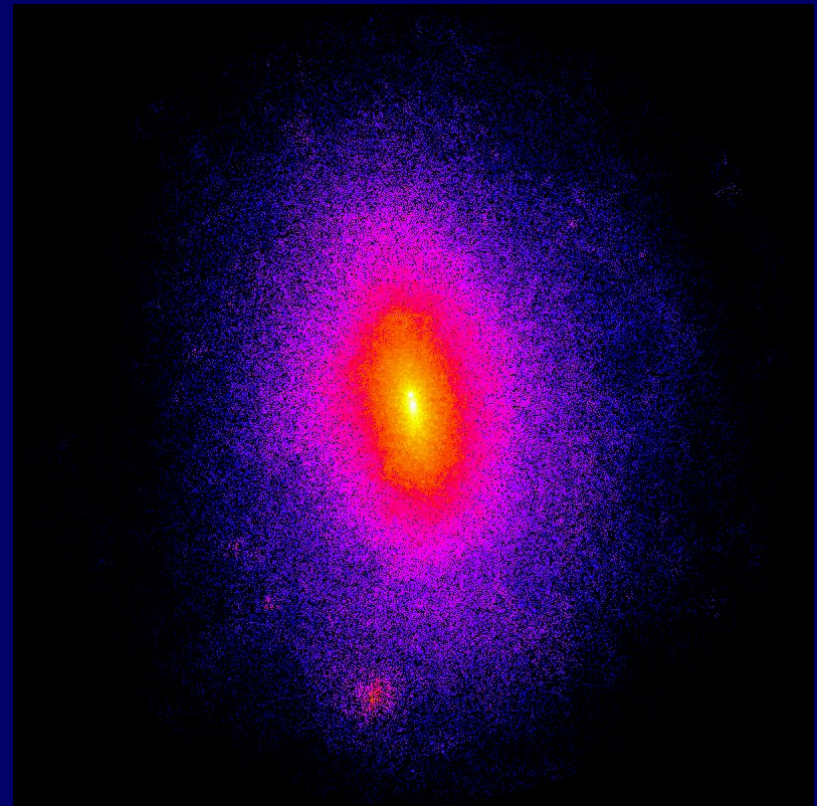
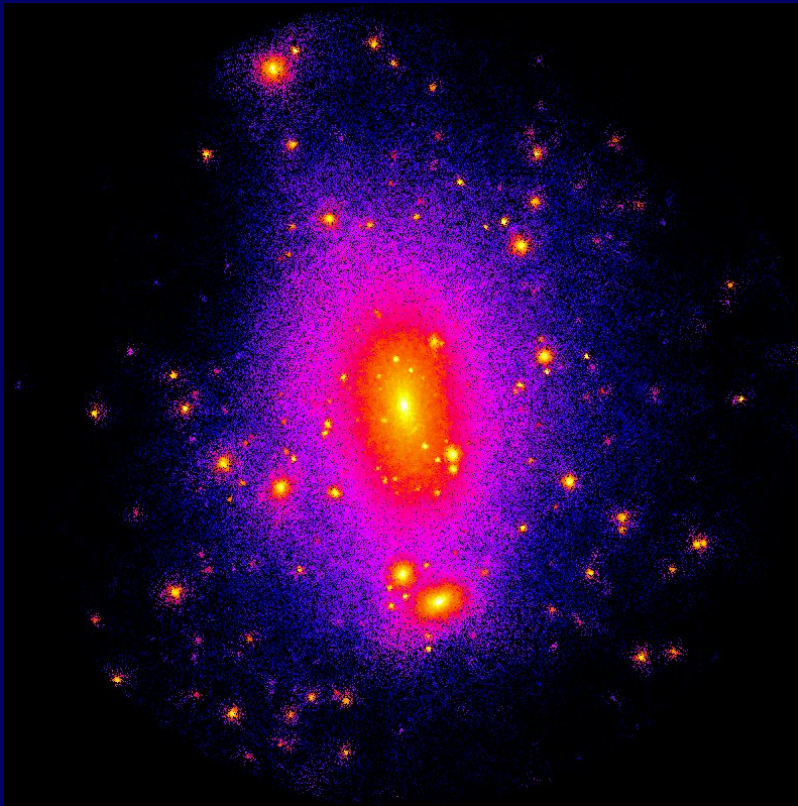


radius

Is $v(f) \propto f^{-2.5}$ determined by substructure?

Λ CDM

No short waves



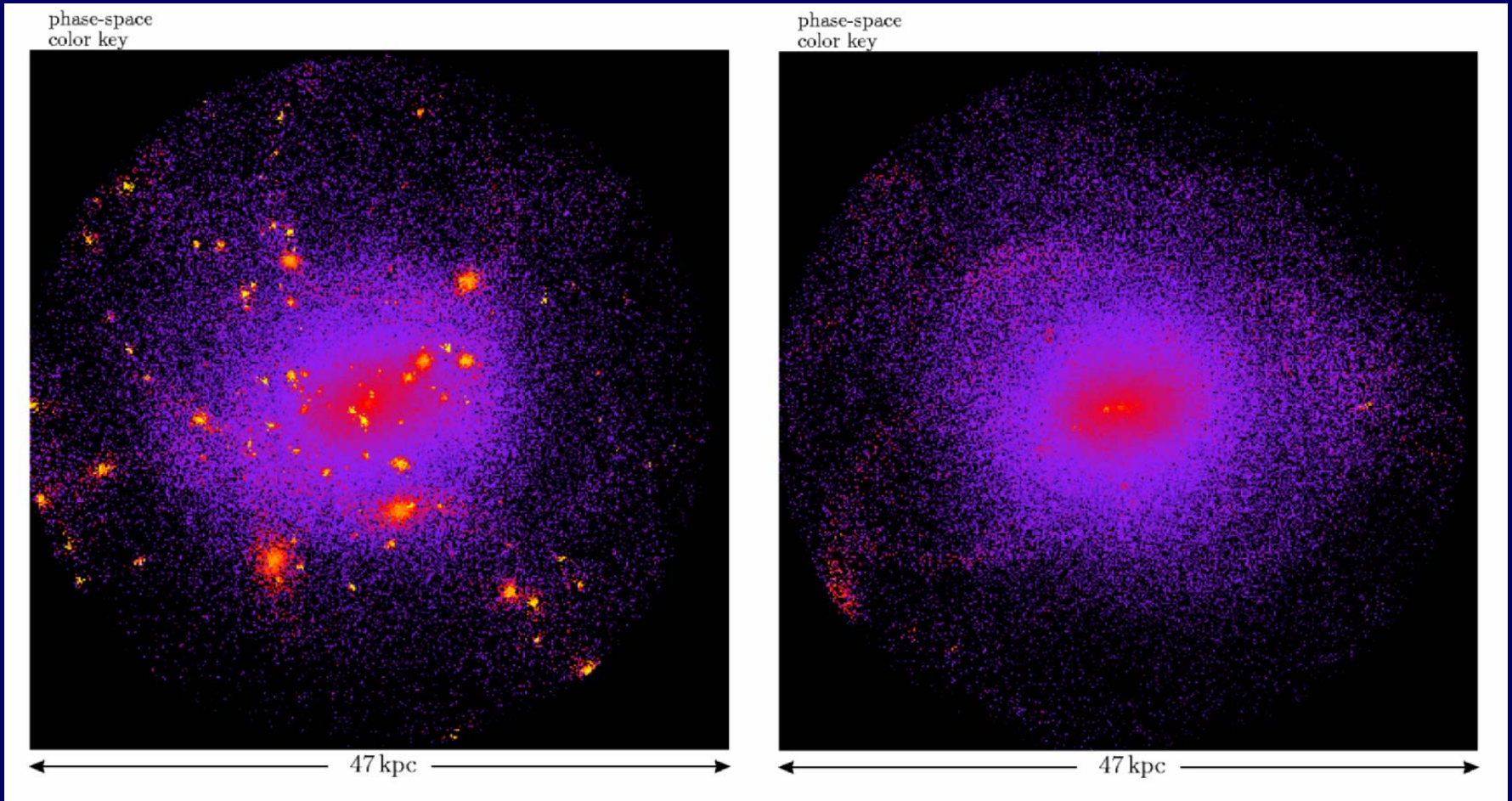
Real-Space Density

Moore et al.

Phase-Space density

Λ CDM

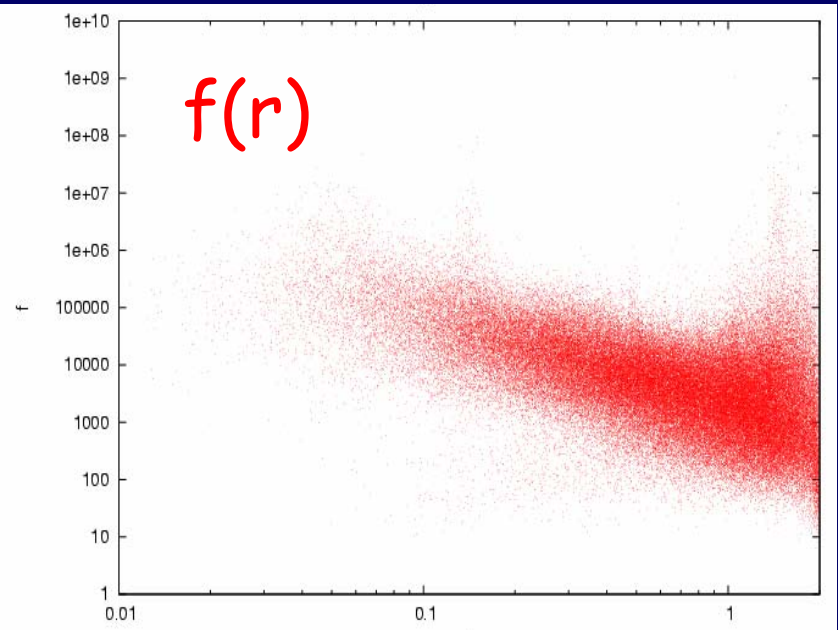
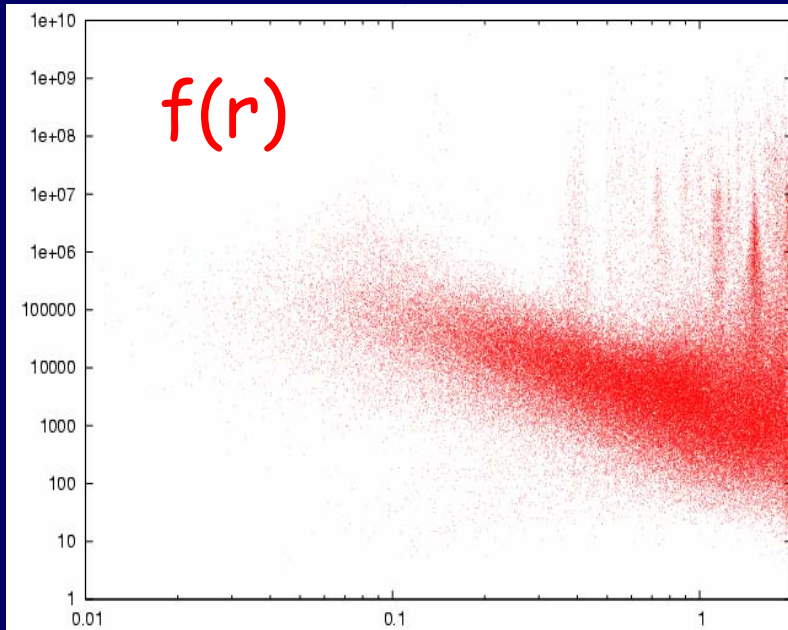
No short waves



Phase-Space Density Profile

Λ CDM

No short waves

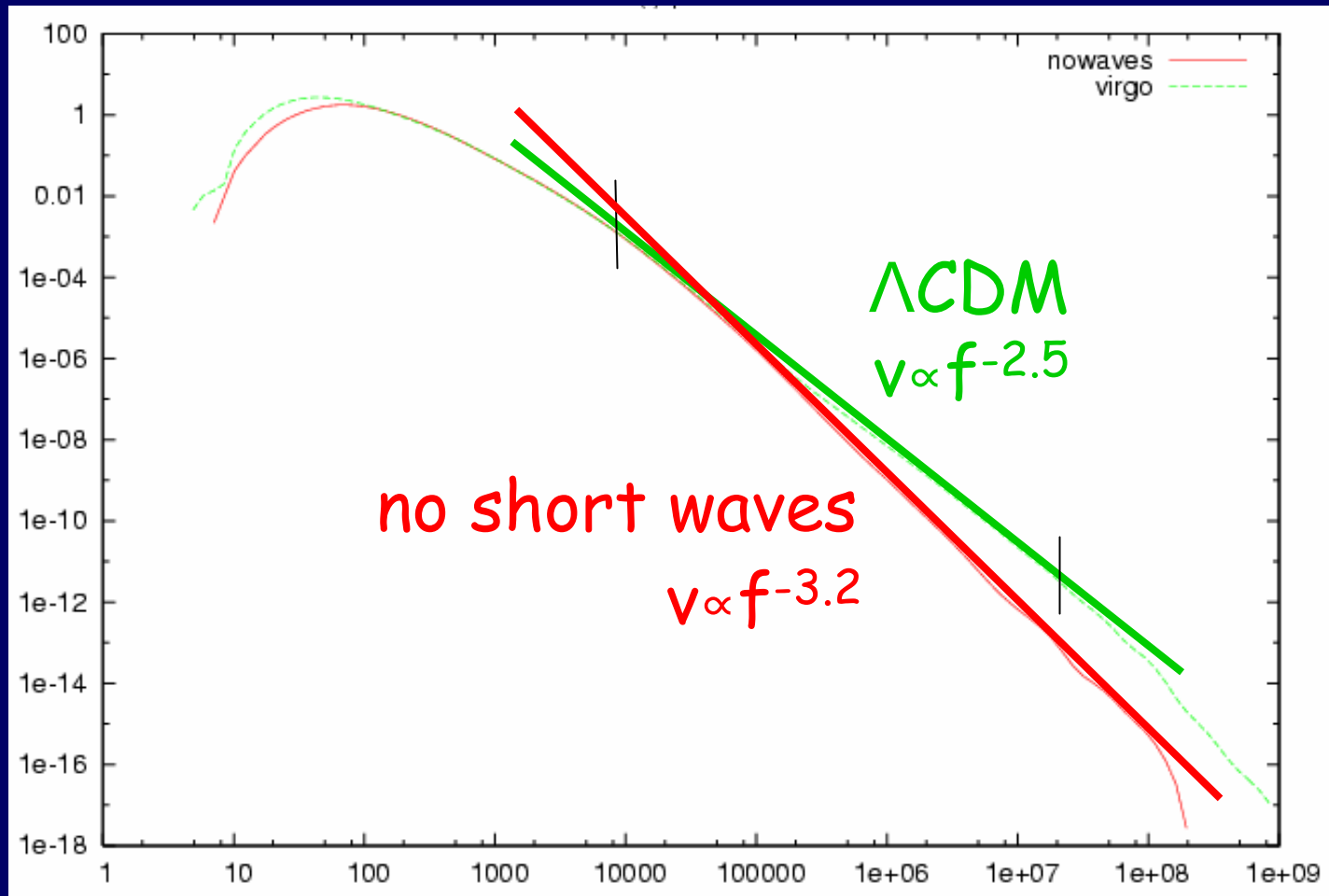


radius

radius

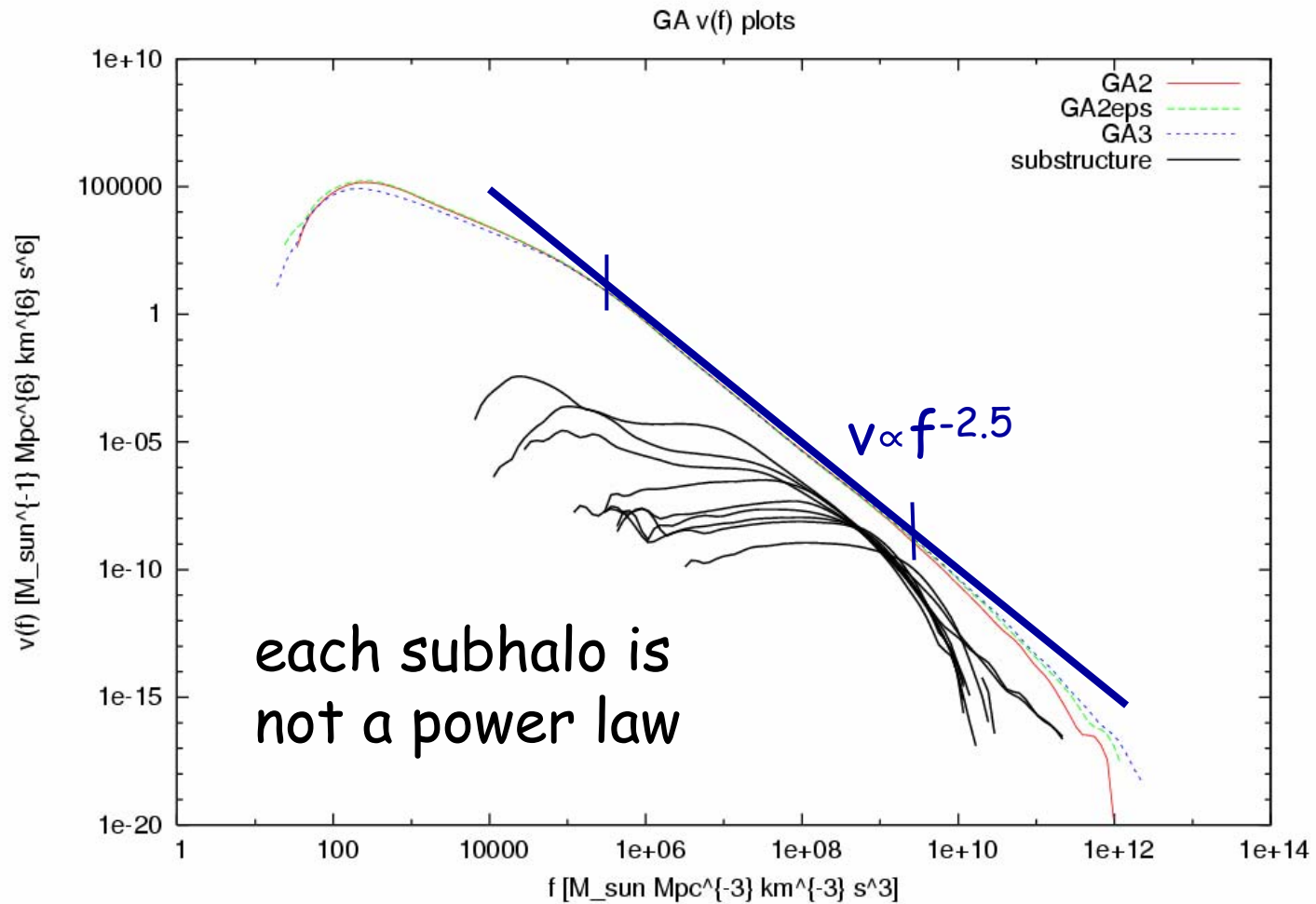
Same power law $v(f)$?

$v(f)$

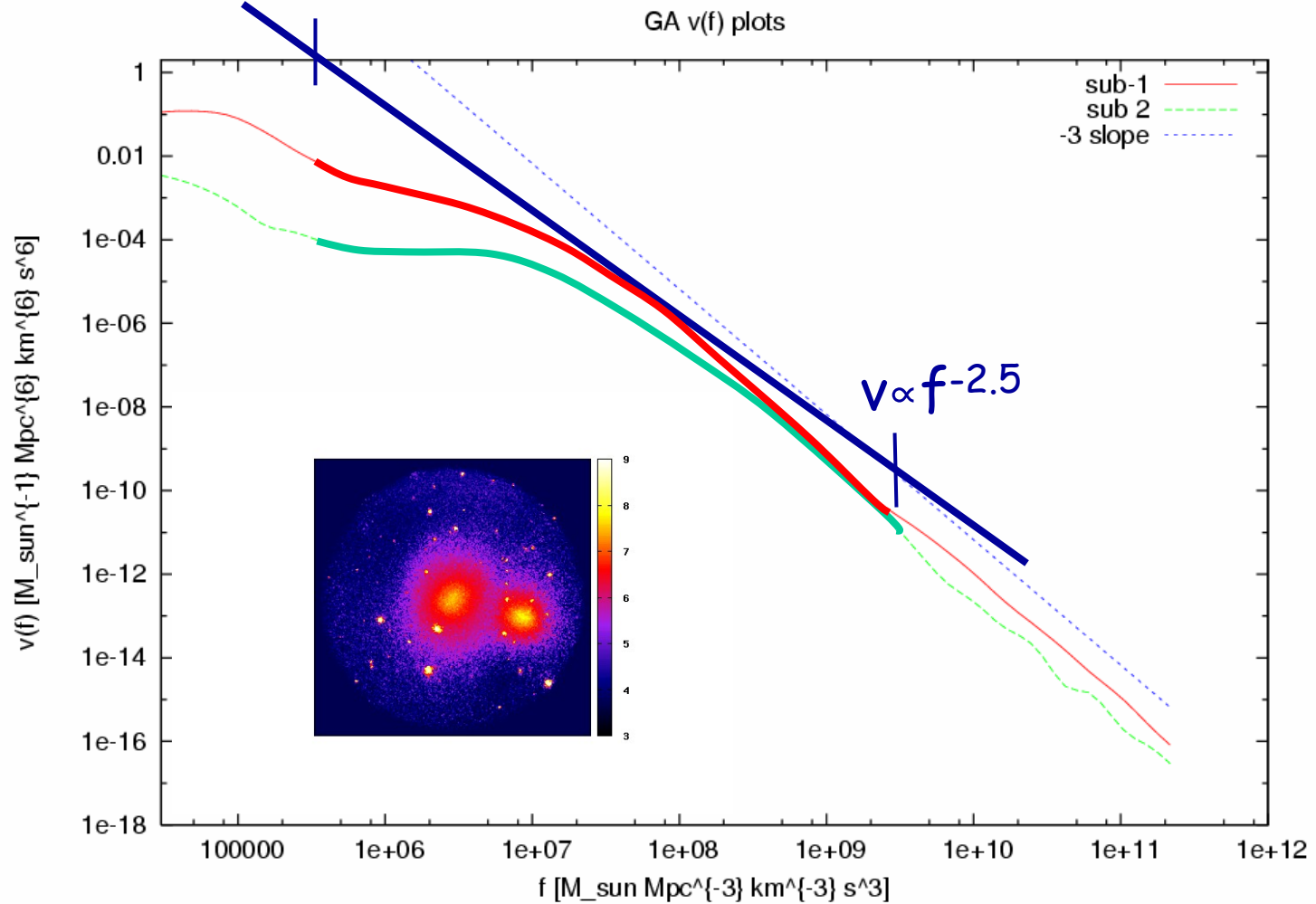


f

Additive Contribution of Subhalos

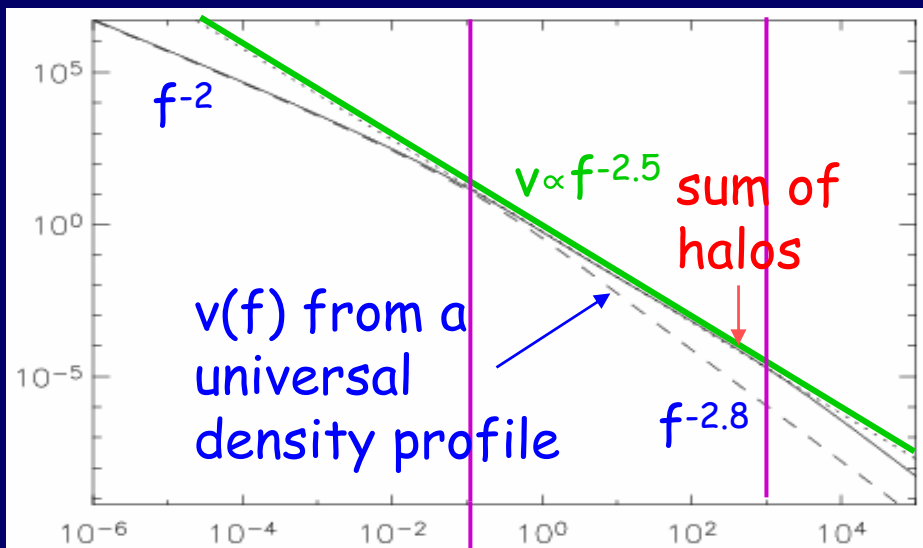


The Two Most Massive Subhalos



Adding up Sub-halos

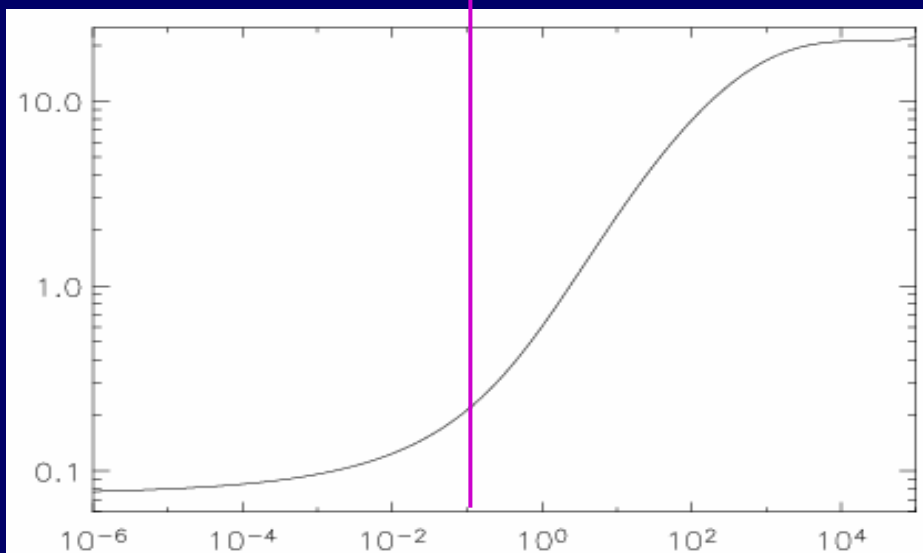
$f v(f)$



$v(f)$ from a universal density profile

f

subs/host



f

$\rho(r) \rightarrow v(f)$

halo mass function:

$\phi(m) \propto m^{-\gamma} \quad \gamma \approx 1.8$

Scaling of halos:

$\rho \propto m / r^3 = \text{const.}$

$r \propto m^{1/3} \quad \sigma \propto m^{1/3}$

Boylan-Kolchin, Ma, Arad, Dekel

Tentative Conclusions

In hierarchical clustering, robust PDF: $v(f) \propto f^{-2.5}$
doesn't depend on power-spectrum slope,
or on method of simulation

The power-law $v(f)$ is driven by substructure.
How exactly? Yet to be understood!

Phase-space density is a unique tool for
studying substructure and its evolution

Adding up small CDM halos leads to $v(f) \propto f^{-2.5}$?
How robust? How dependent on subhalo
density profile and mass function?