

# Lecture 9

## Angular Momentum

Tidal-Torque Theory

Halo spin

Angular-momentum distribution within halos

Gas Condensation and Disk Formation

The AM Problem(s)

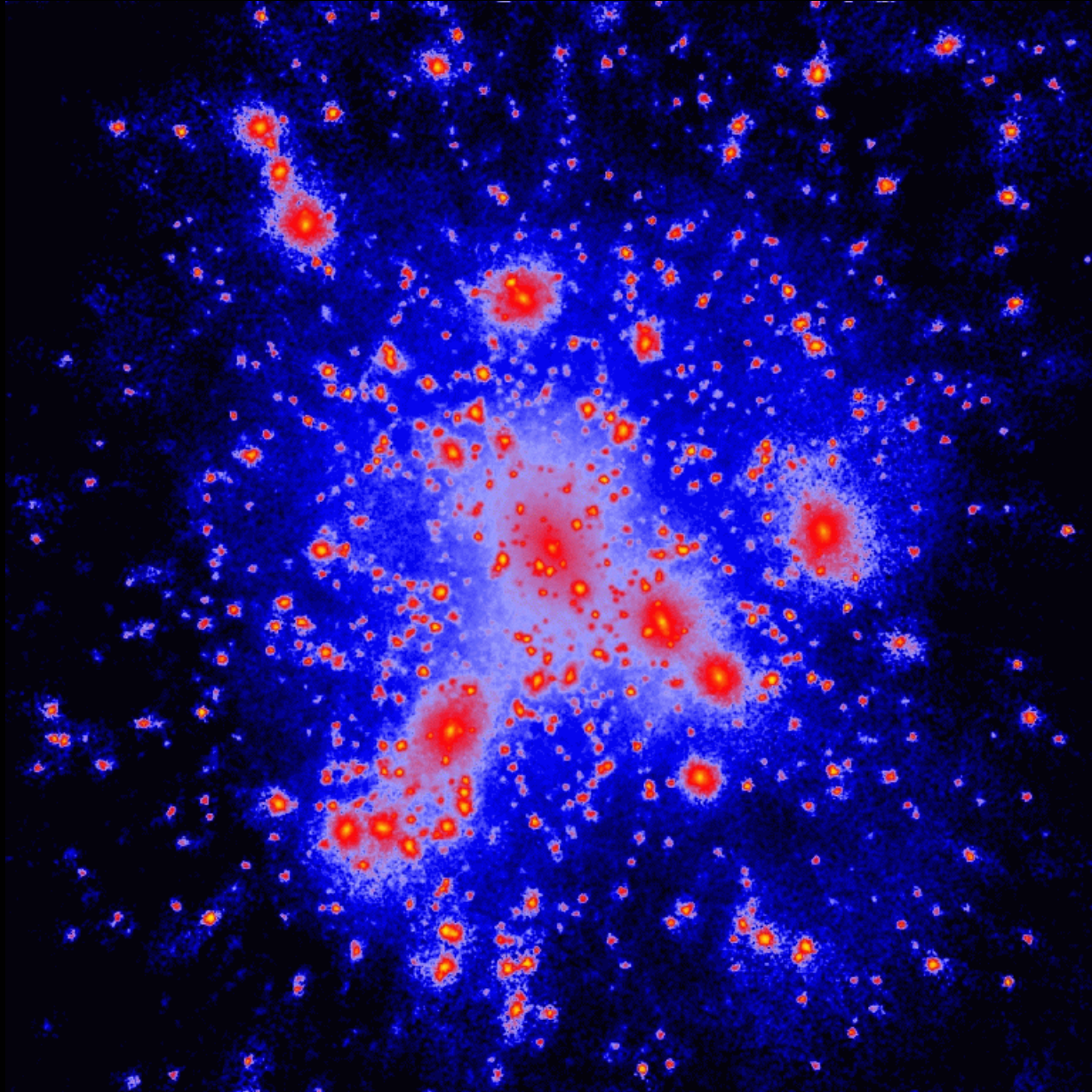
Thin disk, thick disk, bulge



# Tidal-Torque Theory (TTT)

Peebles 1976 White 1984

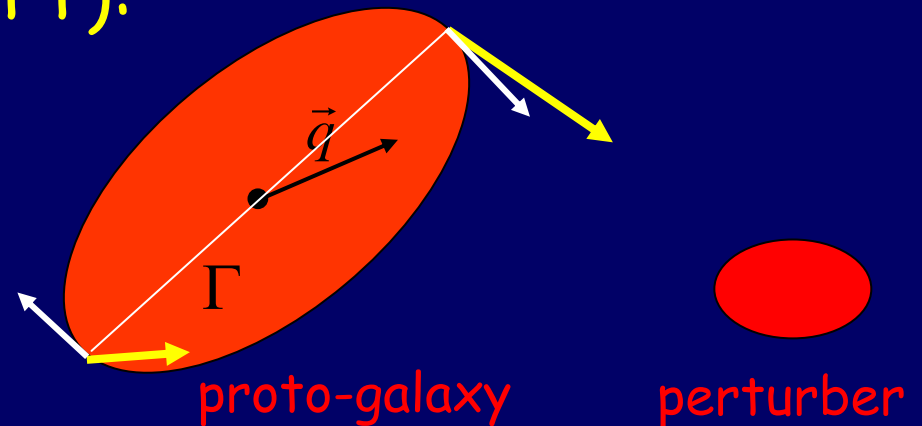
# N-body simulation of Halo Formation



# Origin of Angular Momentum

## Tidal Torque Theory (TTT):

Peebles 1976 White 1984



Result:

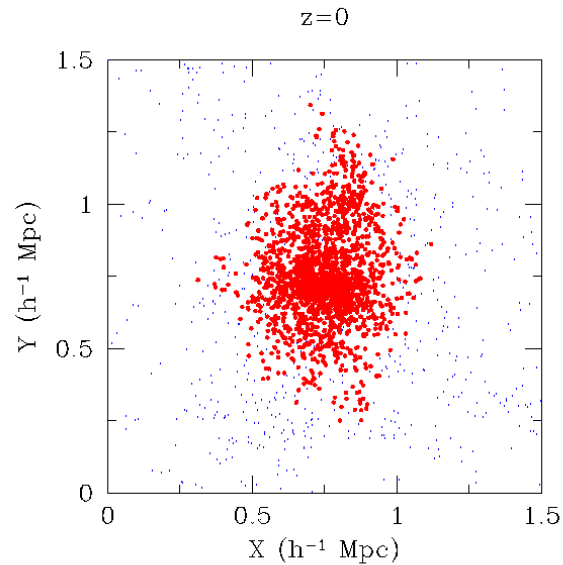
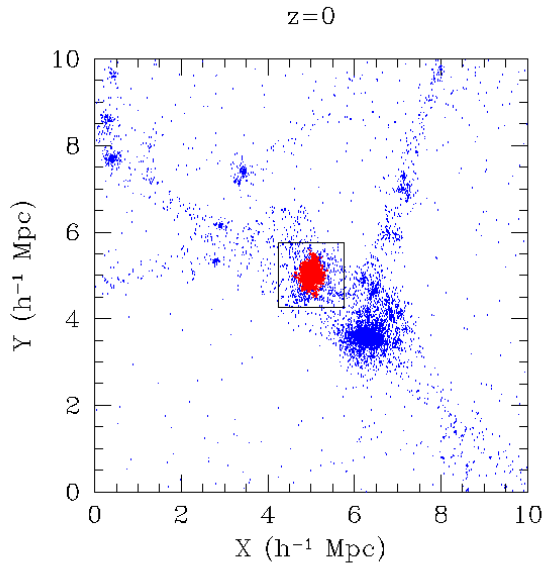
$$J_i \propto t \varepsilon_{ijk} T_{jl} I_{lk}$$

Tidal:  $T_{ij} = -\frac{\partial^2 \phi}{\partial q_i \partial q_j}$

Inertia:  $I_{ij} = \rho_0 a_0^3 \int_{\Gamma} q_i q_j d^3 q$

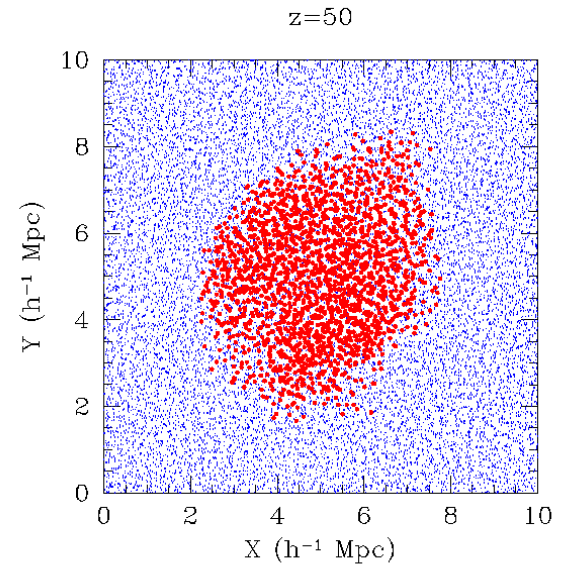
# Tidal-Torque Theory

Halo



Proto-halo:

a Lagrangian patch  $\Gamma$



$\Gamma$



# Tidal-Torque Theory

angular momentum in Eulerian patch  
comoving coordinates

$$\vec{L}(t) = \int_{\gamma \text{ Eulerian}} \rho(\vec{r}, t) [\vec{r}(t) - \vec{R}_{cm}(t)] \times [\vec{v}(t) - \vec{V}_{cm}(t)] d^3 r$$

$$\vec{x} \equiv \vec{r} / a \quad \vec{v} \equiv a \dot{\vec{x}} \quad \delta \equiv \rho / \bar{\rho}(t) - 1$$

$$\vec{L}(t) = \bar{\rho}(t) a^3(t) \int_{\gamma} [1 + \delta(\vec{x}, t)] [\vec{x}(t) - \vec{X}_{cm}(t)] \times \dot{\vec{x}} d^3 x$$

const. in m.d.

displacement from  
Lagrangian  $q$  to Eulerian  $x$

$$\vec{q} \rightarrow \vec{x} \quad \vec{x}(\vec{q}, t) = \vec{q} - \vec{S}(\vec{q}, t)$$

laminar flow

$$1 + \delta[x(q, t)] = J_{acobian}^{-1}(q, t) \rightarrow (1 + \delta) d^3 x = d^3 q$$

$$\vec{L}(t) = a^3 \bar{\rho}_0 a_0^3 \int_{\Gamma \text{ Lagrangian}} [(q - \bar{q}) + (S(q, t) - \bar{S})] \times \dot{S}(q, t) d^3 q$$

average over  $q$  in  $I$

Zel'dovich  
approximation

$$S(q, t) = -D(t) \nabla \phi(q) \quad \phi(q) = \phi_{\text{grav}}(q, t) / [4\pi G \rho(t) a^2(t) D(t)] \rightarrow \vec{S} \parallel \dot{\vec{S}}$$

$$\vec{L}(t) = -a^2(t) \dot{D}(t) \bar{\rho}_0 a_0^3 \int_{\Gamma} (q - \bar{q}) \times \nabla \phi(q) d^3 q$$

in a flat universe  $a^2 \dot{D} \propto D^{3/2} \propto t$  in EdS

2<sup>nd</sup>-order Taylor expansion  
of potential about  $q_{cm}=0$

$$\phi(\vec{q}) \approx \phi(0) + \left. \frac{\partial \phi}{\partial q_i} \right|_{\vec{q}=0} q_i + \frac{1}{2} \left. \frac{\partial^2 \phi}{\partial q_i \partial q_j} \right|_{\vec{q}=0} q_i q_j \quad \vec{q} \equiv \vec{q} - \bar{\vec{q}}$$

$$L_i(t) = a^2(t) \dot{D}(t) \varepsilon_{ijk} D_{jl} I_{lk}$$

Deformation  
tensor

$$D_{jl} \equiv - \left. \frac{\partial^2 \phi}{\partial q_j \partial q_l} \right|_{q=q_{cm}=0}$$

Inertia  
tensor

$$I_{lk} \equiv \bar{\rho}_0 a_0^3 \int_{\Gamma} q_l q_k d^3 q$$

antisymmetric  
tensor

$$\varepsilon_{ijk}$$

# Tidal-Torque Theory

$$L_i(t) = a^2(t) \dot{D}(t) \varepsilon_{ijk} T_{jl} I_{lk}$$

$$D_{jl} \equiv - \left. \frac{\partial^2 \phi}{\partial q_j \partial q_l} \right|_{q=q_{cm}=0}$$

Deformation tensor

$$I_{lk} \equiv \bar{\rho}_0 a_0^3 \int_{\Gamma} q_l q_k d^3 q$$

Inertia tensor

$\varepsilon_{ijk}$

antisymmetric

Tidal tensor = Shear tensor

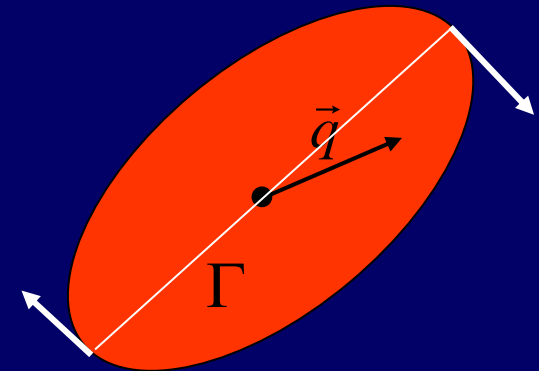
$$T_{ij} \equiv D_{ij} - D_{ii} \delta_{ij} / 3$$

Only the trace-less part contributes

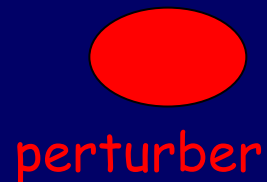
Quadrupolar Inertia

$$I_{ij} - I_{ii} \delta_{ij} / 3$$

L by gravitational coupling of  
Quadrupole moment of  $\Gamma$  with  
Tidal field from neighboring fluctuations  
→ T and I must be misaligned.



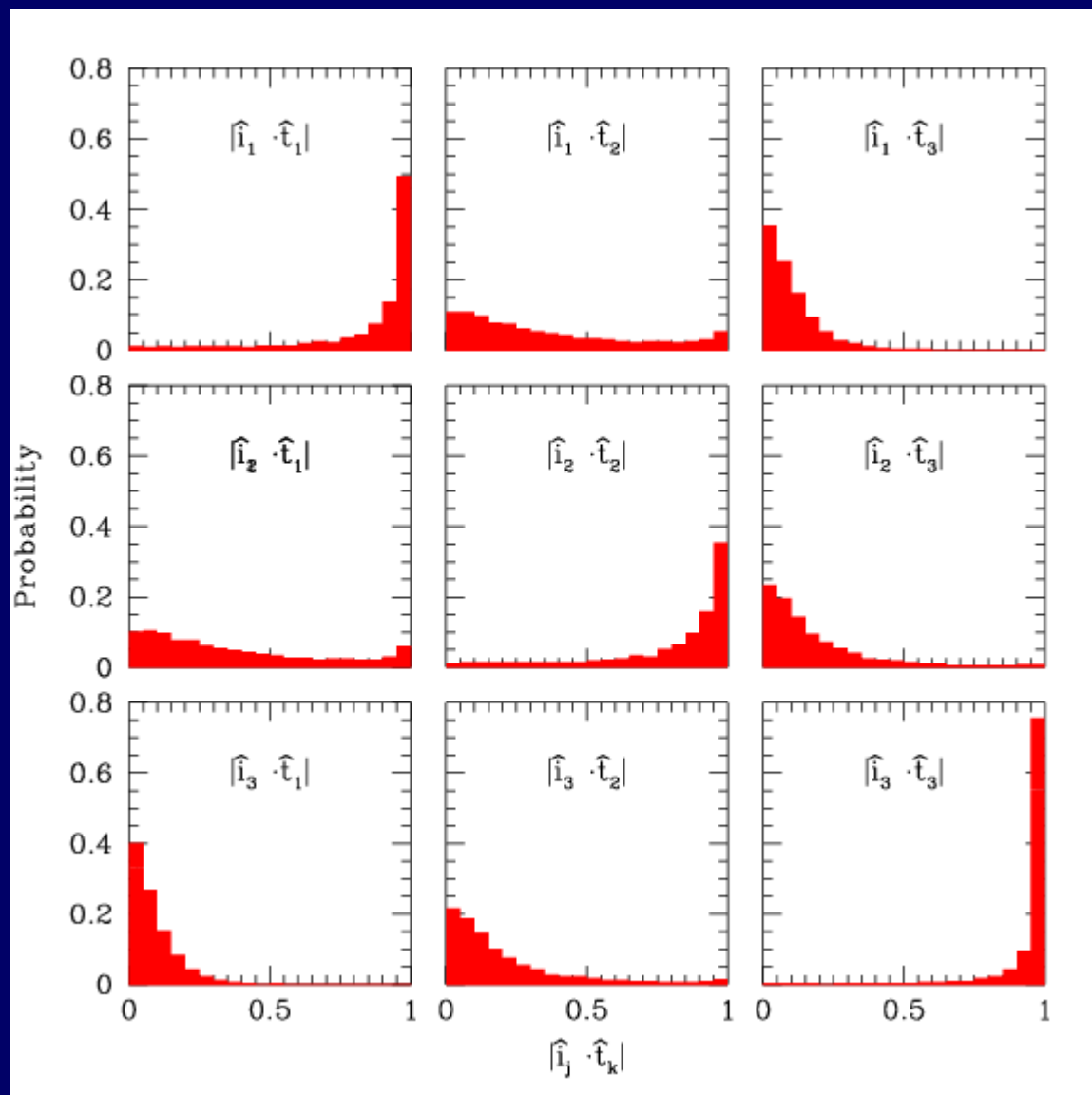
$L \propto t$  till ~turnaround





# TTT vs Simulations

(Porciani, Dekel & Hoffman 2002)



Alignment of T and I:  
Spin originates from the  
residual misalignment.

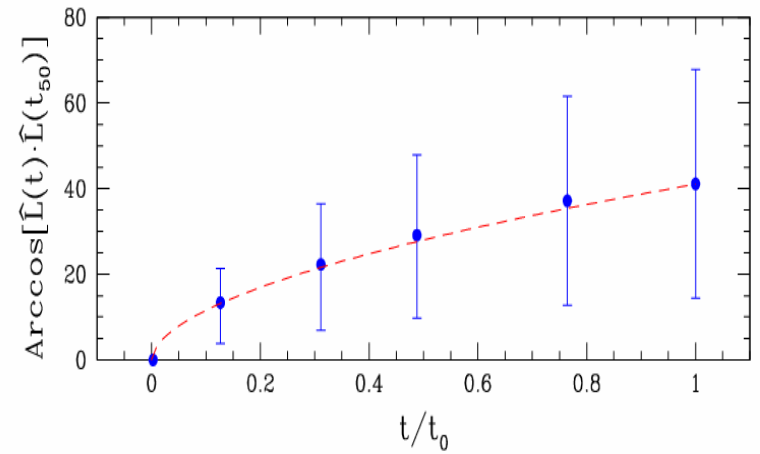
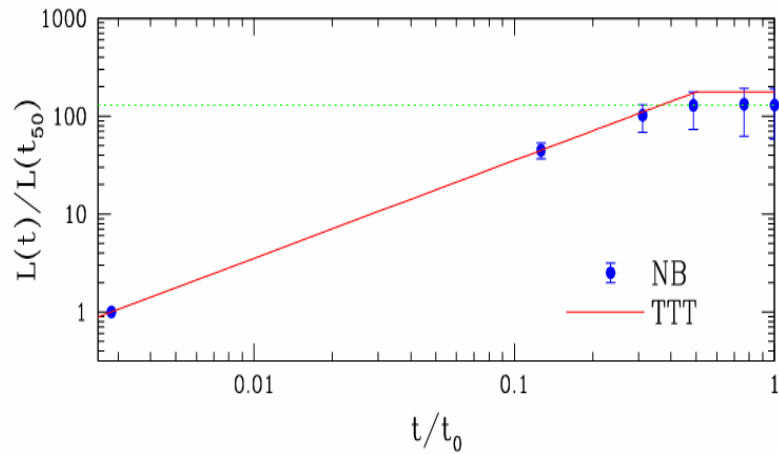
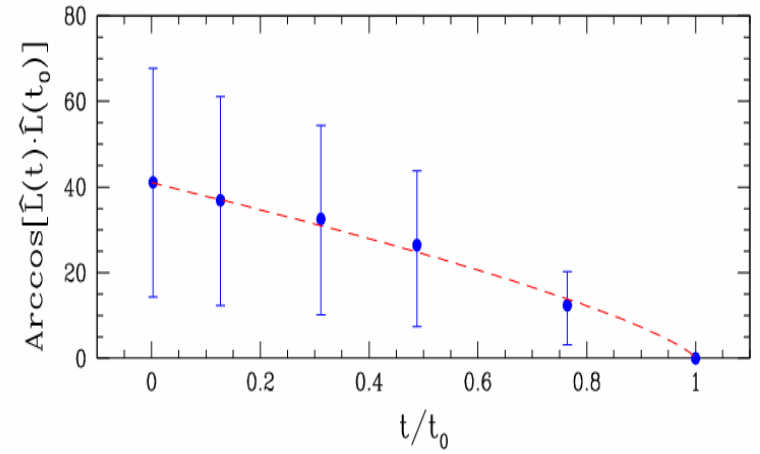
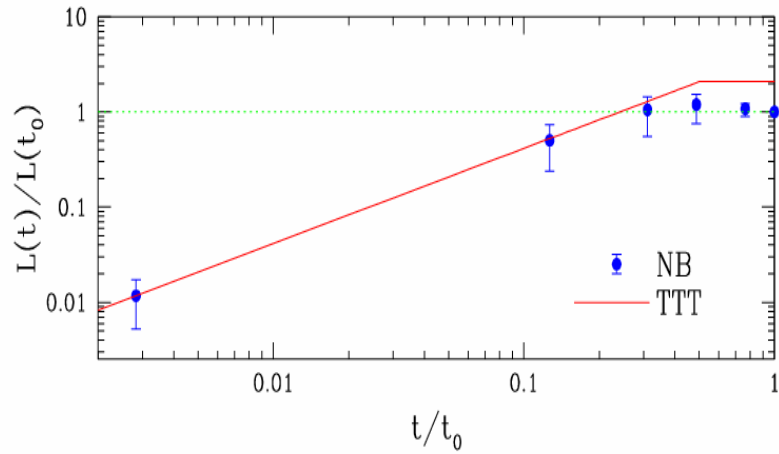
→ Small spin !

# TTT vs. Simulations: Amplitude Growth Rate

Porciani, Dekel & Hoffman 02

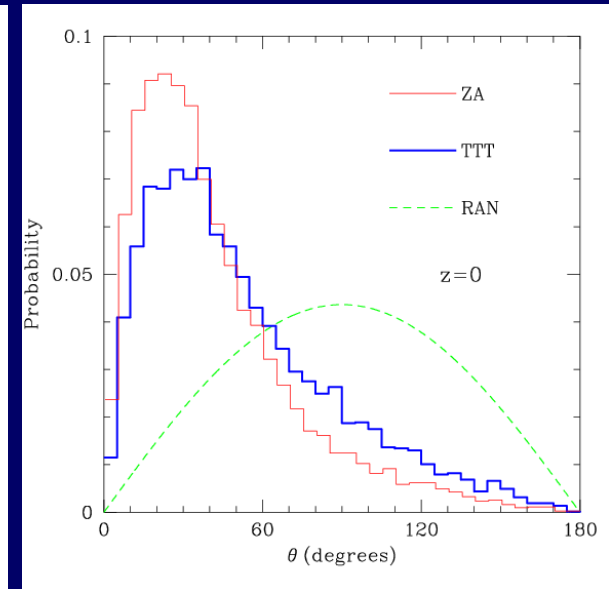
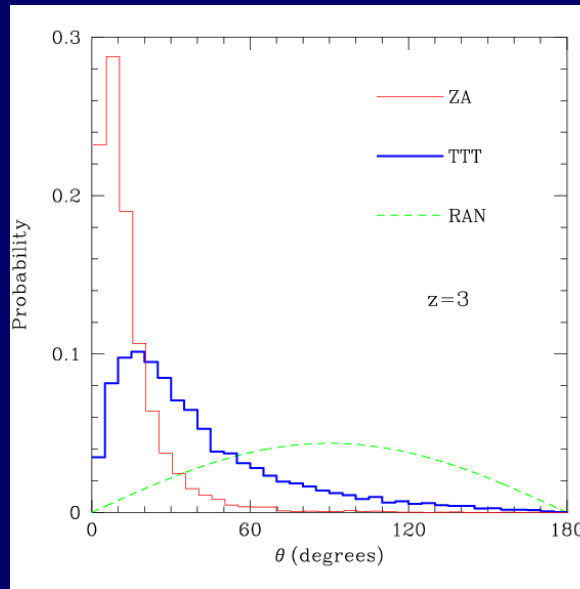
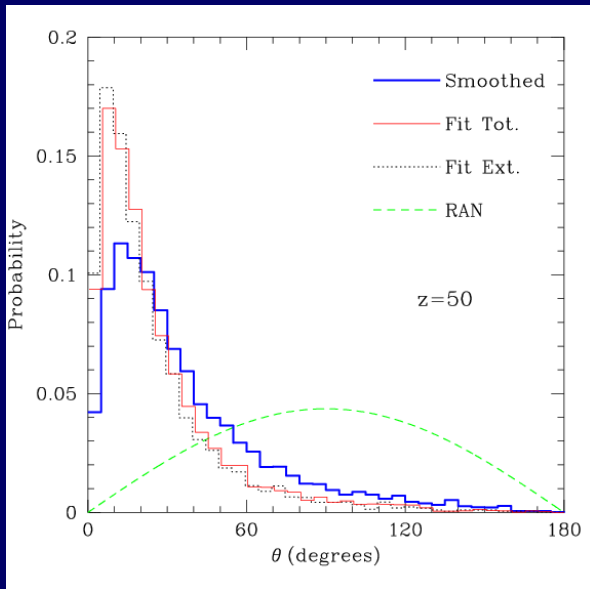
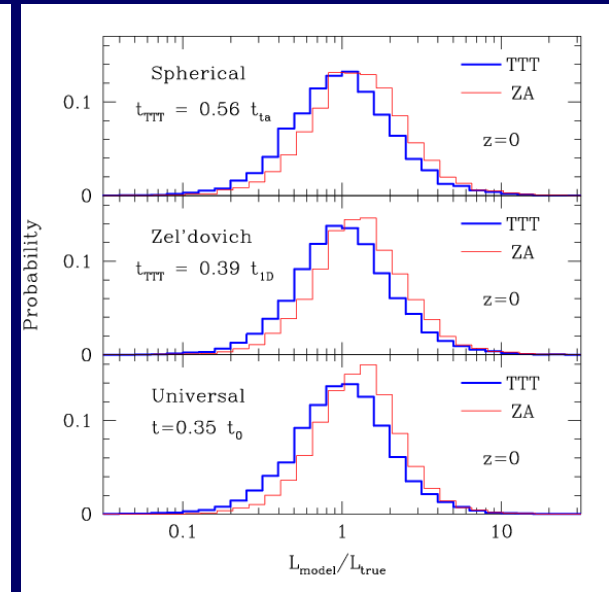
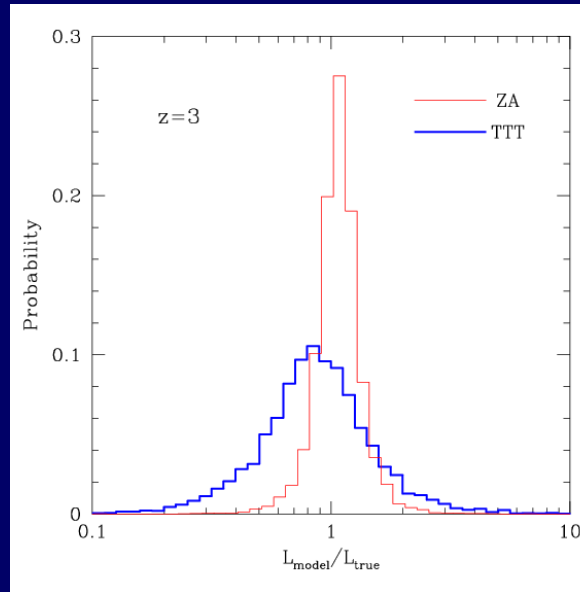
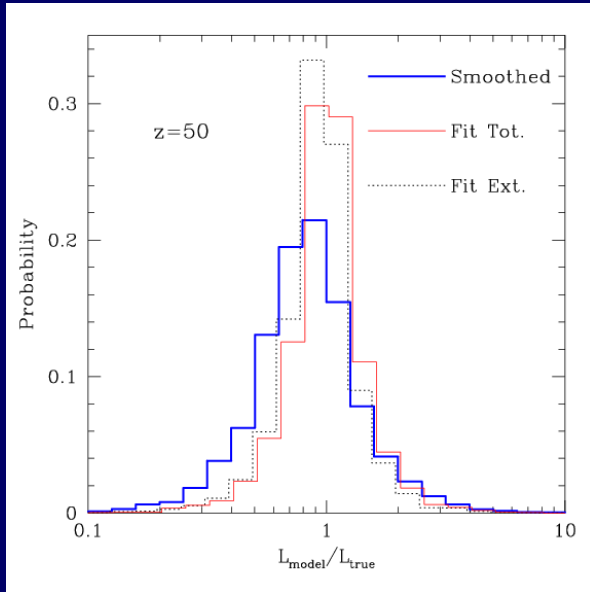
Amplitude

Direction



# TTT vs Simulations: Scatter

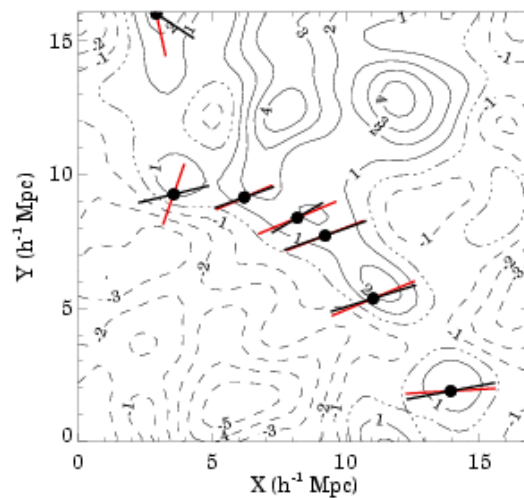
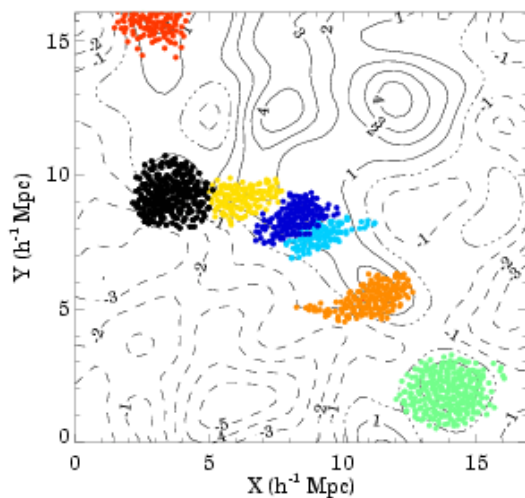
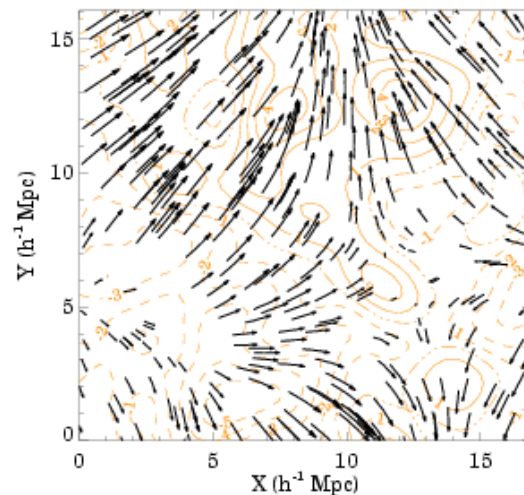
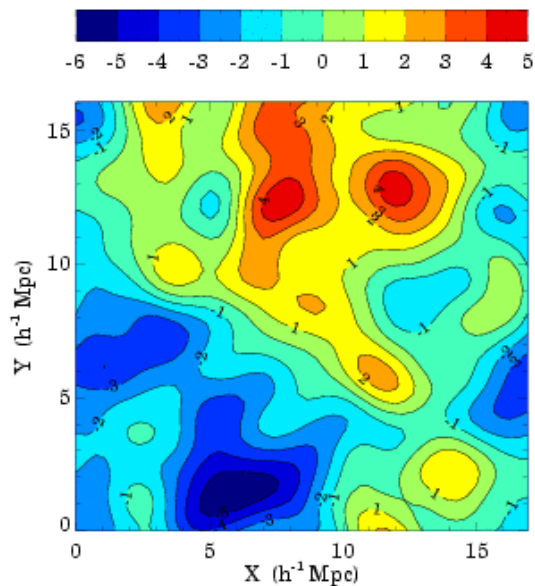
(Porciani, Dekel & Hoffman 2002)



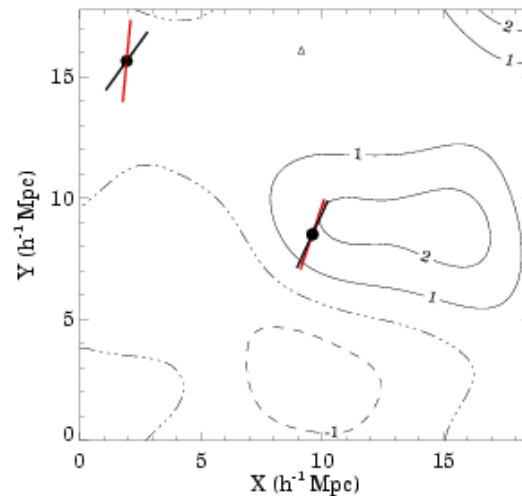
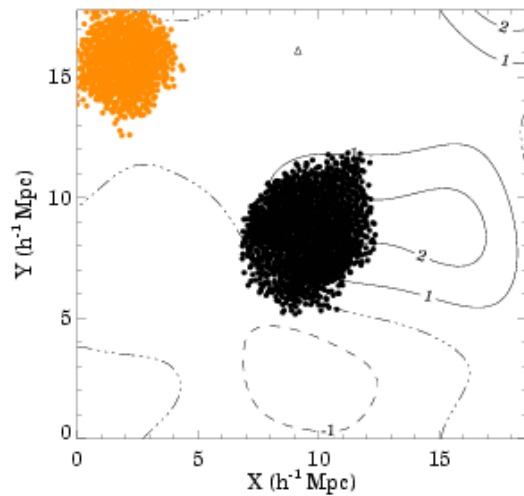
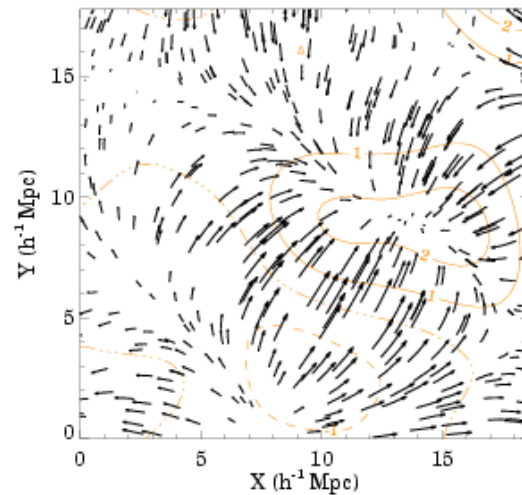
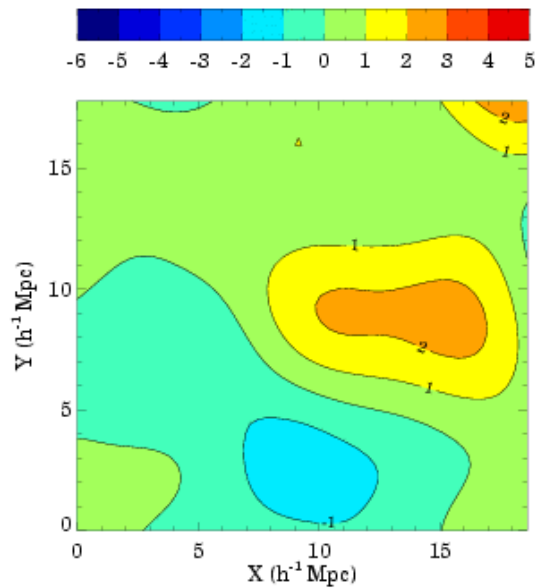
TTT predicts the spin amplitude  
to within a factor of  $\sim 2$ ,

but it is not a very reliable  
predictor of spin direction.

# Alignment of I and T: Protohalos and Filaments

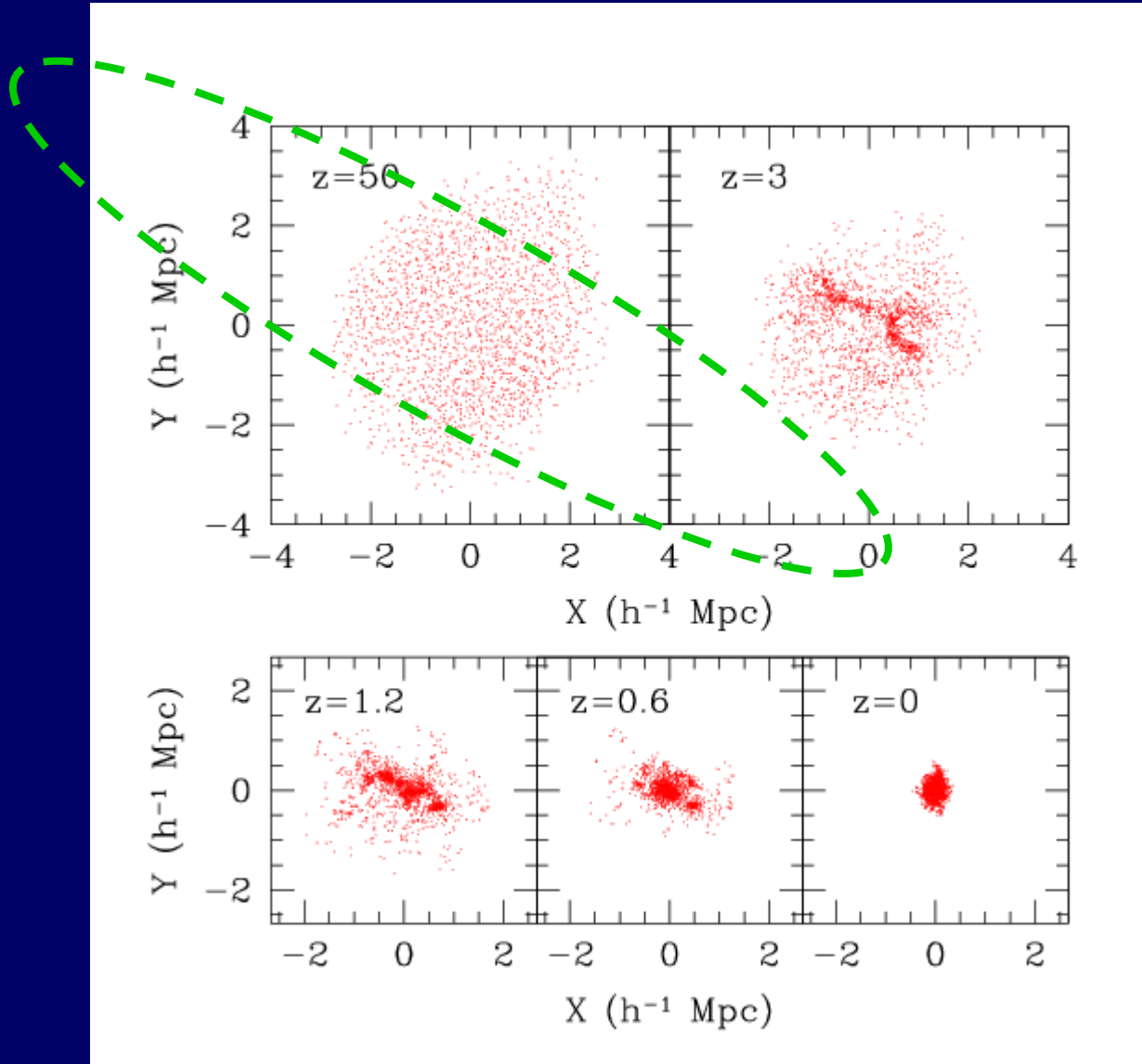


# Alignment of I and T: Protohalos and Filaments





# Stages in Halo Formation



# Spin axis and Large-Scale Structure

TTT:

$$J_x = \frac{\partial^2 \phi}{\partial y \partial z} (I_{yy} - I_{zz})$$

$$J_y = \frac{\partial^2 \phi}{\partial x \partial z} (I_{xx} - I_{zz})$$

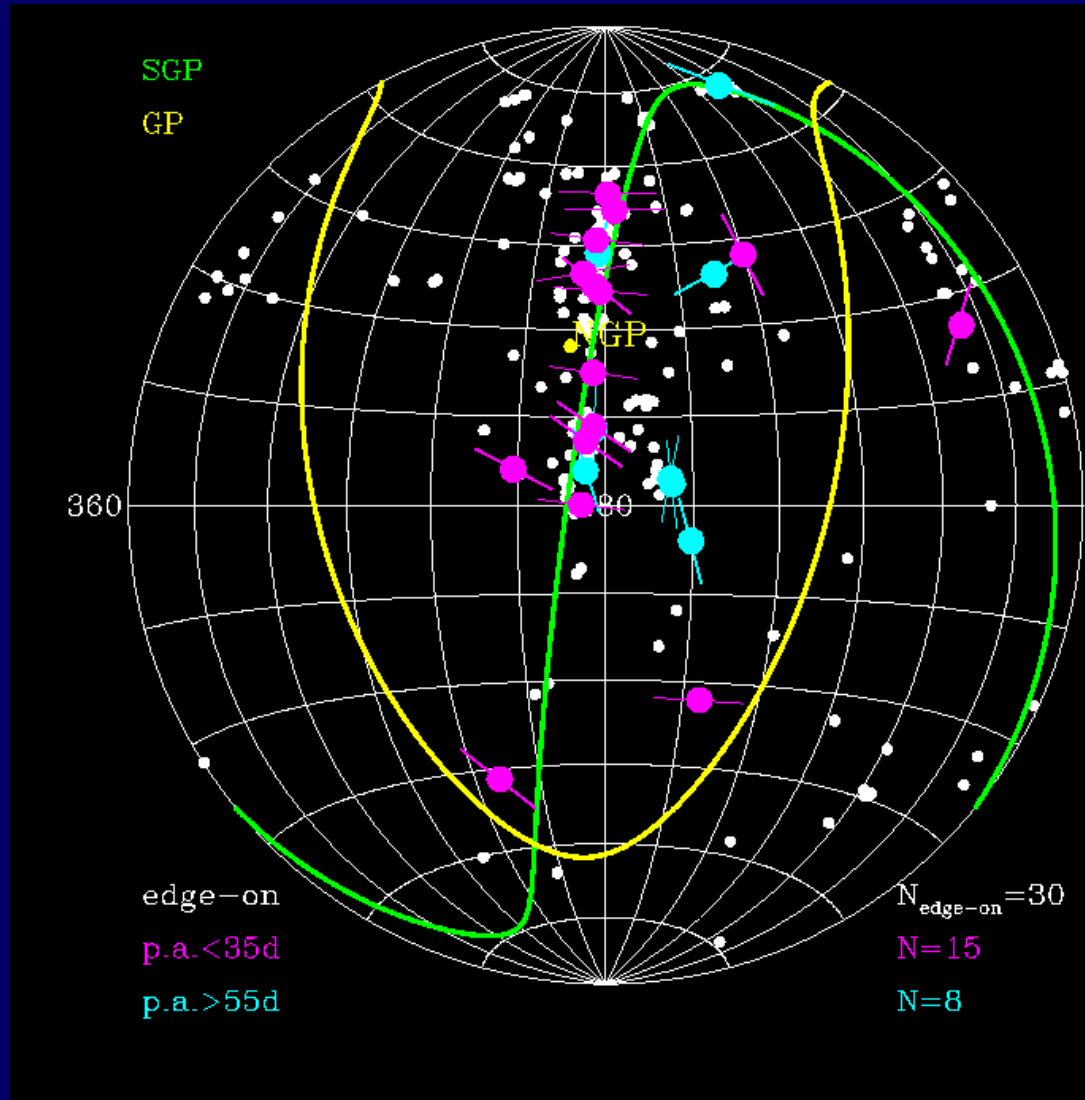
$$J_z = \frac{\partial^2 \phi}{\partial x \partial y} (I_{xx} - I_{yy})$$

$$I_{xx} > I_{yy} > I_{zz}$$

The spin direction is correlated with the intermediate principal axis of the  $I_{ij}$  tensor at turnaround.

In a large-scale pancake: the spin axis should tend to lie in the plane.

# Disk-Pancake Alignment in the Local Supercluster



# Halo Spin Parameter

Fall & Efstathiou 1980

Barnes & Efstathiou 1984

Steinmetz et al. 1994-...

Bullock et al. 2001b

# Halo Spin Parameter

Peebles 76: dimensionless

$$\lambda \equiv \frac{J|E|^{1/2}}{GM^{5/2}}$$

Bullock et al. 2001

$$\lambda \equiv \sqrt{\frac{3}{4}} \frac{J/M}{RV}$$

same for isothermal sphere

$$|E| = \frac{3}{2} M \sigma^2 \quad \sigma^2 = \frac{1}{2} \frac{GM}{R} \quad V^2 = 2\sigma^2$$

TTT:

$$J \sim a^2 \dot{D} \nabla^2 \phi_0 MR_0^3 \sim a^{1/2} M^{5/3}$$

$$a^2 \dot{D} \sim t \sim a^{3/2}$$

J determined at turnaround

$$\delta \sim D \nabla^2 \phi \rightarrow \text{when } \delta \sim 1: \nabla^2 \phi_0 \sim D^{-1} \sim a^{-1}$$

$$\text{comoving } R_0^3 \sim M / \bar{\rho}_0 \sim M$$

$$E \sim M^2 / R \sim a^{-1} M^{5/3}$$

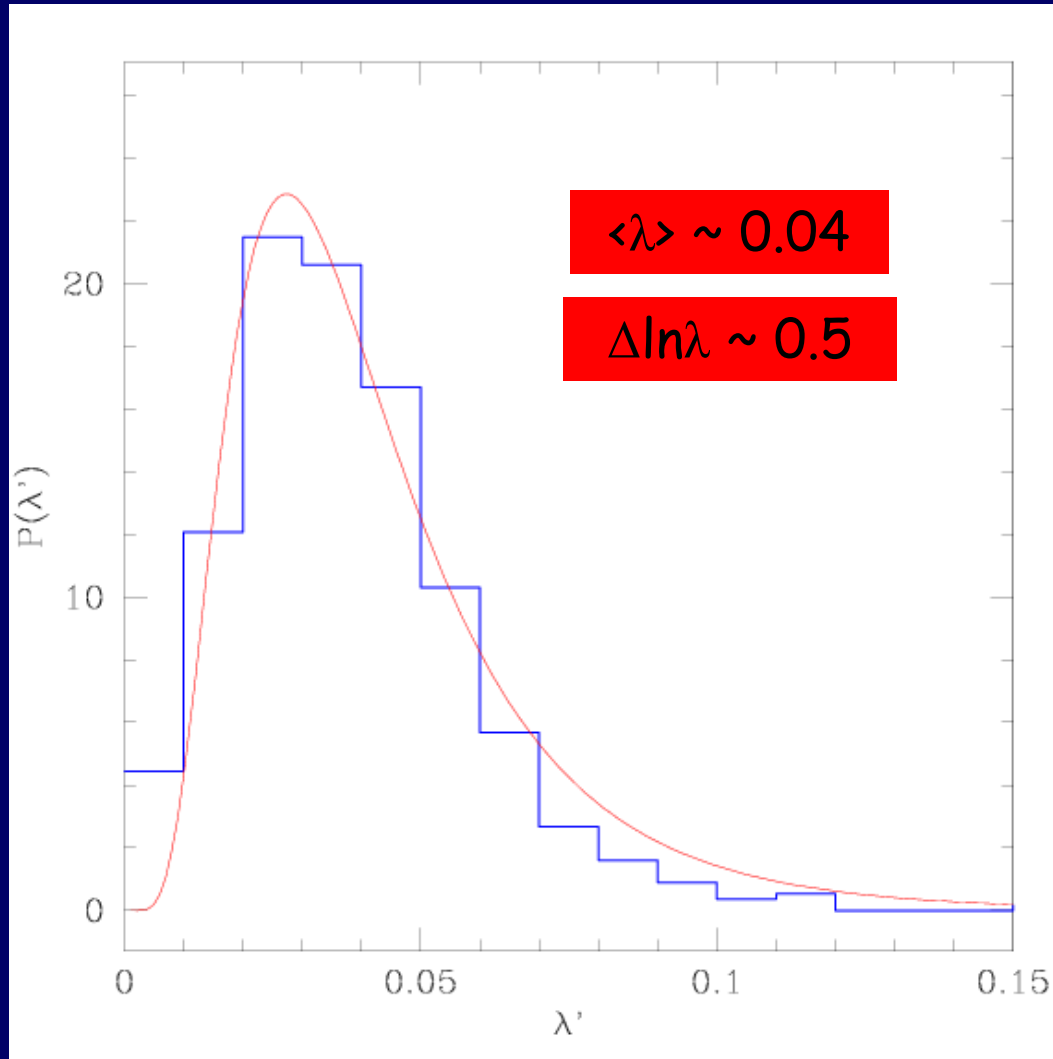
$$\text{Physical } R^3 \sim \rho^{-1} M \sim a^3 M$$



$\lambda$  is constant, independent of  $a$  or  $M$

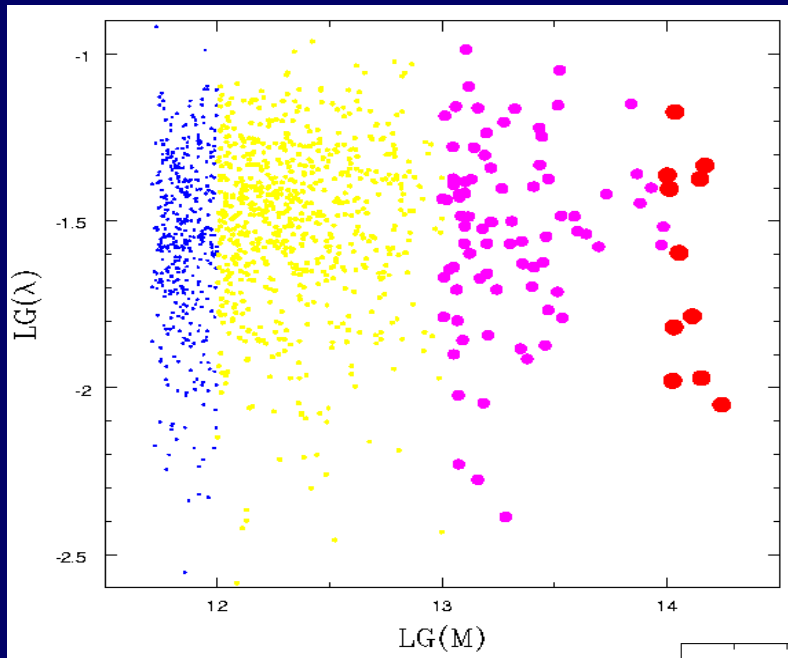
simulations:  $\lambda \sim 0.05$

# Distribution of Halo Spins



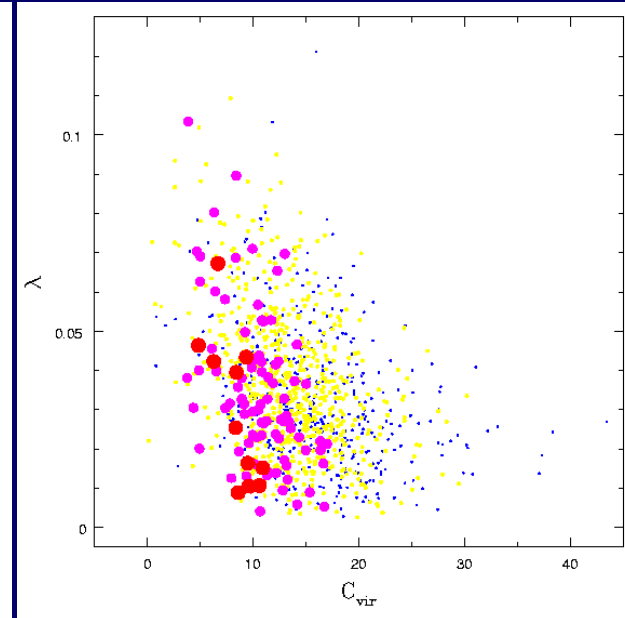
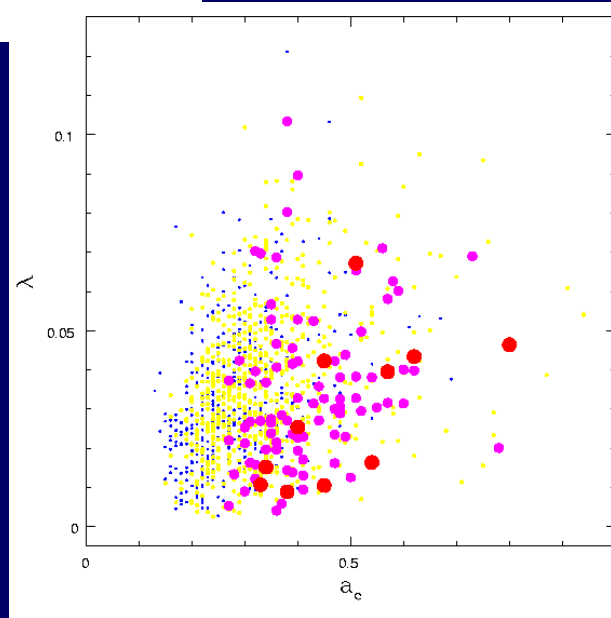


# Spin vs Mass, Concentration, History



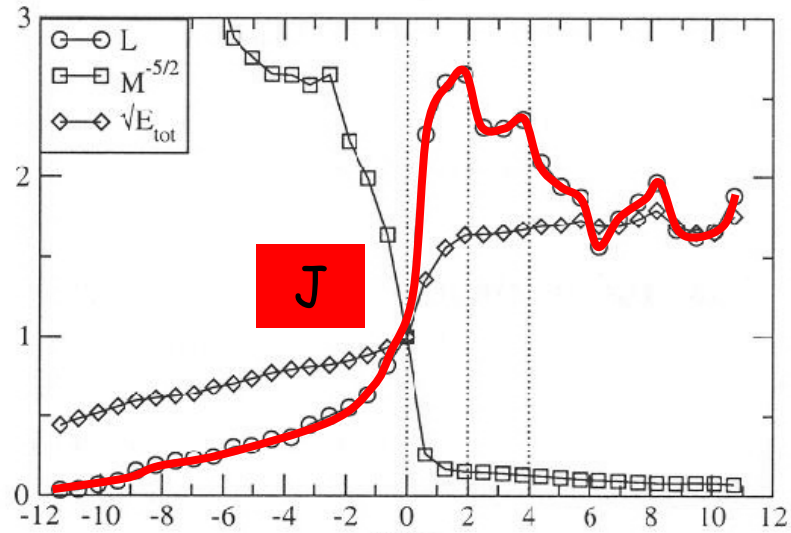
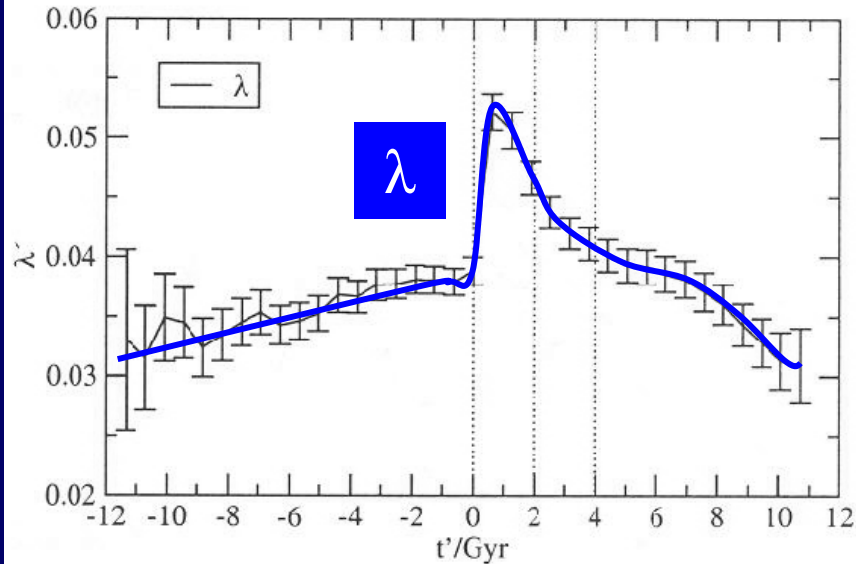
$\lambda$  distribution is universal

$\lambda$  correlated with  $a_c$ ,  
anti-correlated with  $C$



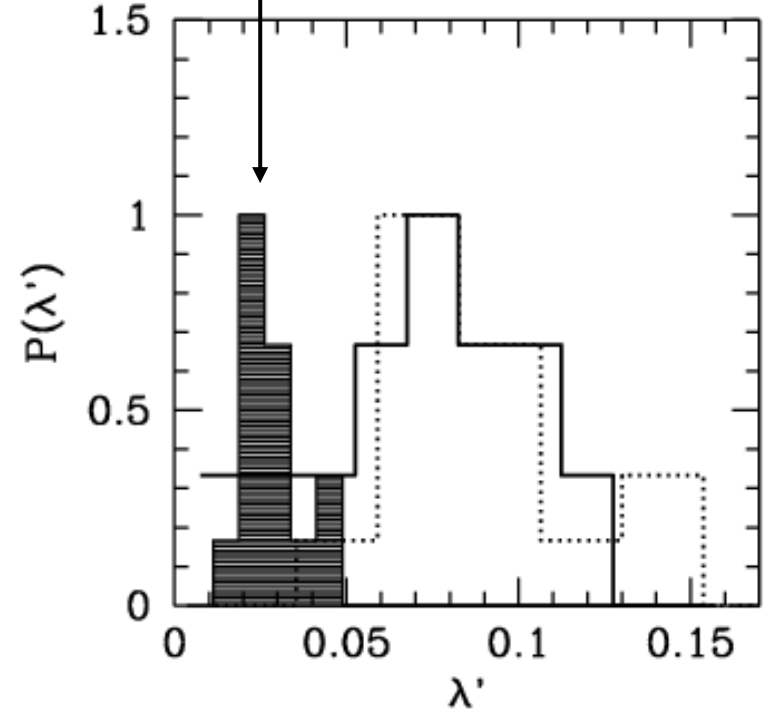
# Spin Jump in a Major Merger

Burkert & D'onghia 04



time

quiet halos with no recent major merger



# J Distribution inside Halos

Bullock et al. 2001b

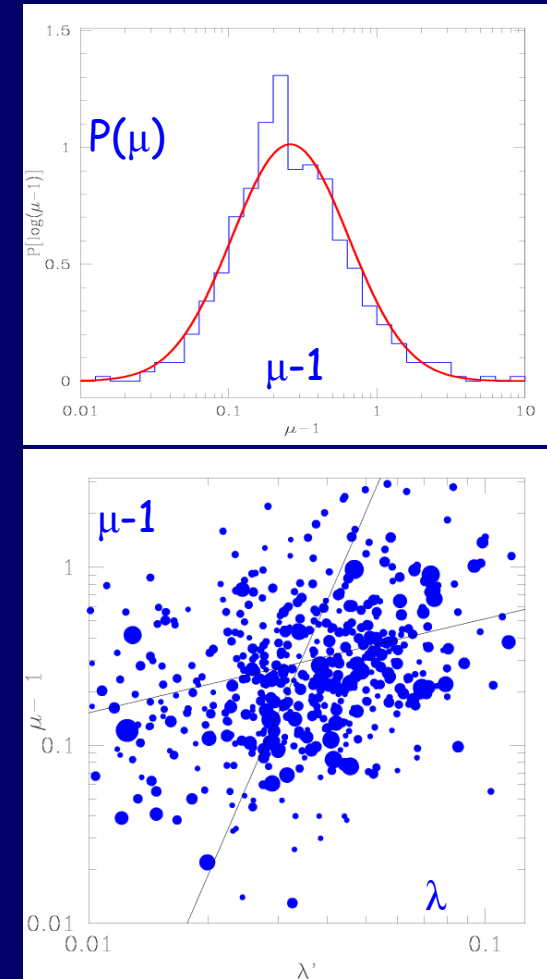
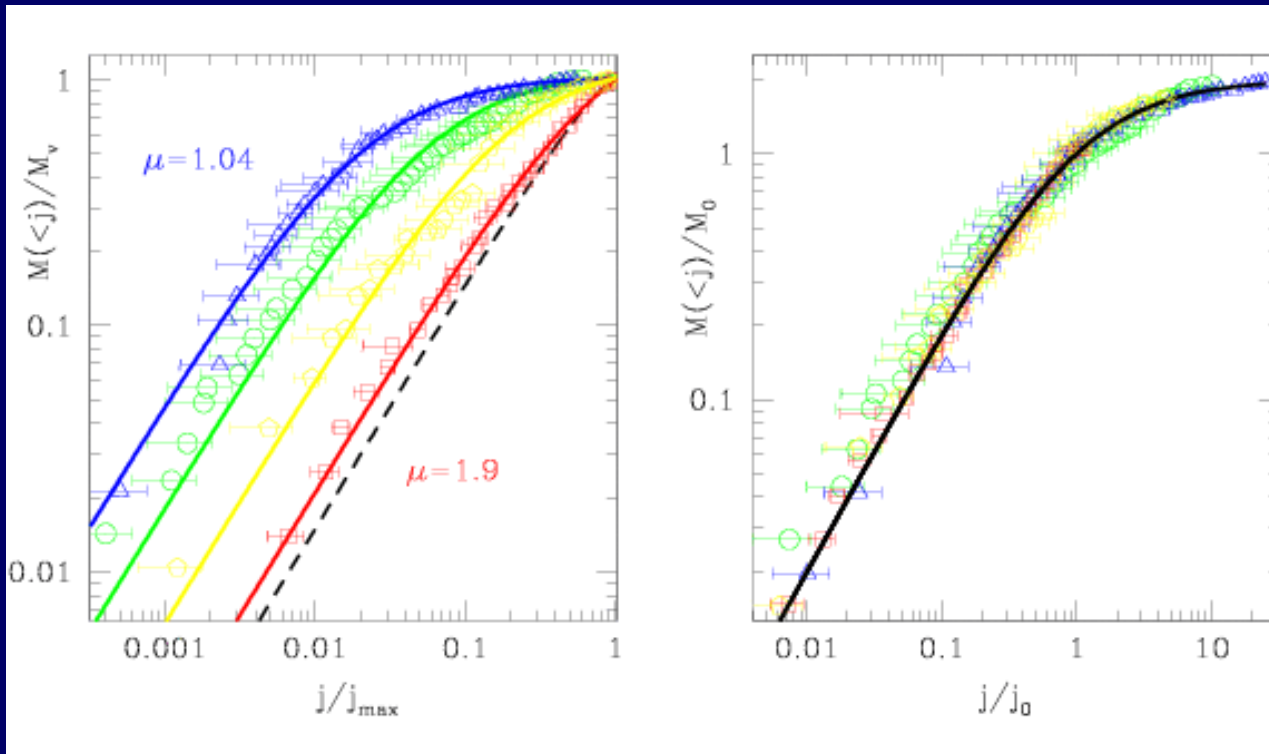
# Universal Distribution of $J$ inside Halos

Bullock et al. 2001b

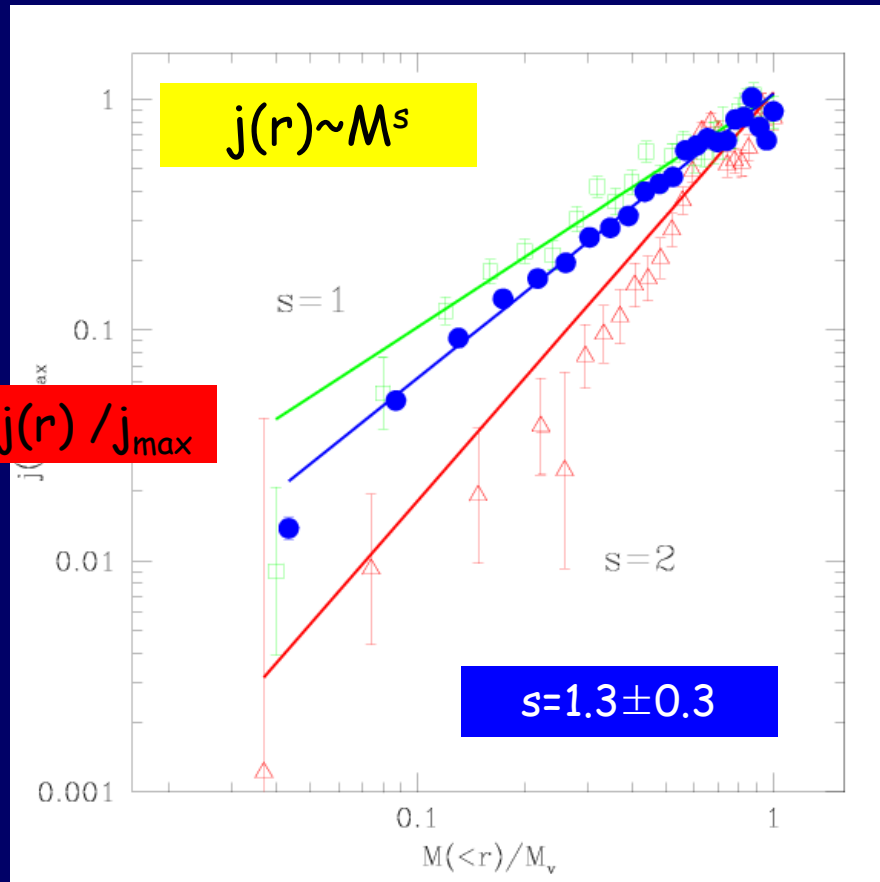
$$M(< j) = M_{vir} \frac{\mu j}{j_0 + j} \quad \mu > 1$$

$$j_{max} = \frac{j_0}{\mu - 1} \quad J/M = j_0 b(\mu) = \sqrt{2} VR \lambda' \quad b(\mu) \equiv -\mu \ln(1 - \mu^{-1}) - 1$$

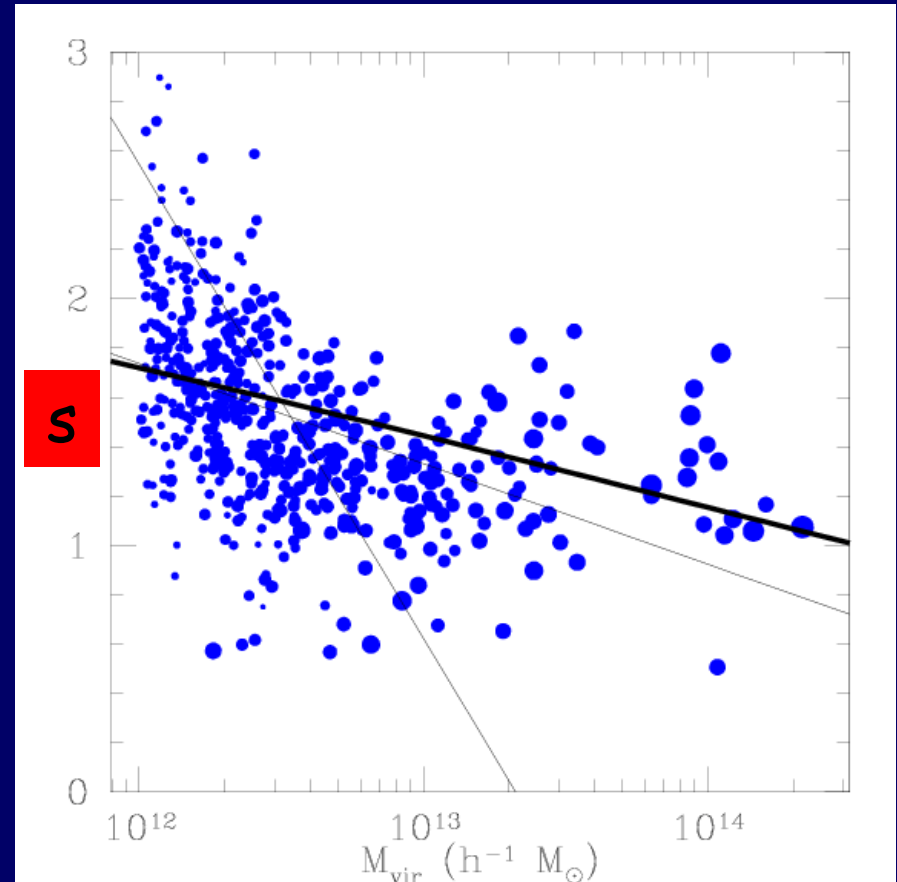
Two parameter family:  
spin parameter  $\lambda$  and shape parameter  $\mu$



# Distribution of $J$ with radius: a power-law profile



$M(<r) / M_v$



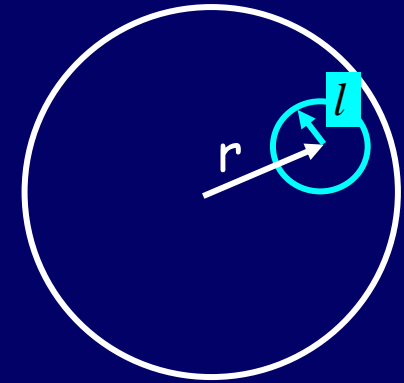
$M_{\text{vir}}$

# Distribution of J in space

Toy model: J by minor mergers

Tidal radius  $\frac{m(l_t)}{l_t^2} = \frac{l_t}{r^3} \left( 2M(r) - r \frac{dM}{dr} \right)$

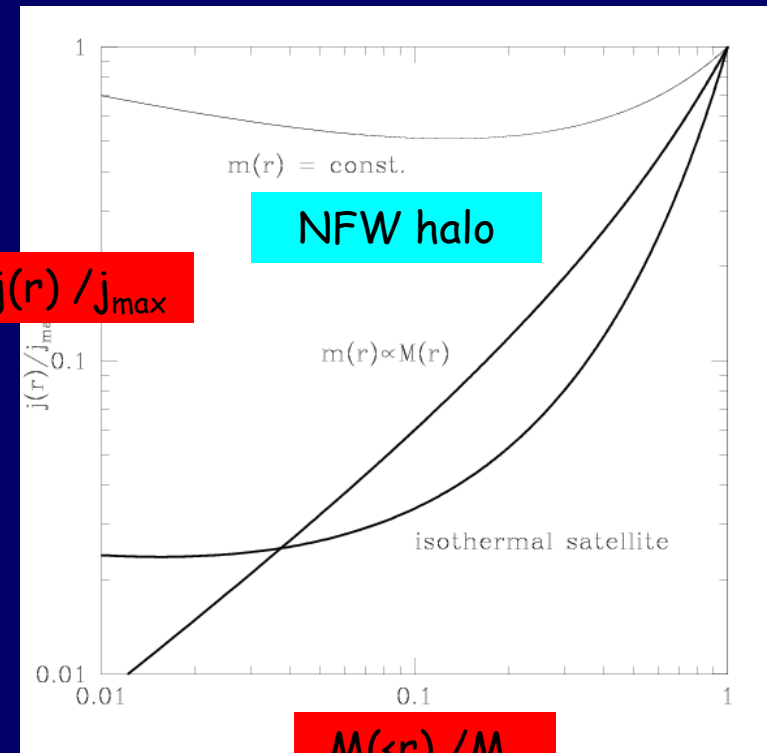
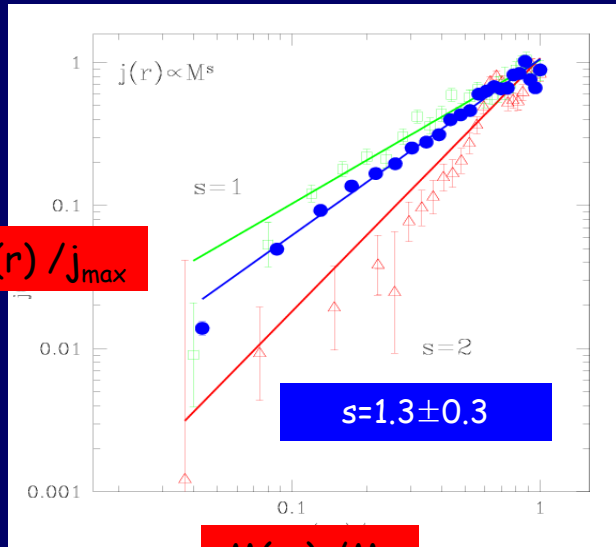
$M \propto r^\alpha \rightarrow m[l_t(r)] \propto M(r)$



Assume  $m$  and  $j$  are deposited locally in a shell  $r$

$4\pi r^2 \rho(r) j(r) = m(r) \frac{d[rV(r)]}{dr} + \frac{dm}{dr} rV(r)$

$M \propto r, \quad m \propto l \rightarrow j(M) \propto M \propto r$



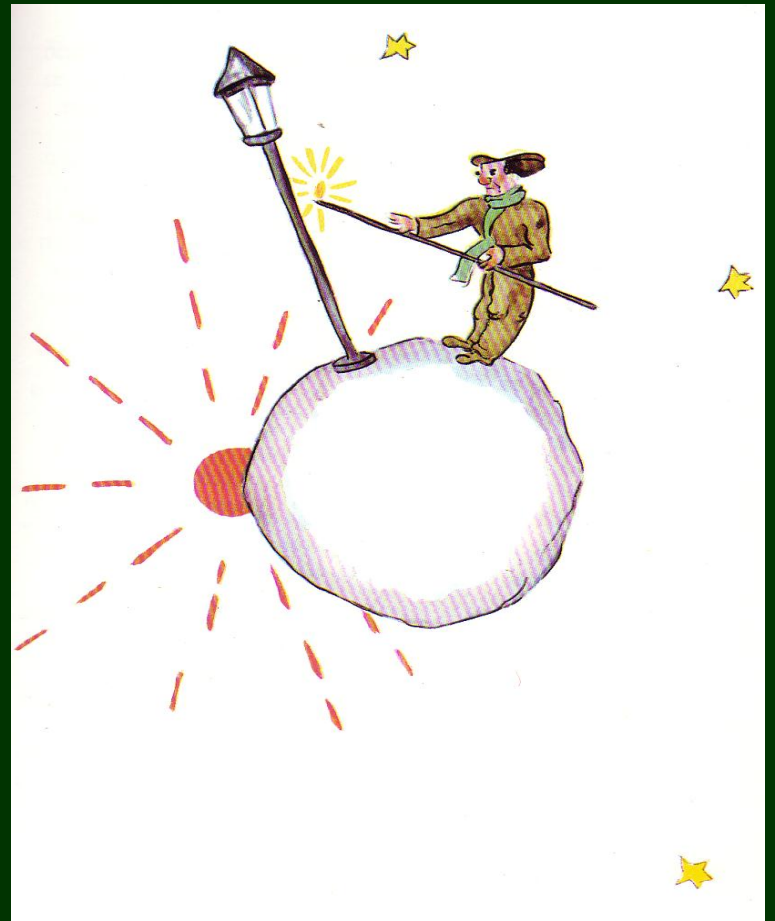


# Formation of Stellar Disks and Spheroids inside DM Halos

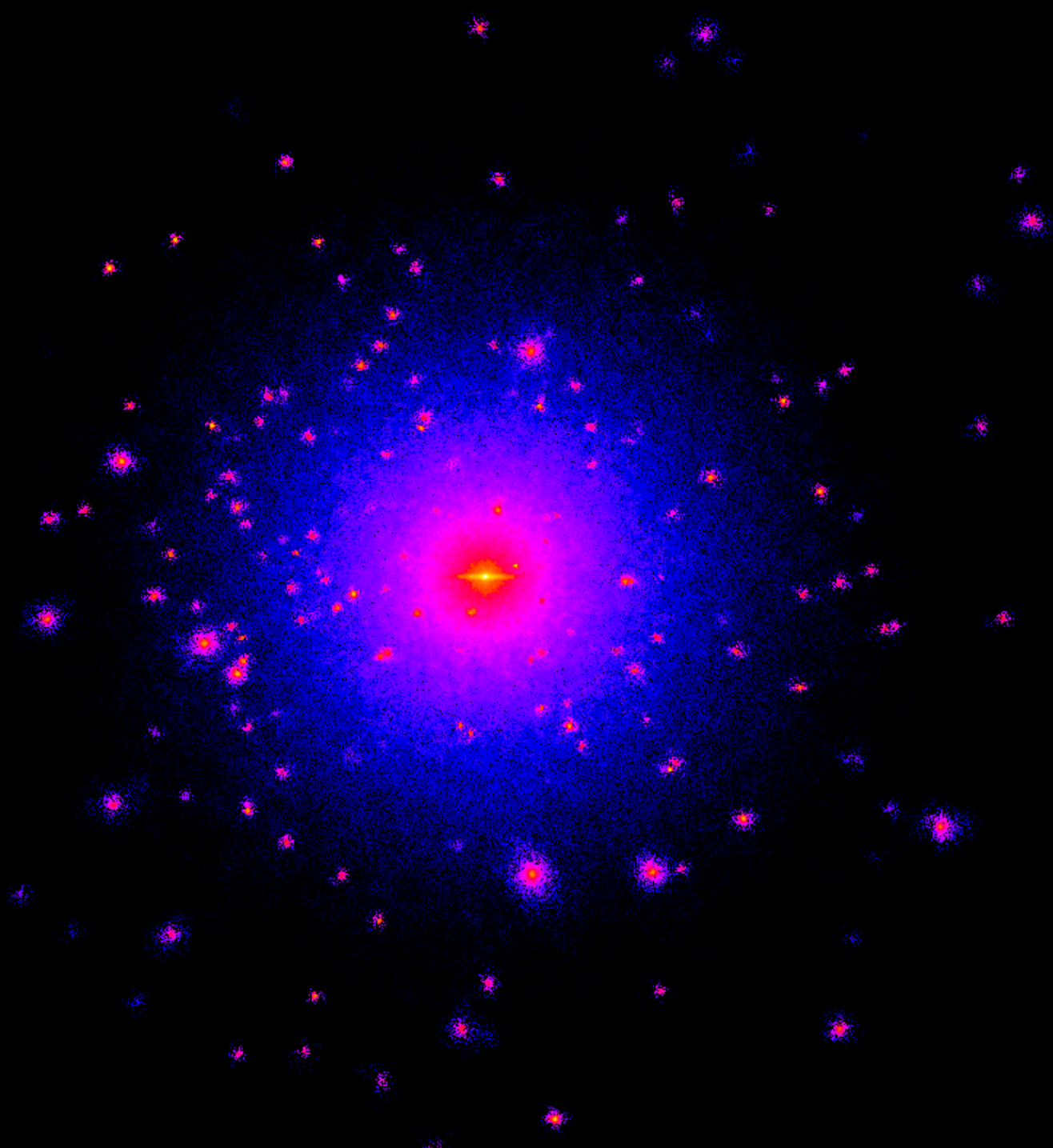
White & Rees 1978

Fall & Efstathiou 1980

Mo, Mao & White

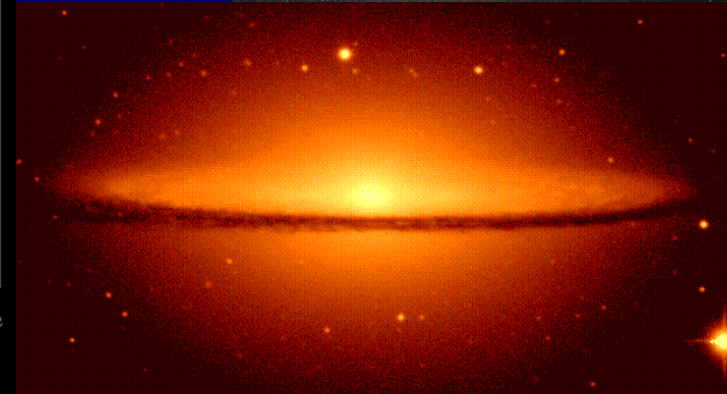
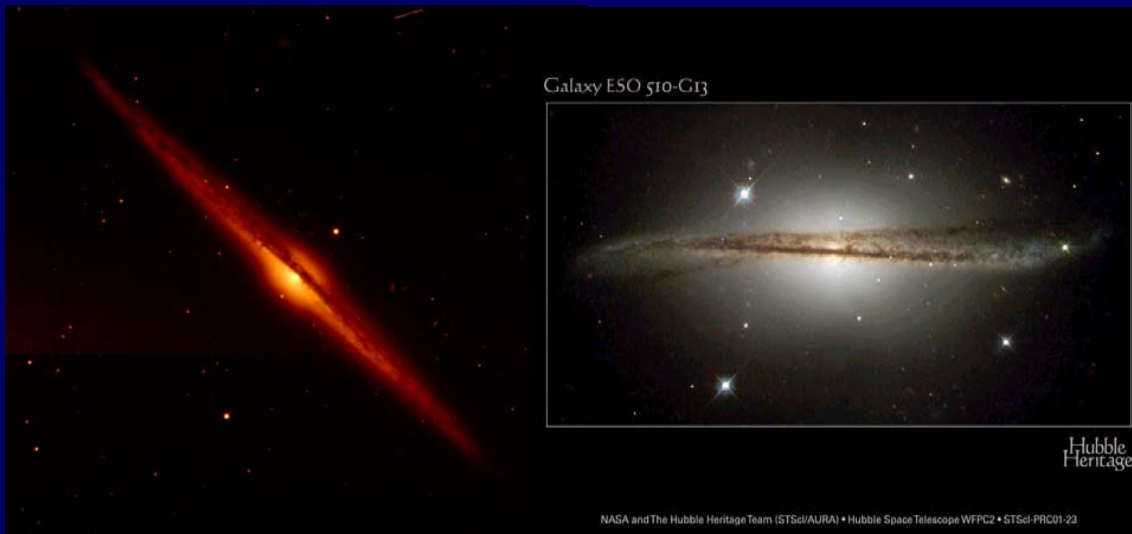


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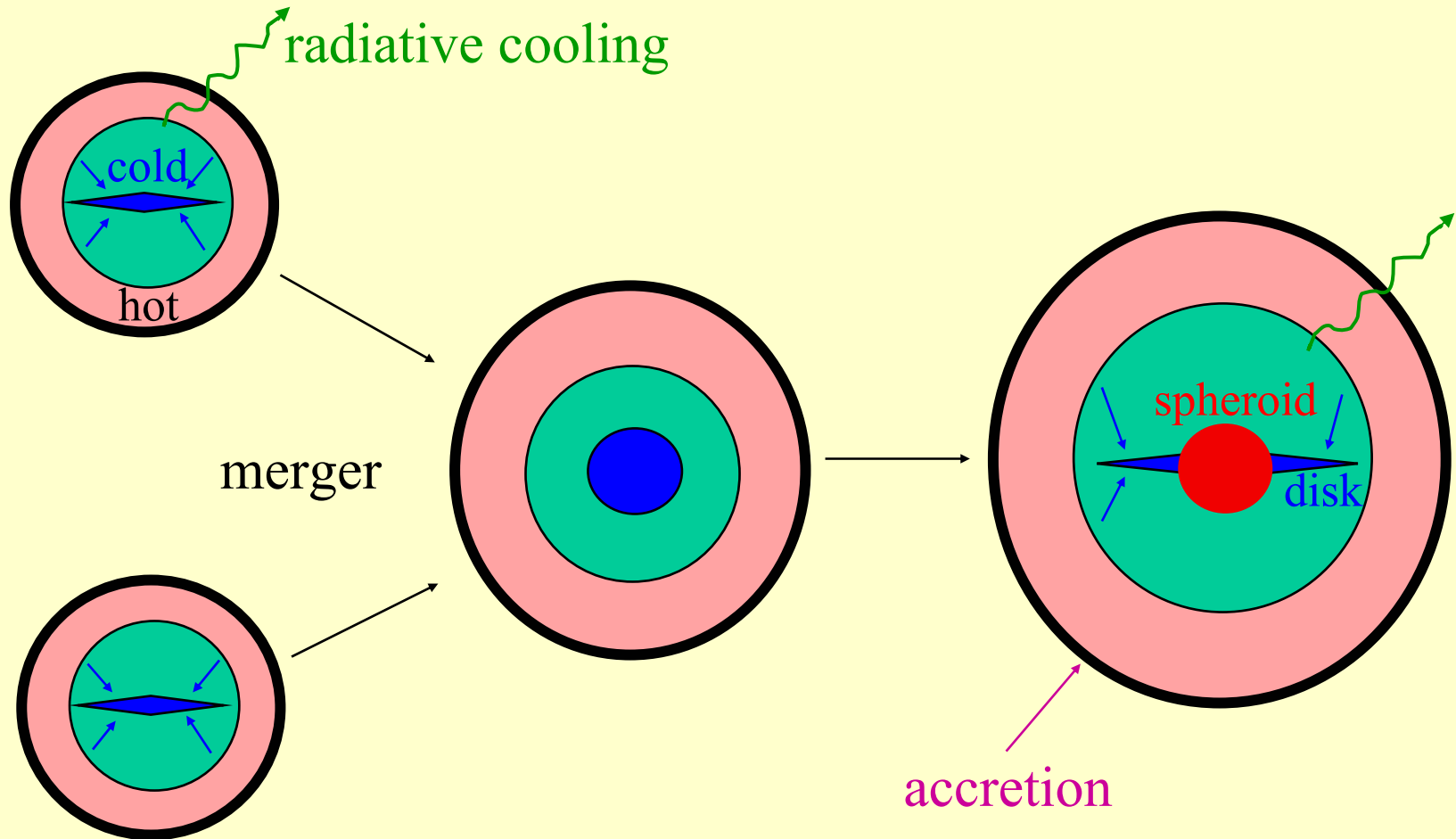


# Galaxy Types: Disks and Spheroids

- The morphology of a galaxy is a transient feature dictated by the mass accretion history of its dark matter halo
  - most stars form in disks; spheroids result from subsequent mergers
  - disks result from smooth gas accretion; oldest disk stars are often used to date the last major merger event



# Galaxy Formation in halos



halos    cold gas → young stars → old stars

# Disk Size

Spin parameter

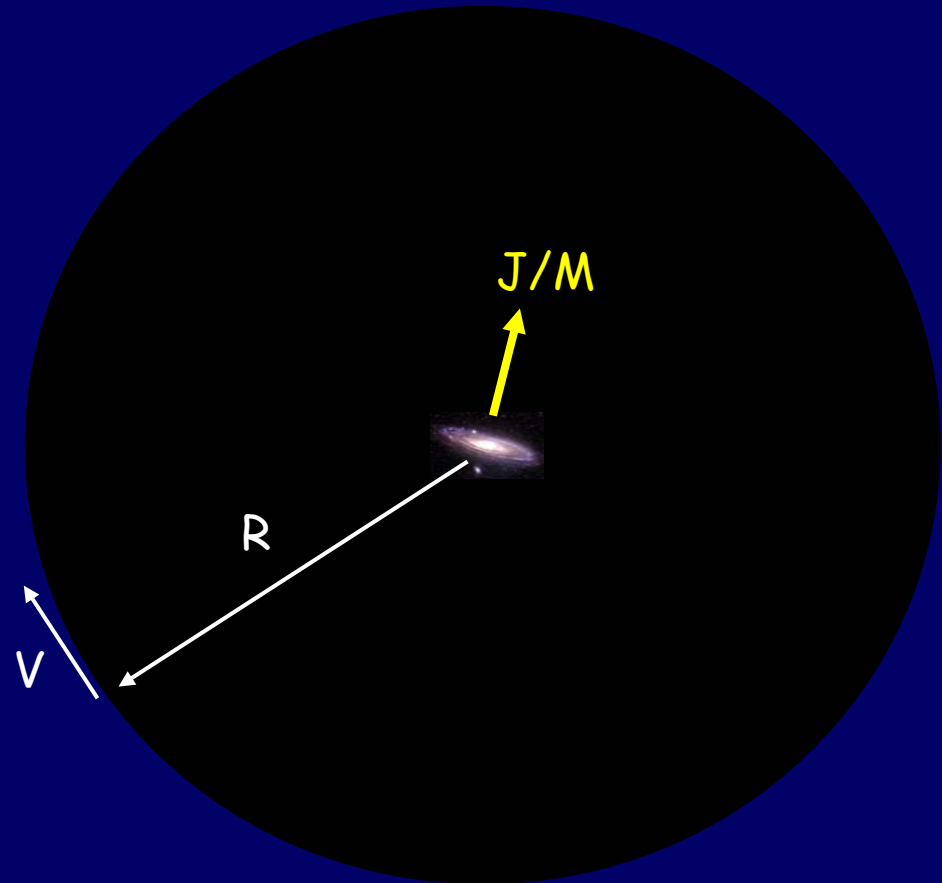
$$\lambda \sim \frac{J/M}{RV}$$

Conservation of specific angular momentum

$$\text{const.} = J/M \sim \lambda R_{\text{virial}} V \sim R_{\text{disk}} V$$



$$\frac{R_{\text{disk}}}{R_{\text{virial}}} \sim \lambda$$



# Disk Profile from the Halo J Distribution

Assume the gas follows the halo  $j$  distribution

$$M_{gas}(< j) = f M(< j)$$

Assume conservation of  $j$  during infall from halo to disk.

In disk:

$$j(r) = Vr = [GM(r)r]^{1/2}$$

In disk: lower  $j$  at lower  $r$

$$M_{halo}(< j) \rightarrow m_{disk}(r)$$

$$M_{halo}(< j) = M_{vir} \frac{\mu j}{j_0 + j} \quad \mu > 1$$



$$m_d(r) = f\mu M_v \frac{j(r)}{j_0 + j(r)} \quad j(r) < j_{max}$$

Assume isothermal sphere  
No adiabatic contraction

$$M \propto r \rightarrow j(r) = rV(r) = rV_{vir}$$



$$m_d(r) = f\mu M_v \frac{r}{r_d + r} \quad r < r_{max}$$

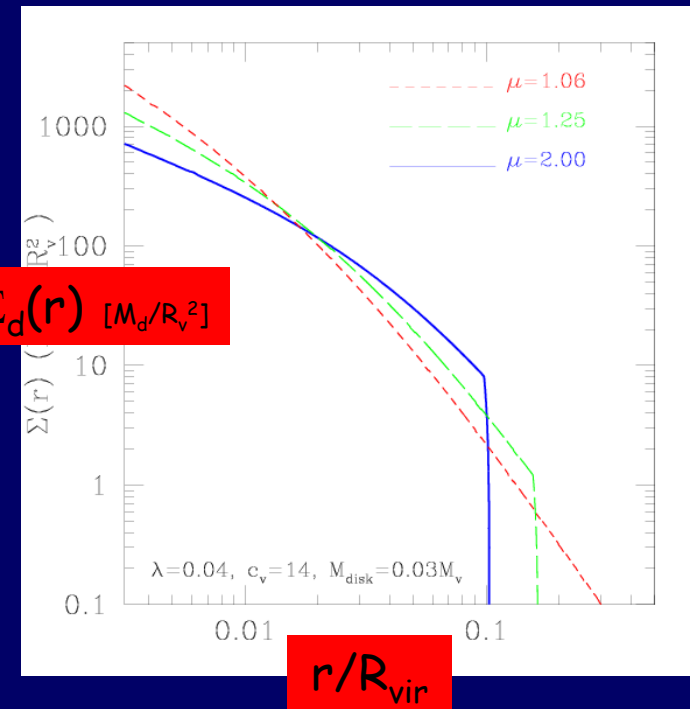
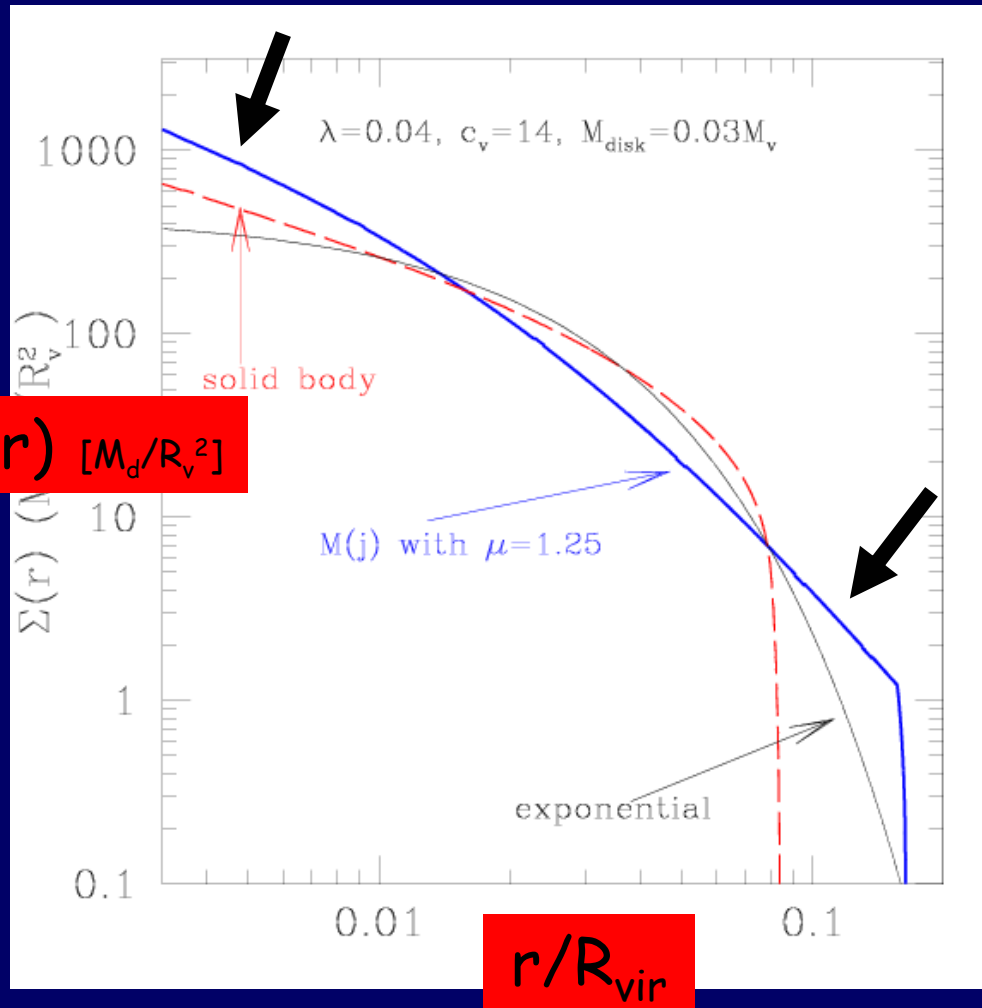
$$r_d = \sqrt{2\lambda'} R_v b^{-1}(\mu)$$

$$r_{max} = r_d / (\mu - 1)$$

$$\Sigma_d(r) = \frac{f\mu M_v}{2\pi} \frac{r_d}{r(r_d + r)^2}$$

# Disk Profile: Shape Problem

Bullock et al. 2001b



# The Angular-Momentum Problem

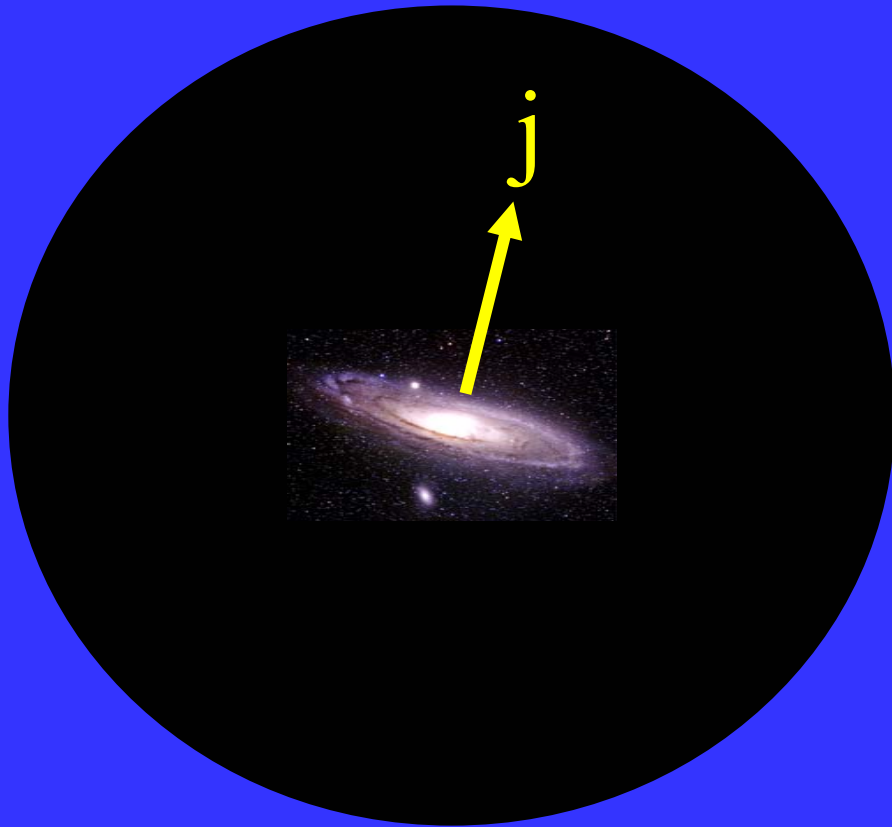
Navarro & Steinmetz



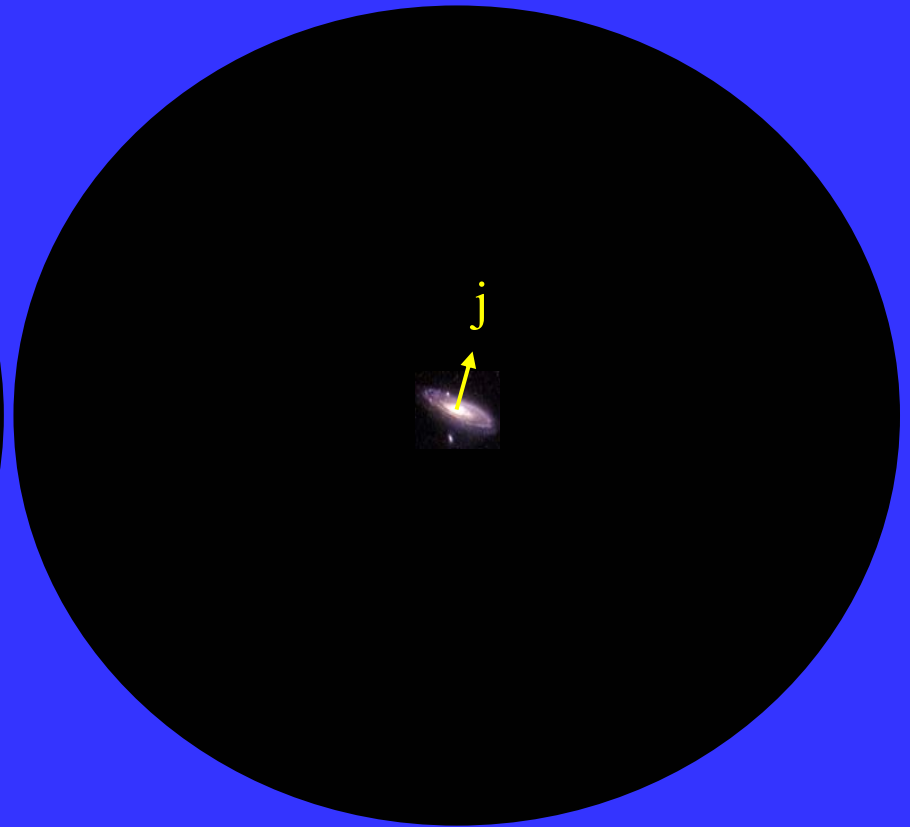
# The Spin Catastrophe

Navarro & Steinmetz et al.

observations

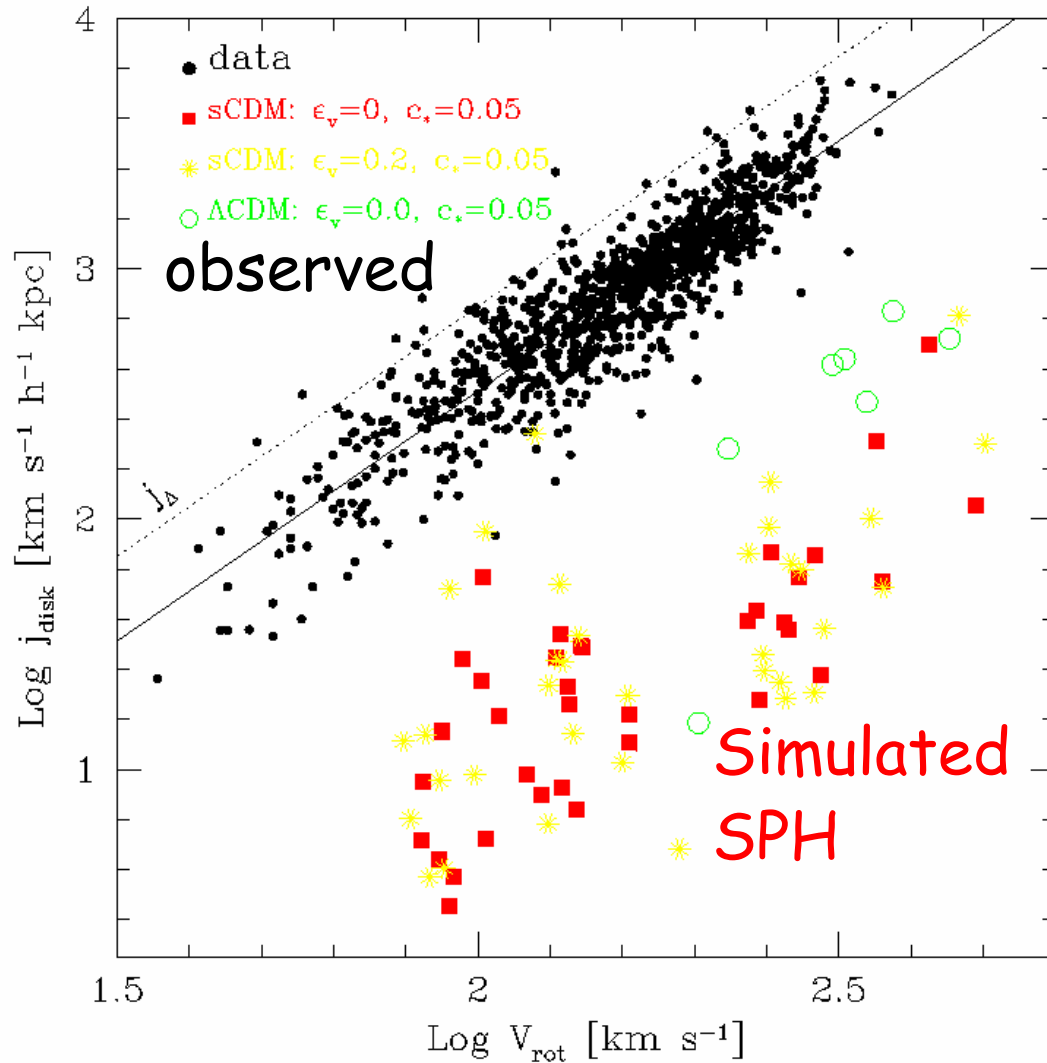


simulations

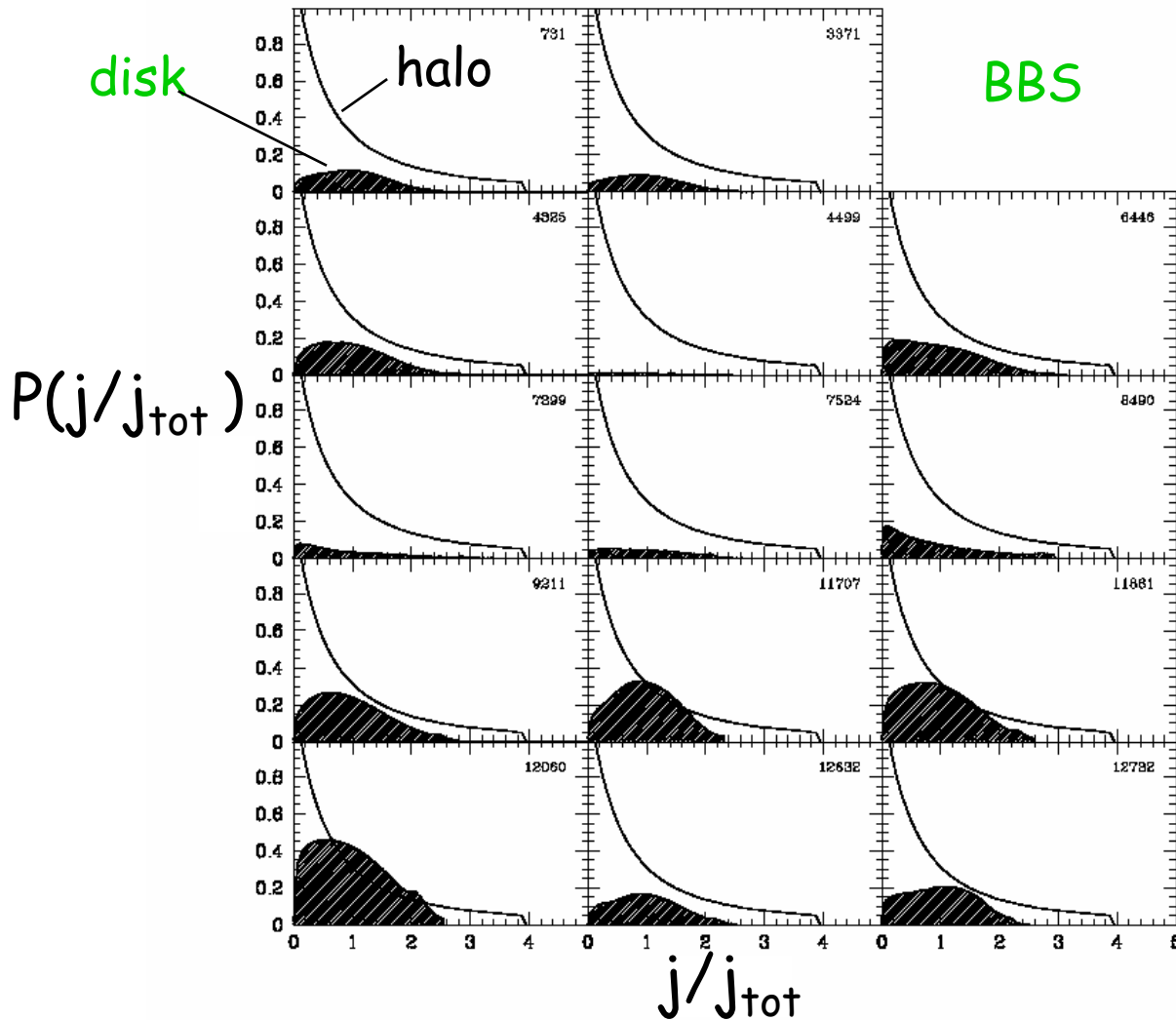


# The spin catastrophe

$j_{\text{disk}}$



# Observed $j$ distribution in dwarfs



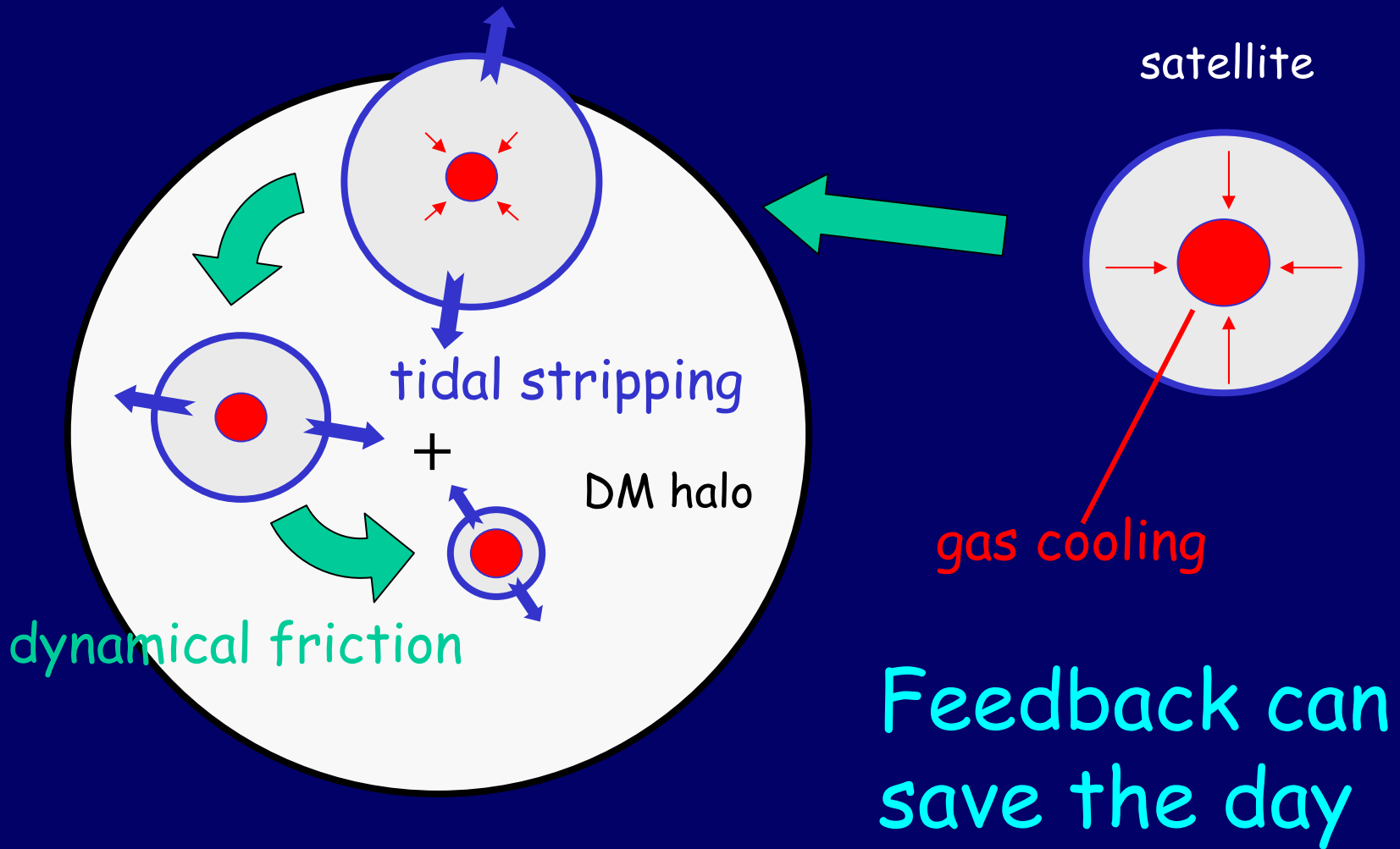
Low  $f_{\text{baryons}} \approx 0.03$

Missing low  $j$

High  $\lambda_{\text{baryons}} \approx 0.07$

# Over-colling → spin catastrophe

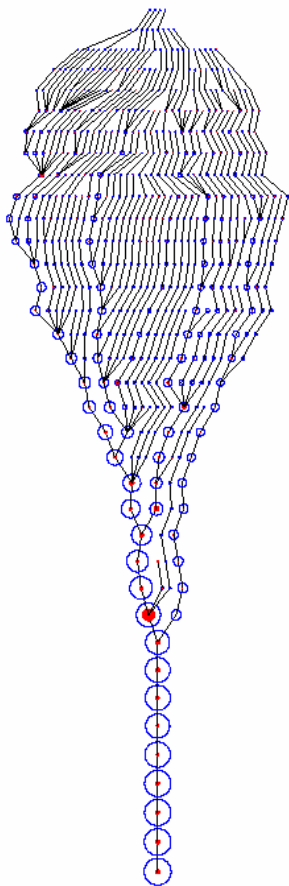
Maller & Dekel 02



# Orbital-merger model:

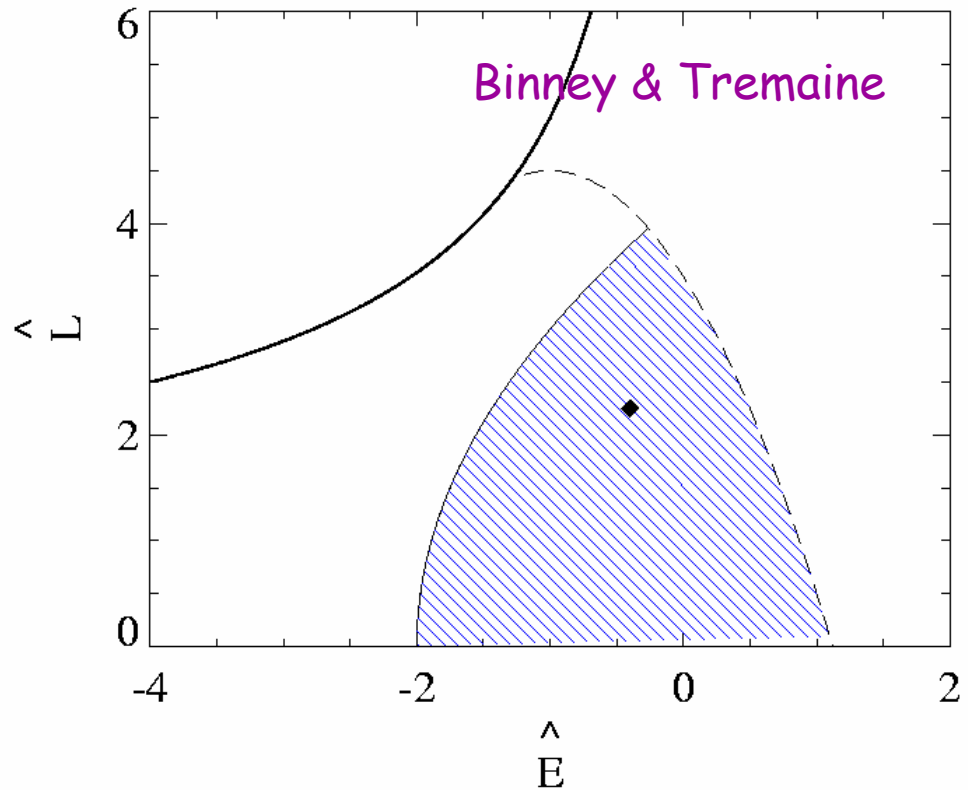
Add orbital angular momentum in merger history

Merger history



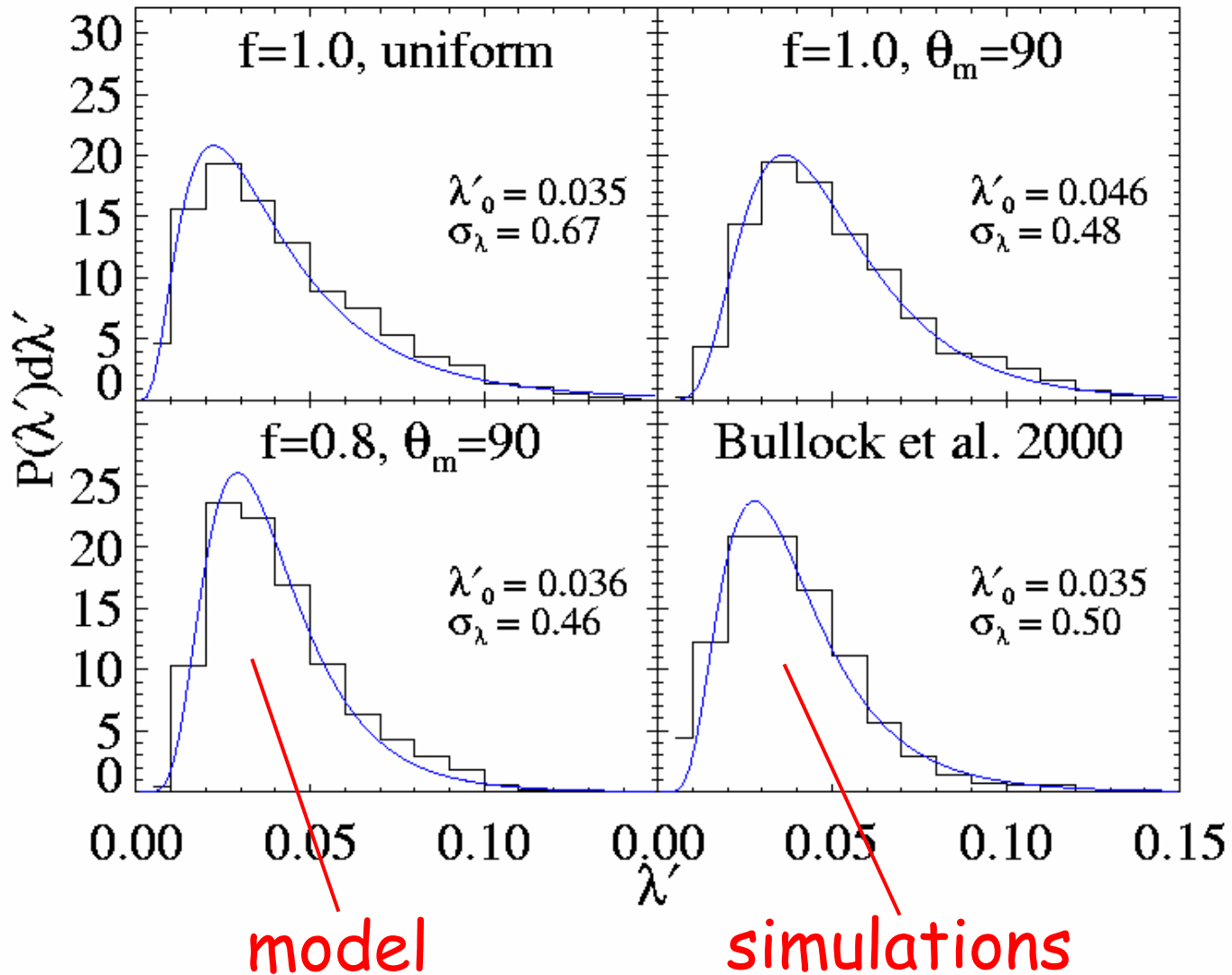
0.122  
0.14  
0.169  
0.182  
0.2  
0.253  
0.287  
0.302  
0.335  
0.377  
0.403  
0.425  
0.455  
0.485  
0.5  
0.529  
0.557  
0.59  
0.628  
0.65  
0.668  
0.71  
0.74  
0.772  
0.8  
0.835  
0.871  
0.893  
0.911  
0.926  
0.941  
0.95  
0.973  
0.982  
0.991  
1.000

Orbit parameters

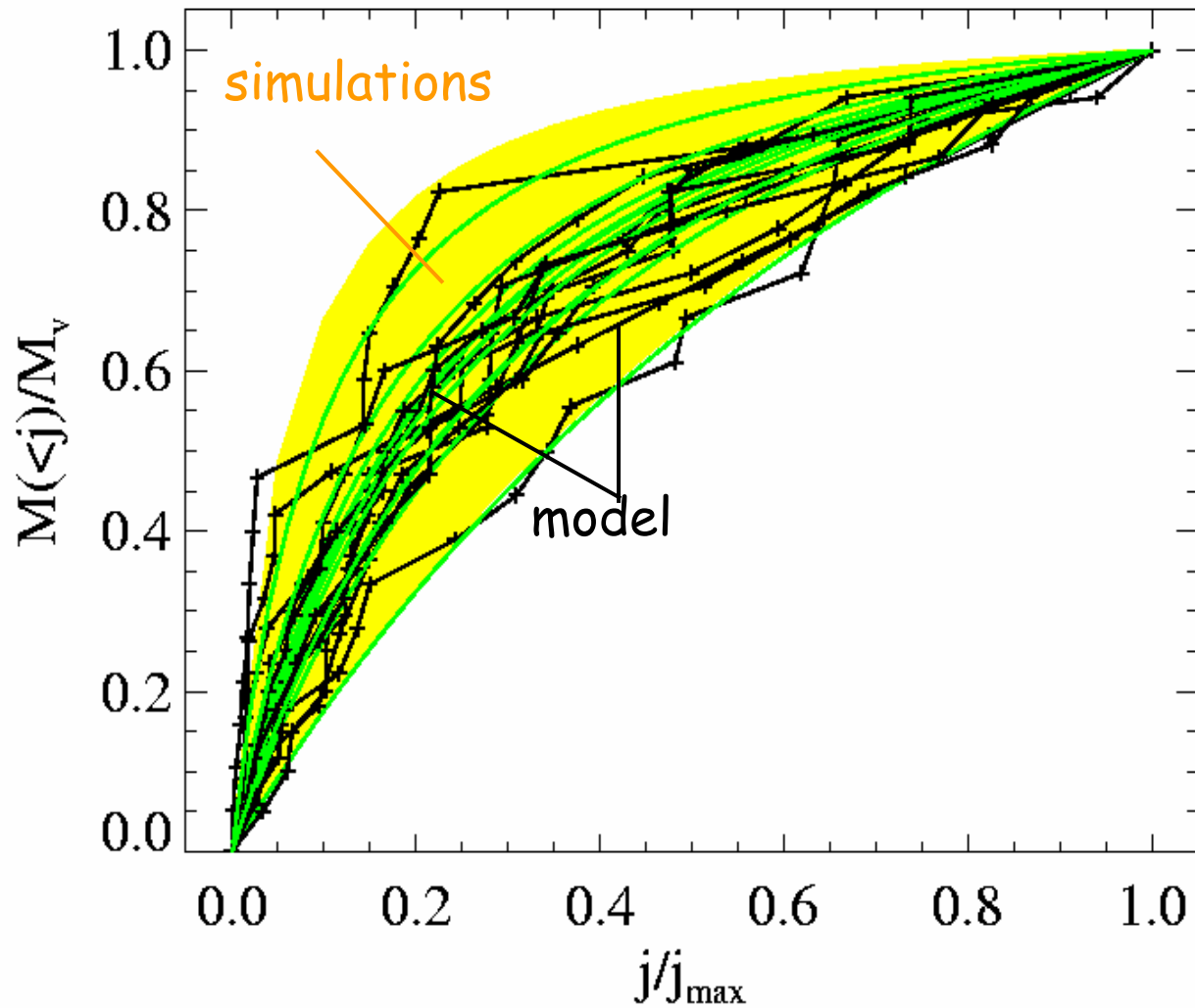


and random orientation

# Success of orbital-merger model

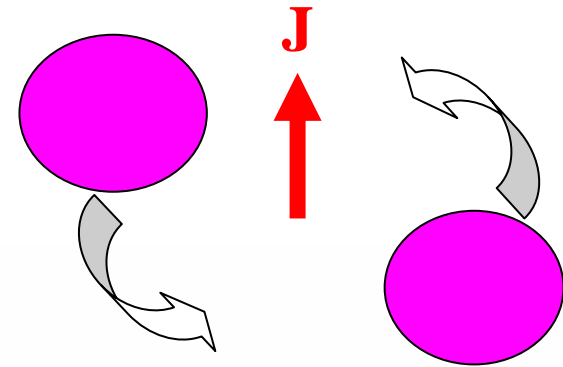
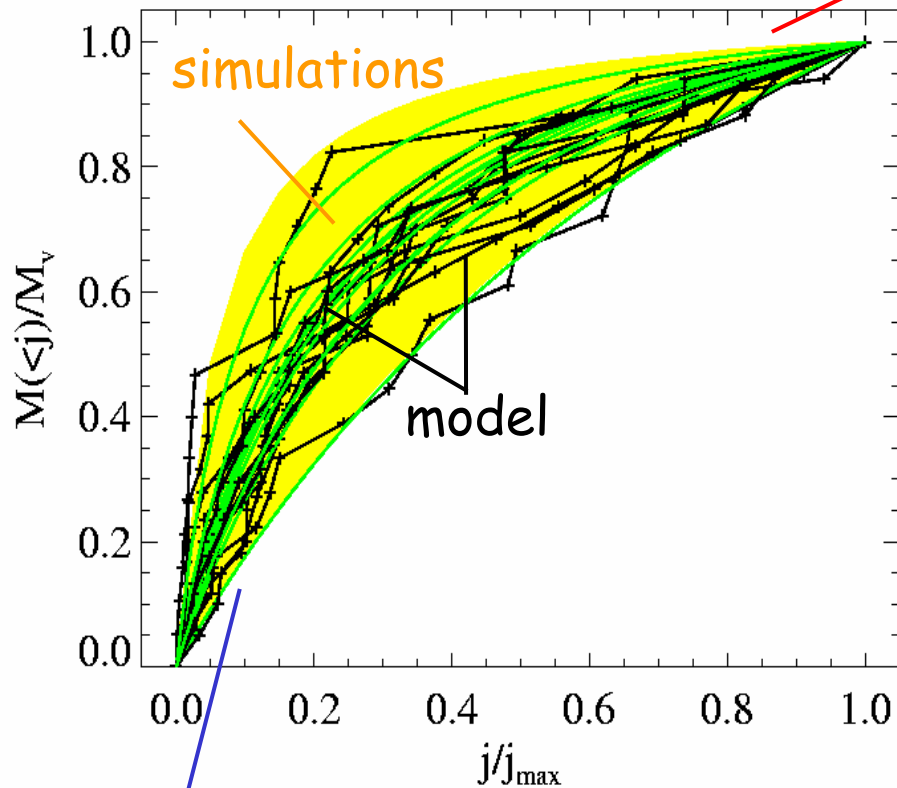


# Model success: $j$ distribution in halos

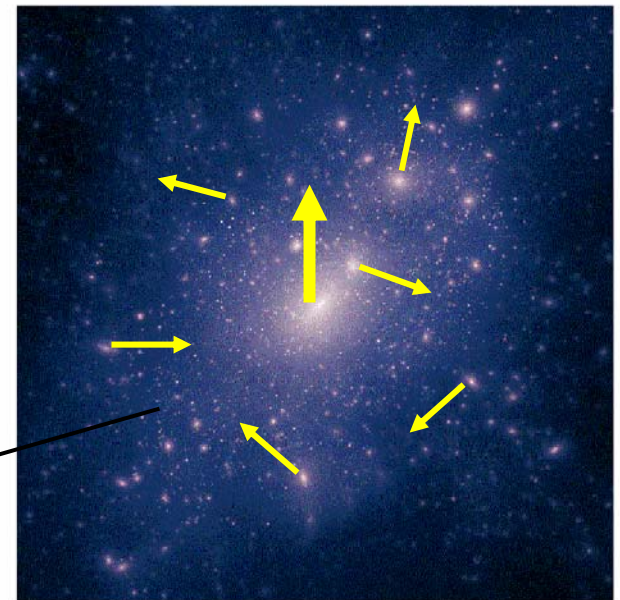


# Low/high- $j$ from minor/major mergers

High- $j$  from major mergers



Low- $j$  from minor mergers





# Supernova Feedback: $V_{\text{SN}}$ (Dekel & Silk 86; Dekel & Woo 03)

Energy fed to the ISM during the “adiabatic” phase:

$$E_{\text{SN}} \approx \nu \varepsilon \dot{M}_* t_{\text{rad}} \propto M_* (t_{\text{rad}} / t_{\text{ff}})$$

$$\dot{M}_* \approx M_* / t_{\text{ff}}$$

$$\approx 0.01$$

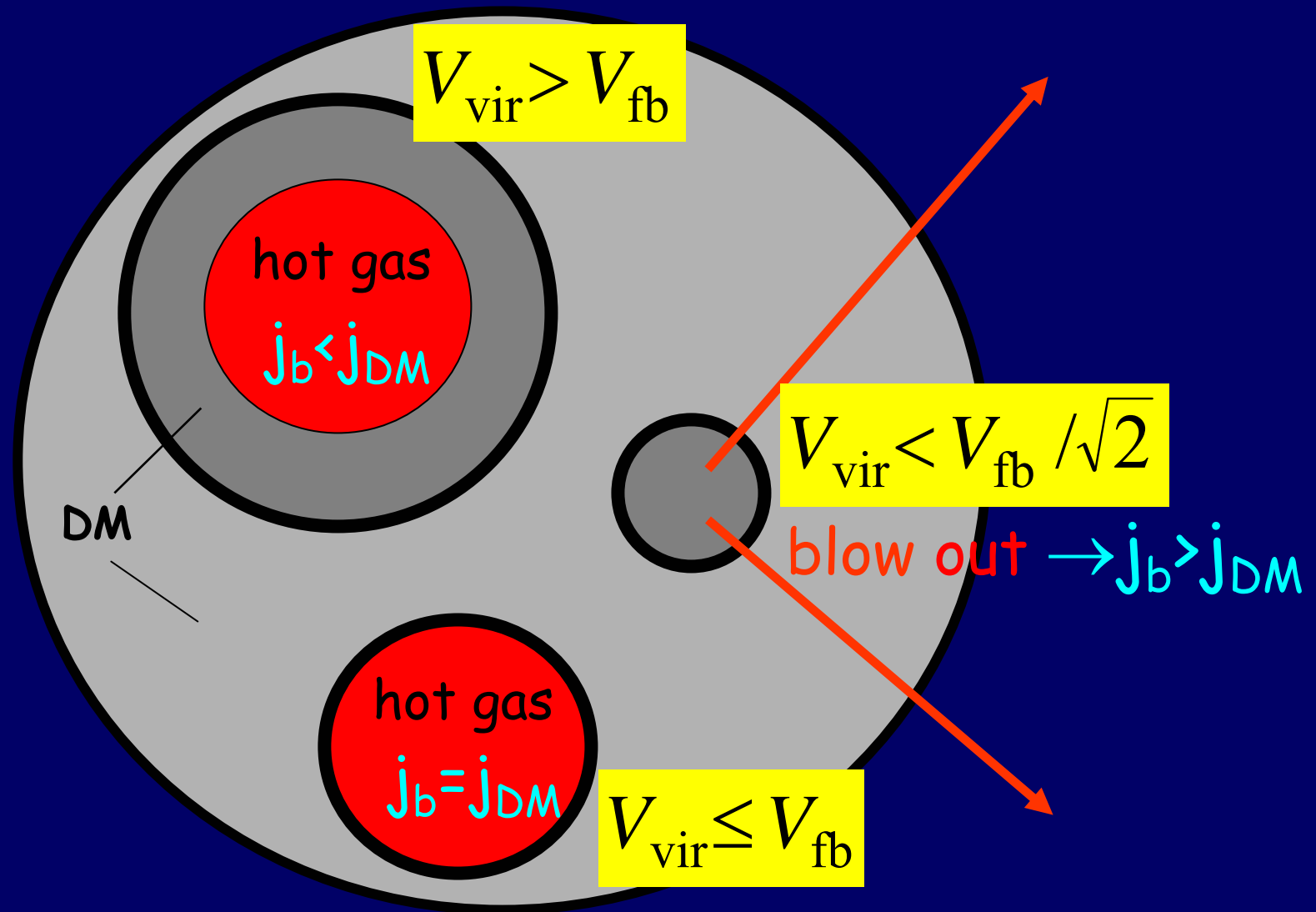
for  $\Lambda \propto T^{-1}$  at  $T \sim 10^5 \text{ K}$

Energy required for blowout:

$$E_{\text{SN}} \approx M_{\text{gas}} V^2$$

$$\rightarrow V_{\text{crit}} \approx 100 \text{ km/s} \rightarrow M_{*\text{crit}} \approx 3 \times 10^{10} M_{\odot}$$

# Feedback in satellite halos

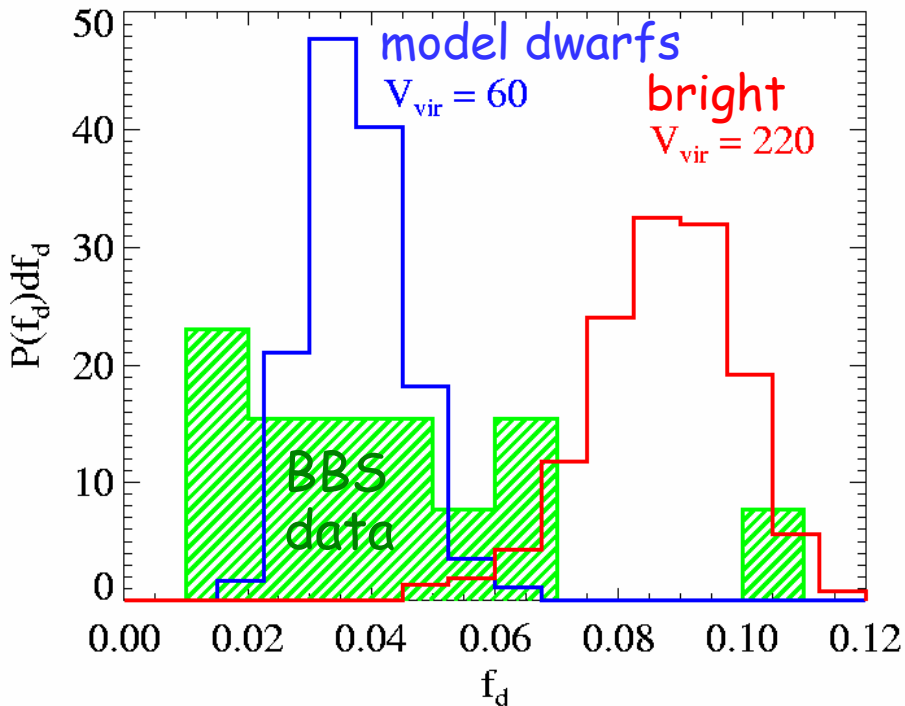


# Model vs Data

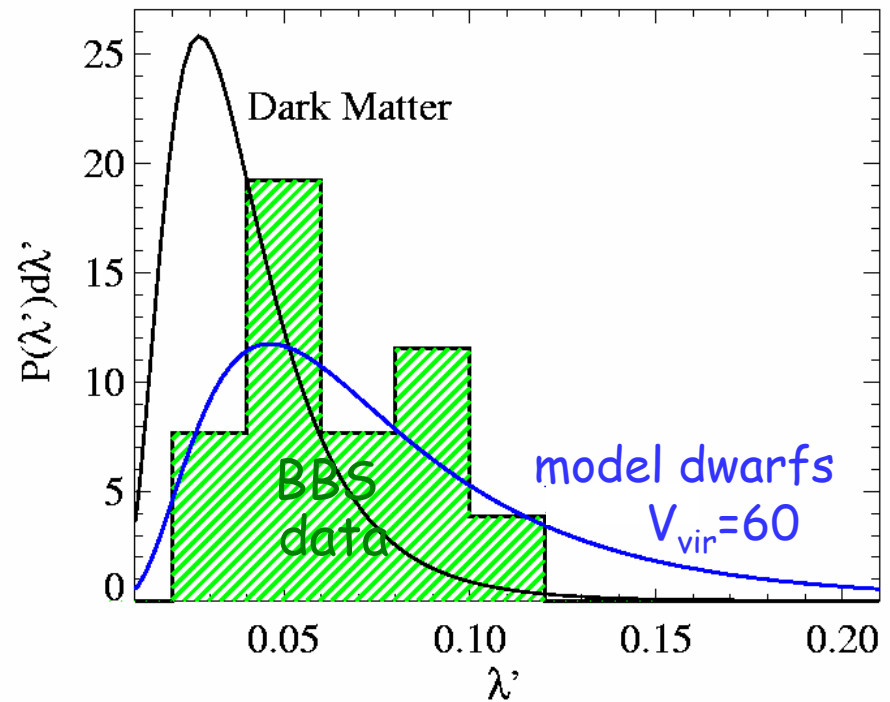
(Maller & Dekel 02)

BBS data: 14 dwarfs, van den Bosch, Burkert & Swaters 02

## baryon fraction

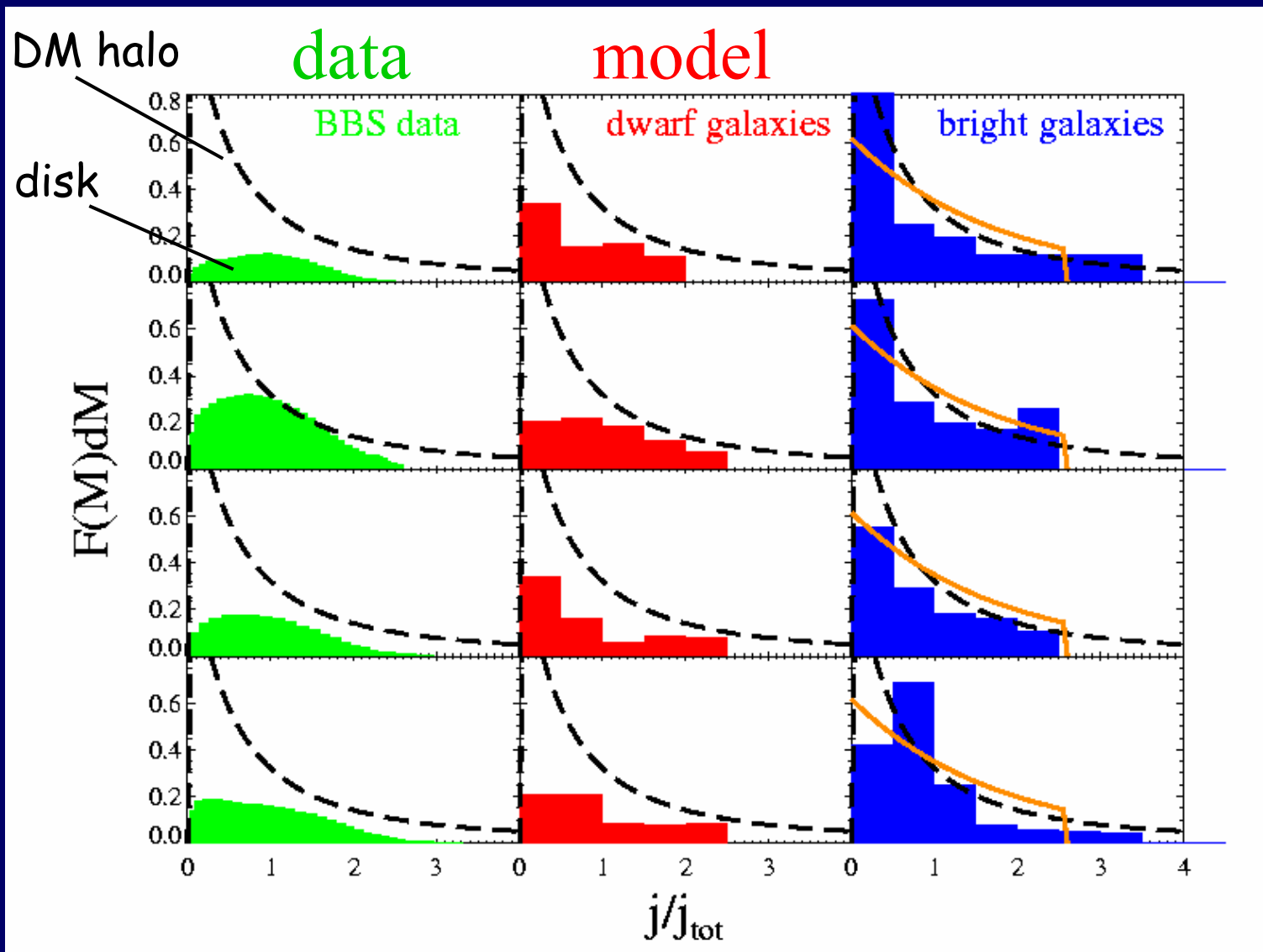


## spin parameter



One free parameter in model:  $V_{\text{feedback}} \cong 90 \text{ km s}^{-1}$

# J-distribution within galaxies



# Summary: feedback effect on spin

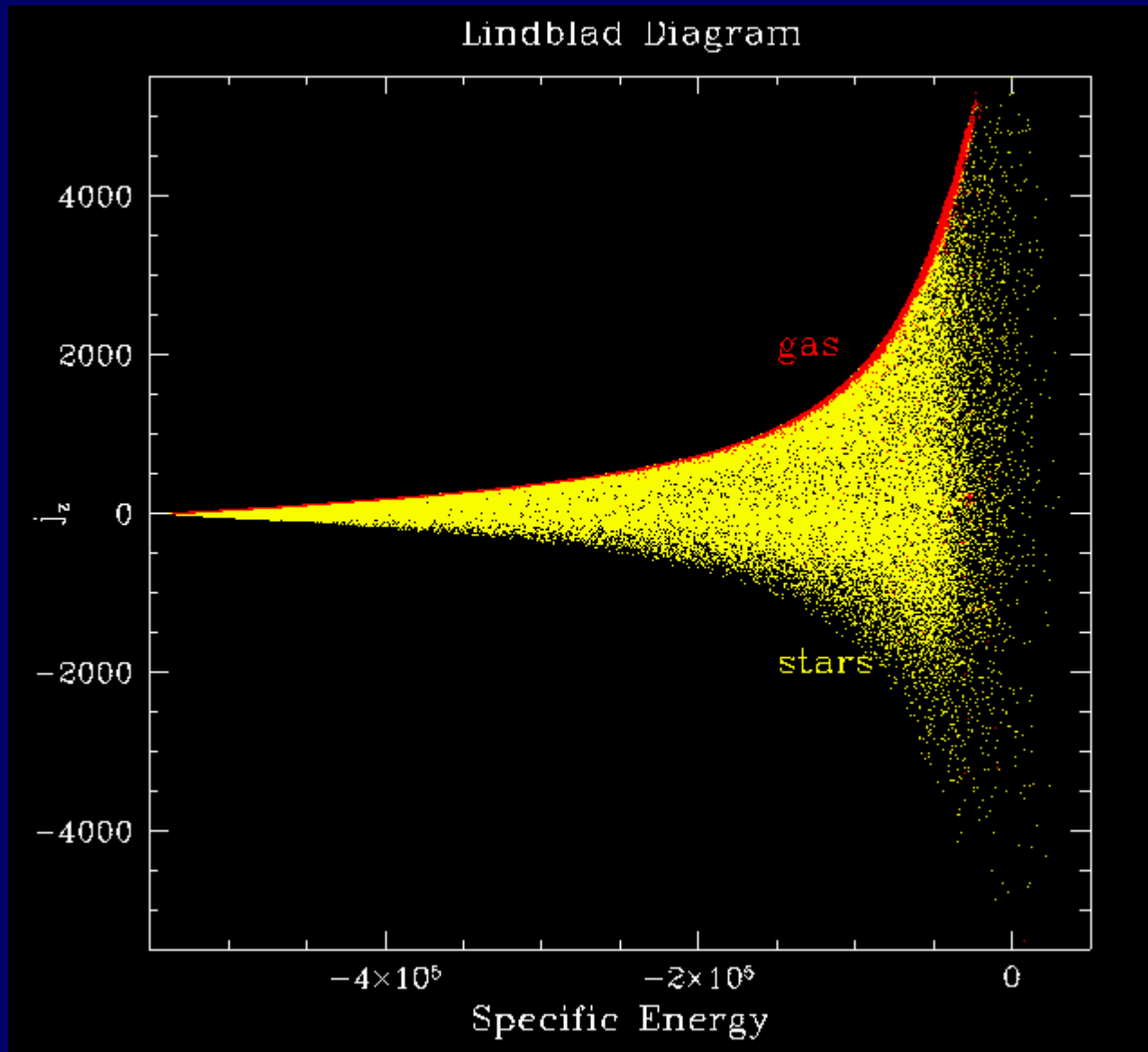
In big satellites (merging to big galaxies)  
heating  $\rightarrow$  gas expansion  $R_b \sim R_{DM}$   
 $\rightarrow$  tidal stripping together  $\rightarrow \lambda_{bar} \sim \lambda_{DM}$

In small satellites (merging to dwarfs)  
gas blowout  $\rightarrow f_{bar}$  down  
blowout of low  $j$  gas  $\rightarrow \lambda_{bar} > \lambda_{DM}$

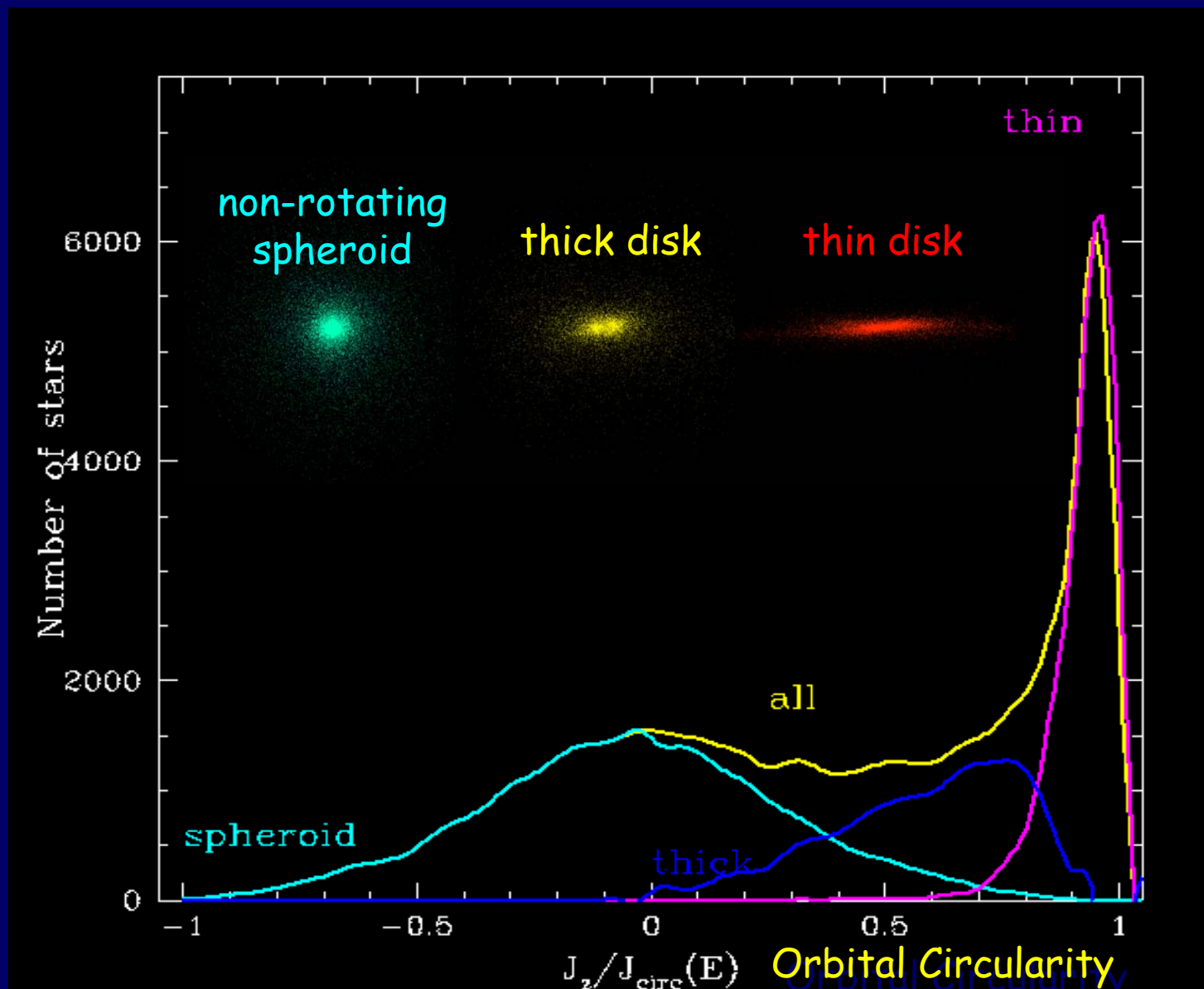
# Thin Disk and Thick Disk

Navarro & Steinmetz

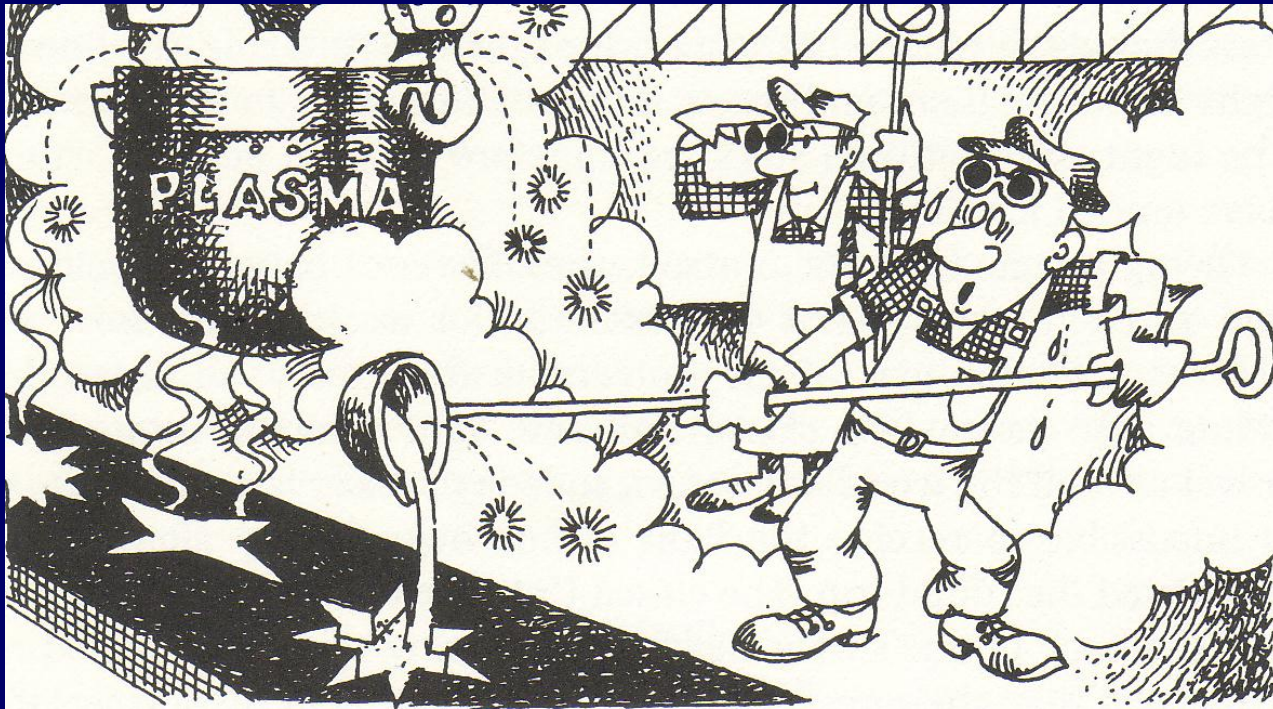
# Dynamical Components of a Simulated galaxy



# Dynamical components of a simulated galaxy







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***“Great PowerPoint, Kevin, but the answer is no.”***