

Joint second- and third-order shear statistics and cosmological constraints with CFHTLS

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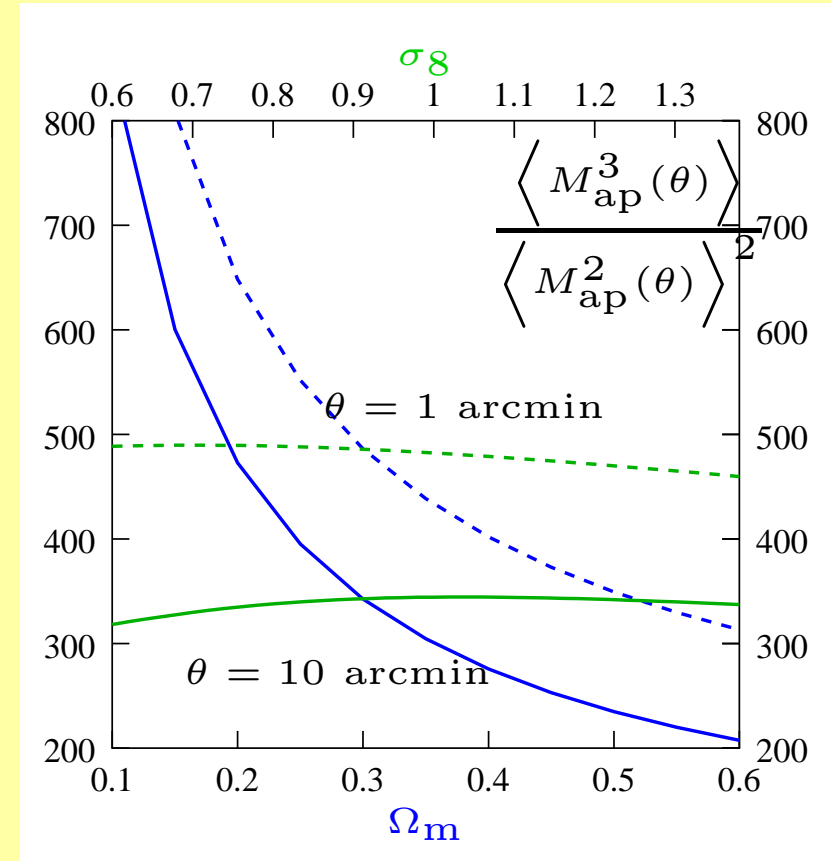
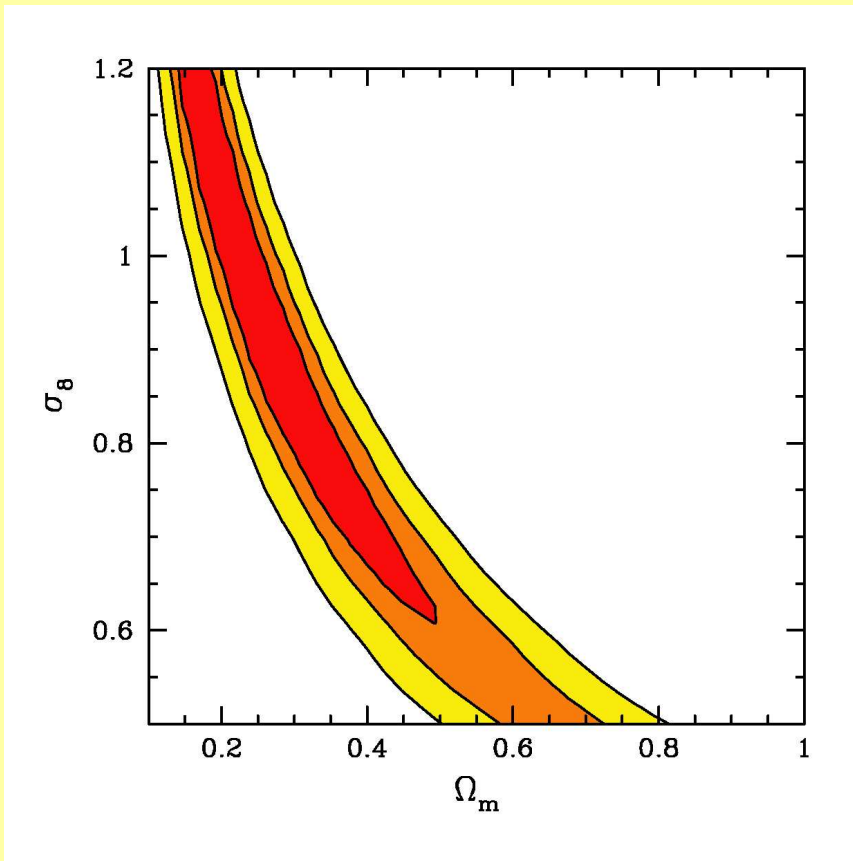
Recent results of (second-order) cosmic shear measurements

		σ_8 ($\Omega_m = 0.3$)
VIRMOS-DESCART	van Waerbeke et al. 2005	0.83 ± 0.07
CFHTLS wide	Hoekstra et al. 2006	0.85 ± 0.06
CFHTLS deep	Semboloni et al. 2006	0.9 ± 0.34
GEMS	Heymans et al. 2005	0.68 ± 0.13
GEMS/GOODS	Schrabback et al. 2006	0.52 ± 0.2
GaBoDS	Hetterscheidt et al. 2006	0.8 ± 0.1
CTIO	Jarvis et al. 2006	0.81 ± 0.15 (+ SN, CMB)

Motivation for 3rd-order shear statistics

- Together with 2nd order: lift parameter near-degeneracies, e.g.:

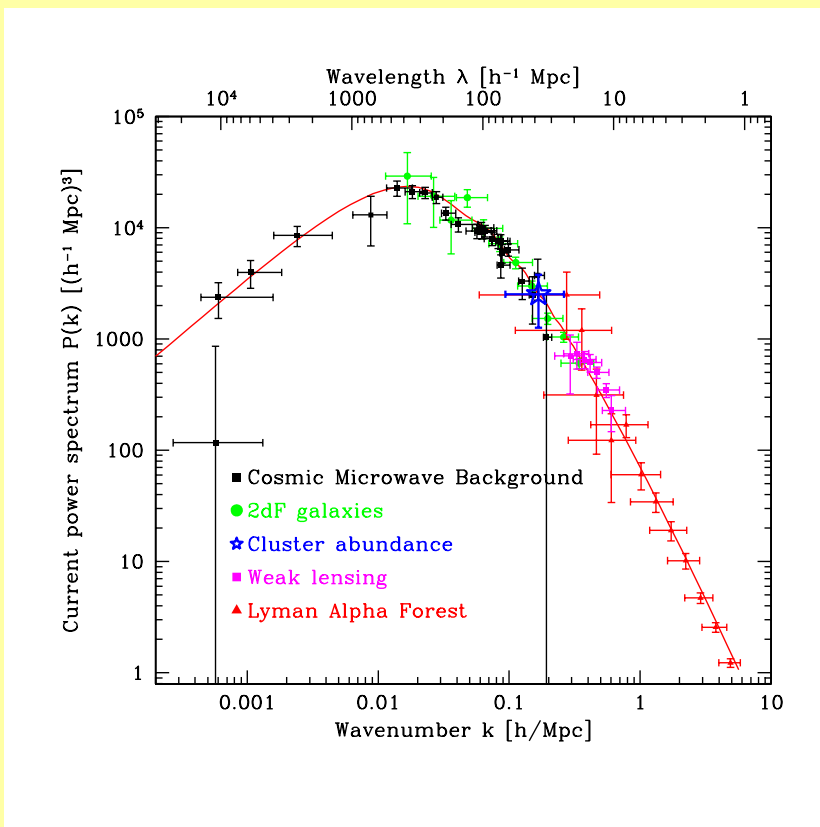
$$s_3(\theta) = \frac{\langle \kappa^3(\theta) \rangle}{\langle \kappa^2(\theta) \rangle^2} \text{ independent of } \sigma_8 \text{ [Bernardeau, van Waerbeke, Mellier 1997]}$$



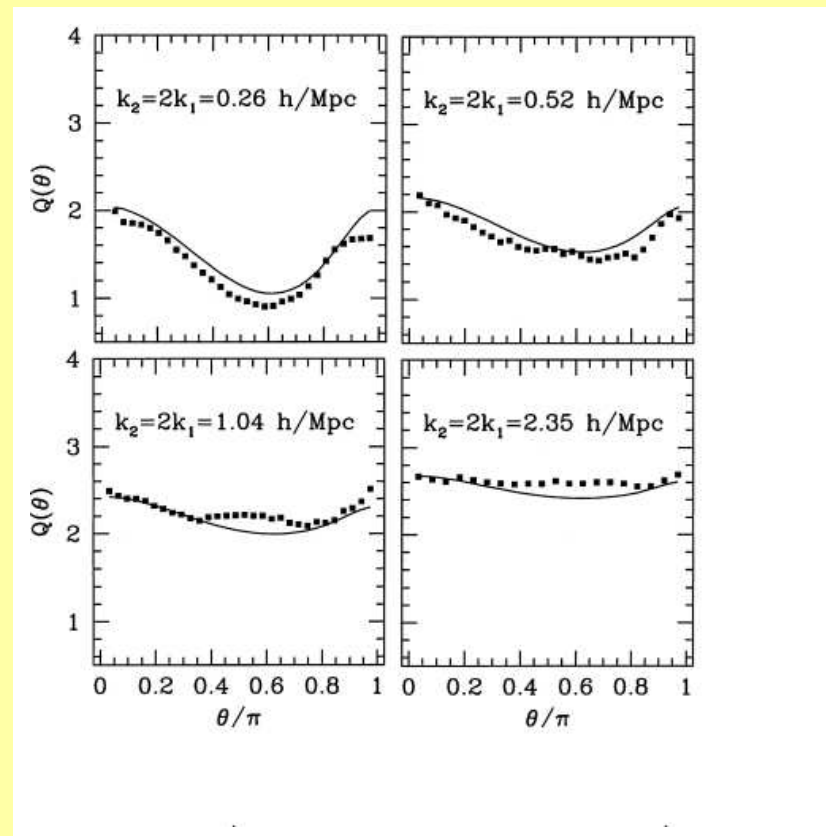
[CFHTLS wide]

- Dark matter power spectrum and bispectrum contain complementary information about cosmology.
- Probe Non-Gaussianity of the LSS on small scales, ($\lesssim 10'$), non-linear gravitational collapse, mode-coupling of the LSS, virialization of halos in hierarchical structure formation

power spectrum

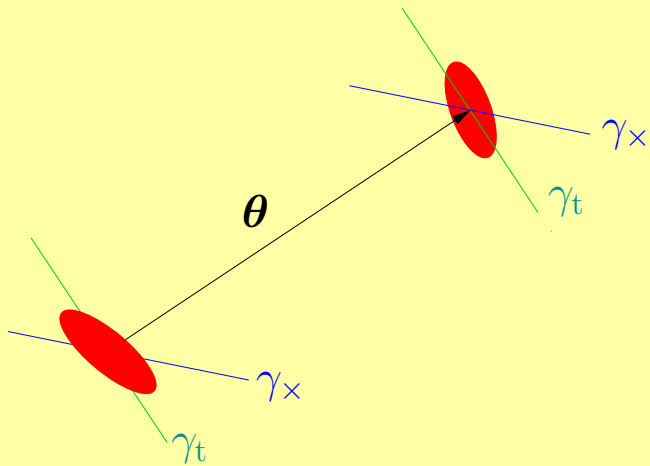


bispectrum



Weak lensing observables

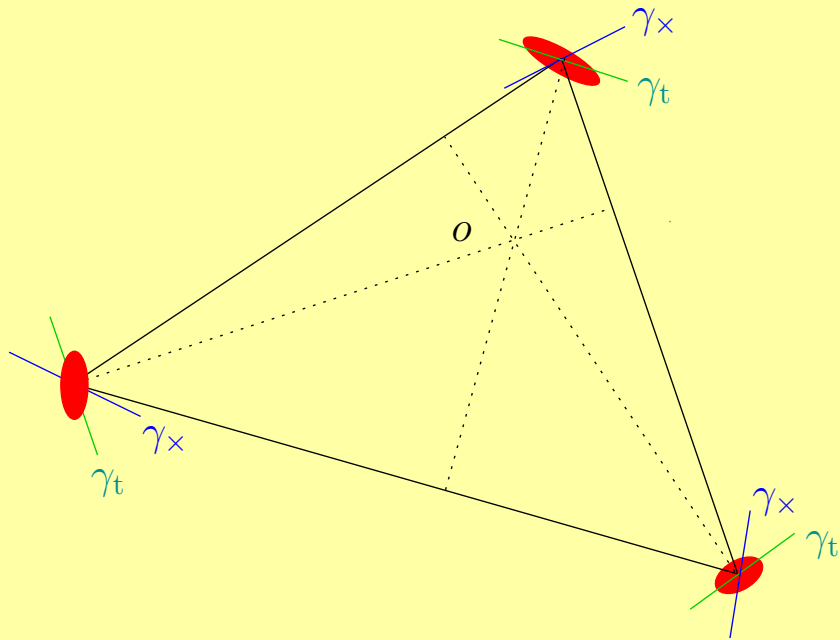
Two-point correlation function



$$\left. \begin{array}{l} \langle \gamma_t \gamma_t \rangle \\ \langle \gamma_x \gamma_x \rangle \\ \langle \gamma_t \gamma_x \rangle \\ \langle \gamma_x \gamma_t \rangle \end{array} \right\} = 0 \quad \text{because of parity}$$

2PCF $\xi_{\pm}(\theta) \equiv \langle \gamma_t \gamma_t \rangle \pm \langle \gamma_x \gamma_x \rangle$, two components

Three-point correlation function

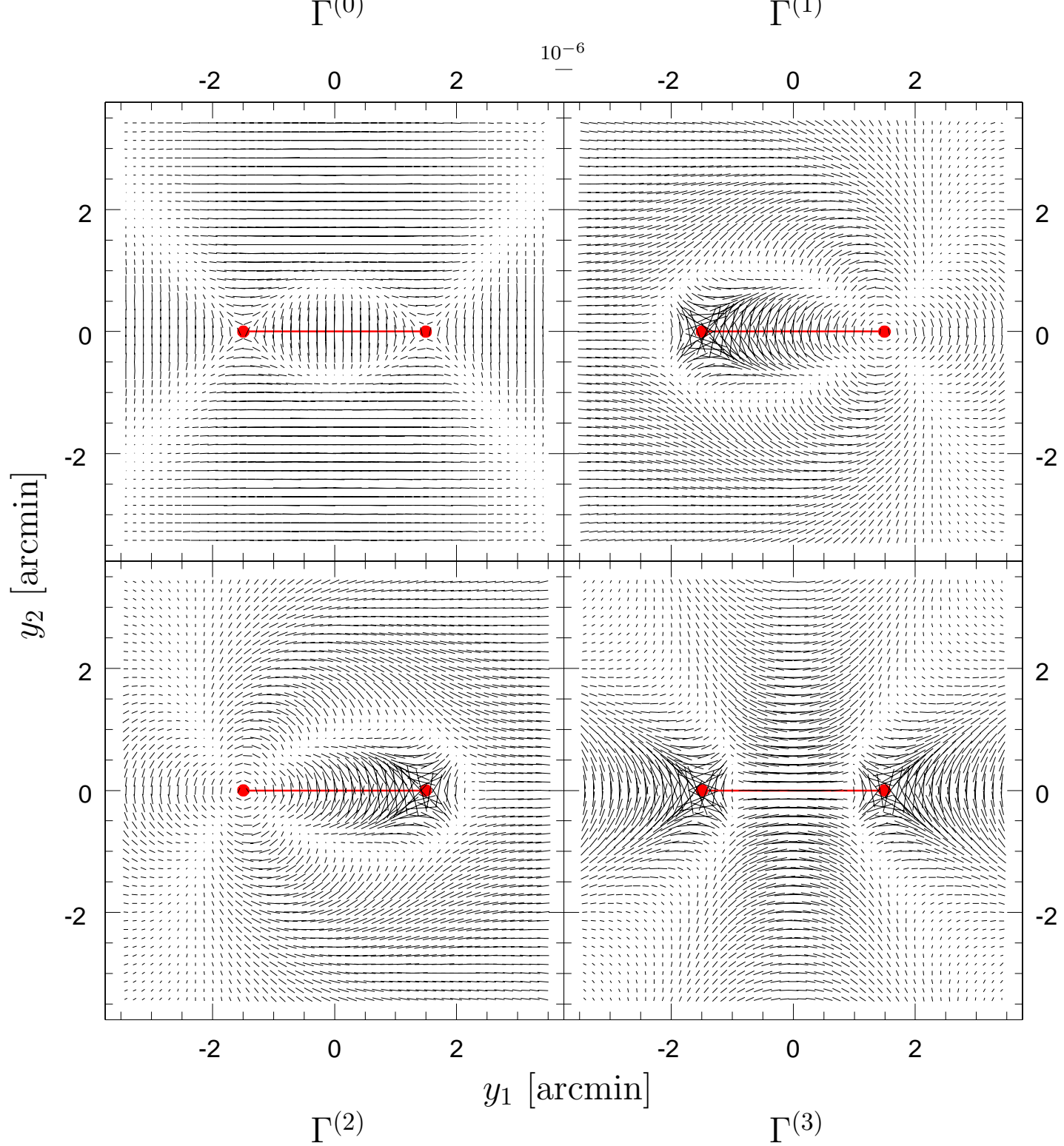


8 components:

$$\begin{array}{ll}
 \langle \gamma_t \gamma_t \gamma_t \rangle & \langle \gamma_t \gamma_t \gamma_x \rangle \\
 \langle \gamma_t \gamma_x \gamma_x \rangle & \langle \gamma_t \gamma_x \gamma_t \rangle \\
 \langle \gamma_x \gamma_t \gamma_x \rangle & \langle \gamma_x \gamma_t \gamma_t \rangle \\
 \langle \gamma_x \gamma_x \gamma_t \rangle & \langle \gamma_x \gamma_x \gamma_x \rangle
 \end{array}$$

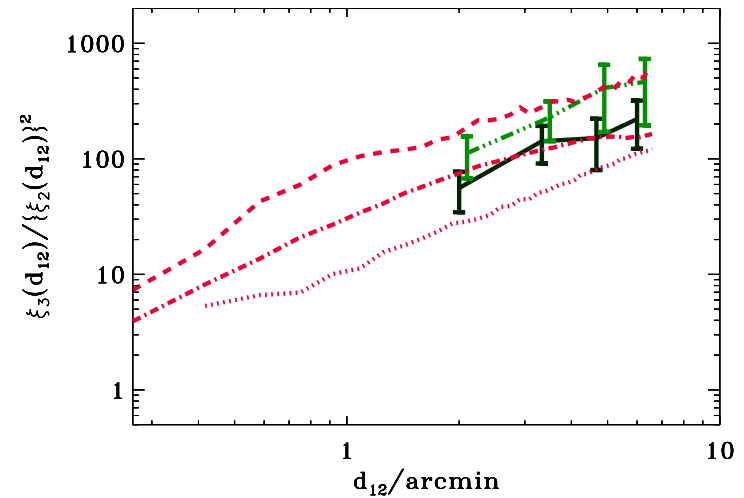
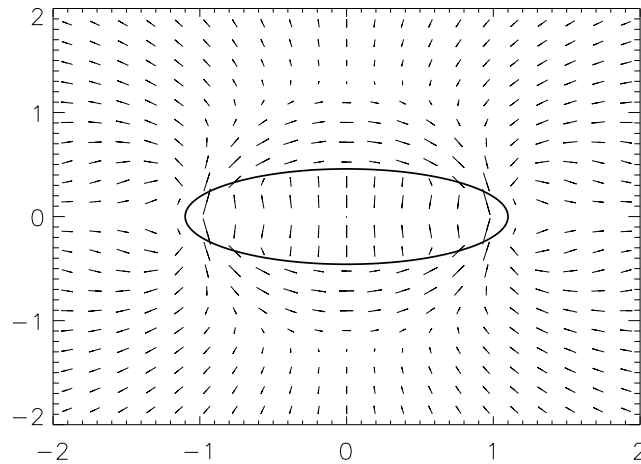
“Natural components” $\Gamma^{(0)}, \Gamma^{(1)}, \Gamma^{(2)}, \Gamma^{(3)} \in \mathbb{C} =$ linear combinations of the $\langle \gamma_\mu \gamma_\nu \gamma_\lambda \rangle$ [Schneider & Lombardi 2003]

3PCF has 8 (non-vanishing) components, depends on 3 quantities and is not a scalar 😞



Flavours of 3rd-order statistics

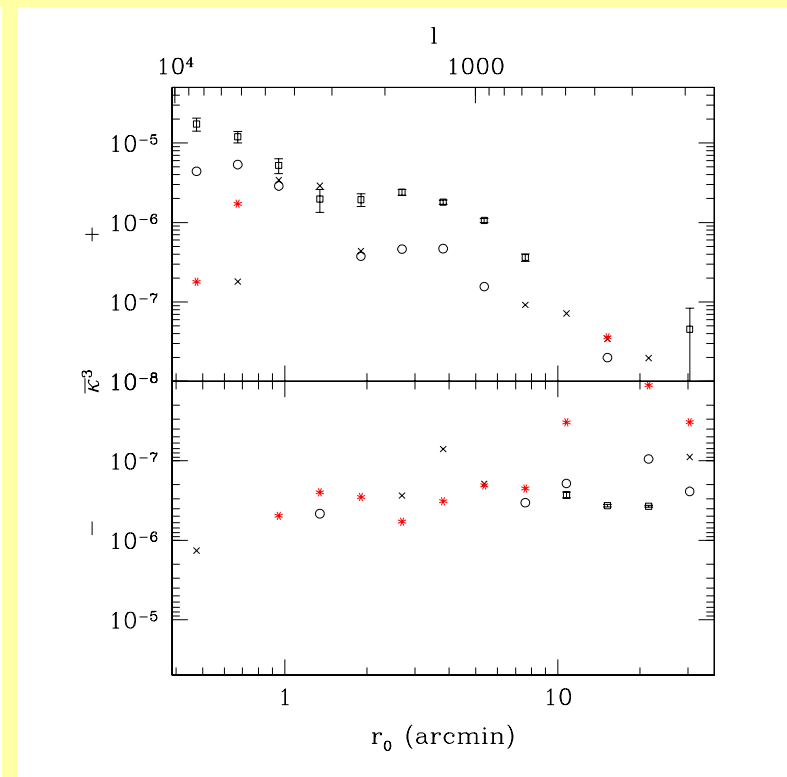
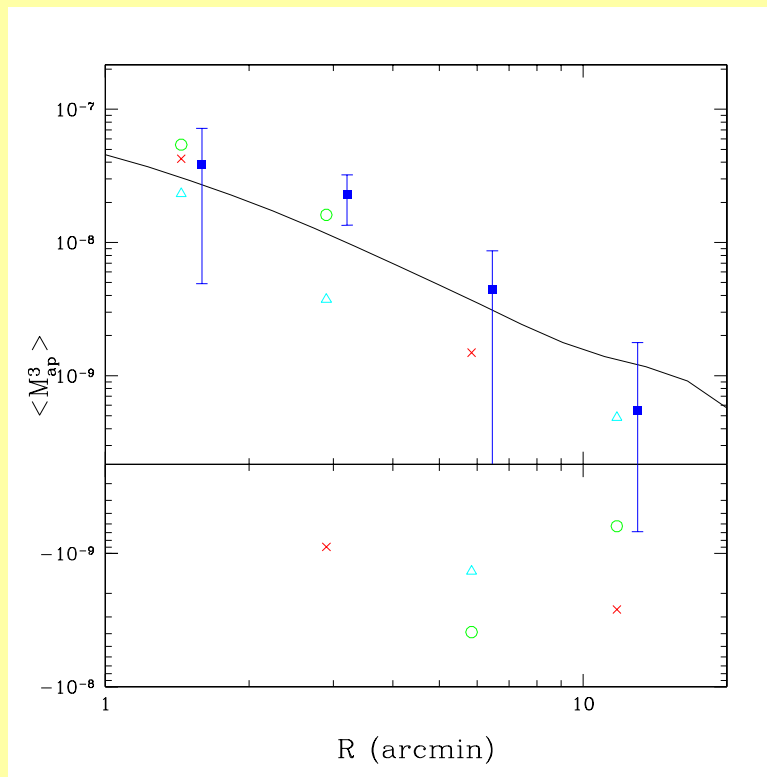
- projected 3PCF, integrated over elliptical region
[Bernardeau, van Waerbeke & Mellier 2002, 2003]



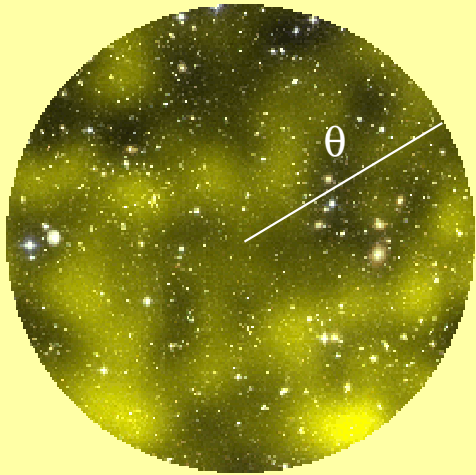
[VIRMOS-DESCART]

Flavours of 3rd-order statistics

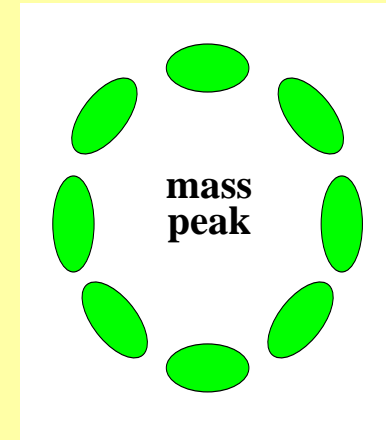
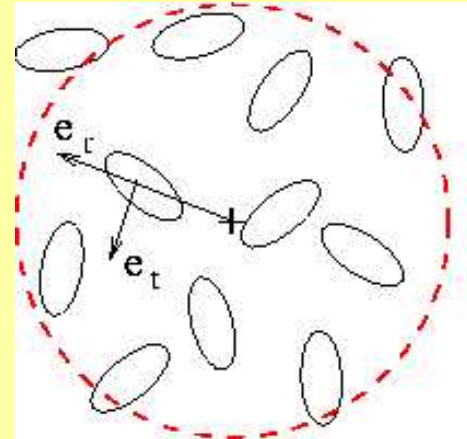
- projected 3PCF, integrated over elliptical region
[Bernardeau, van Waerbeke & Mellier 2002, 2003]
- Aperture-mass $\langle M_{\text{ap}}^3 \rangle$: CTIO [Jarvis et al. 2004] and
VIRMOS-DESCART [Pen et al. 2003]



Aperture-Mass Statistics

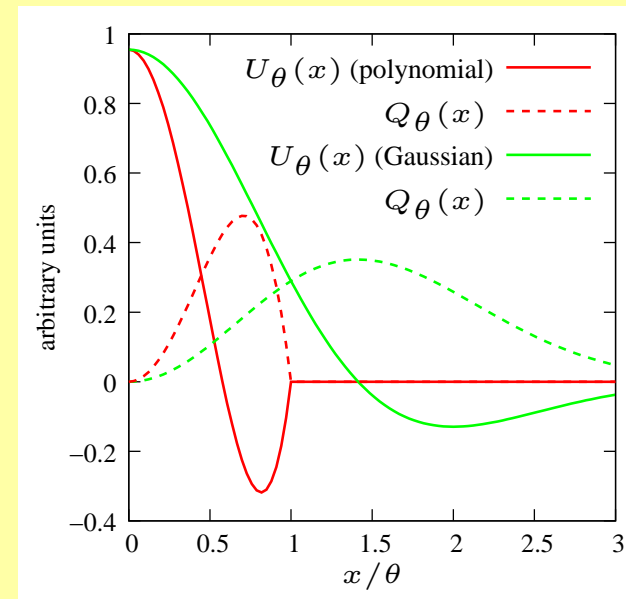


Mass overdensity κ



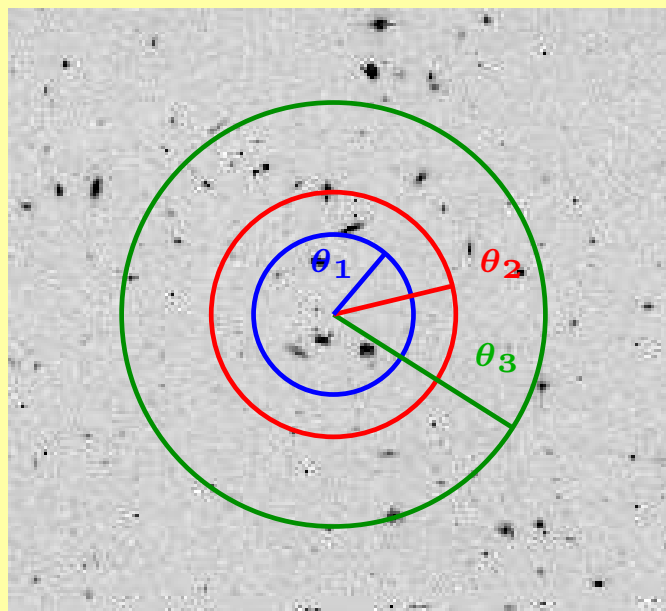
Tangential shear in aperture γ_t

$$\begin{aligned}
 M_{\text{ap}}(\theta) &= \int d^2x U_{\theta}(x) \kappa(\vec{x}) \\
 &= \int d^2x Q_{\theta}(x) \gamma_t(\vec{x})
 \end{aligned}$$



Second- and third-order statistics

- variance $\langle M_{\text{ap}}^2(\theta) \rangle$ probes power spectrum $P_{\kappa}(\ell)$ at a scale $\ell \propto 1/\theta$
- skewness $\langle M_{\text{ap}}(\theta_1) M_{\text{ap}}(\theta_2) M_{\text{ap}}(\theta_3) \rangle$ probes bispectrum $B_{\kappa}(\ell_1 \propto 1/\theta_1, \ell_2 \propto 1/\theta_2, \ell_3 \propto 1/\theta_3)$, cross-correlation or mode coupling of the large-scale structure on different scales [Schneider, MK & Lombardi 2005]



convergence κ

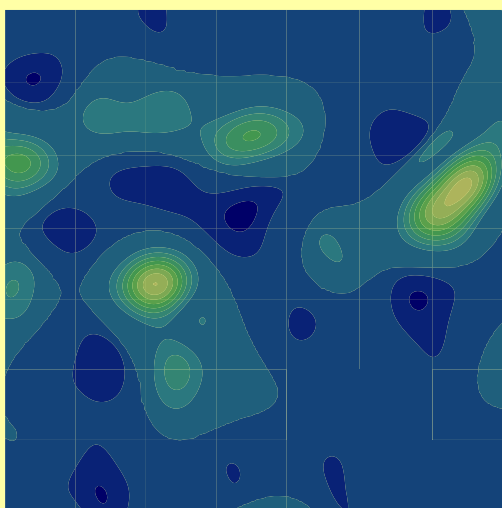
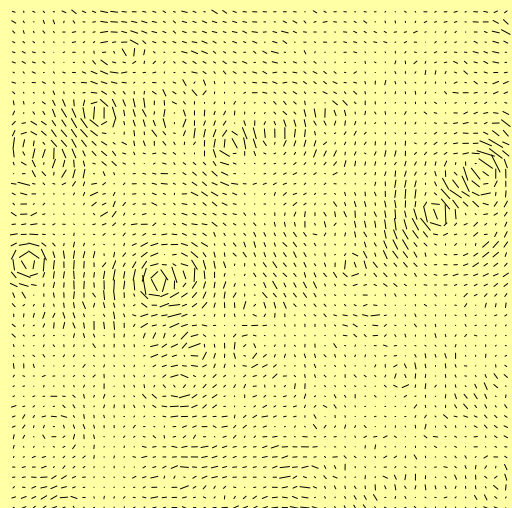
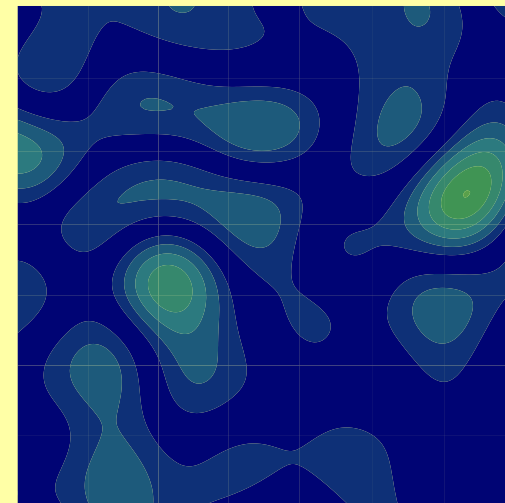
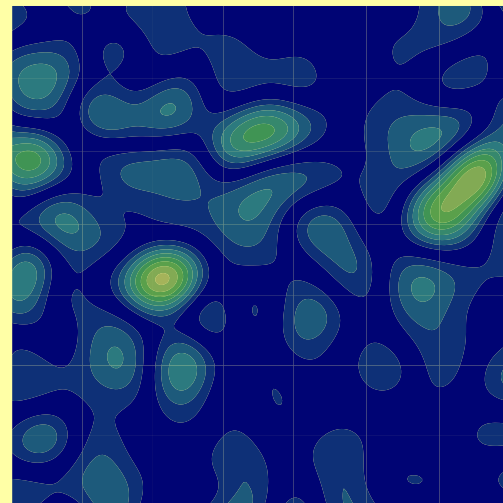
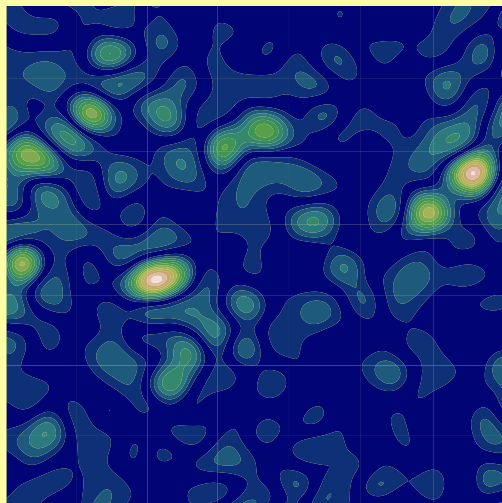
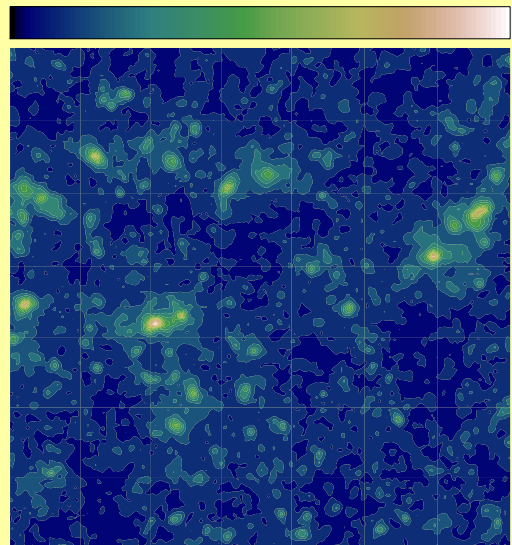
1.4'

$M_{\text{ap}}^2(\theta)$

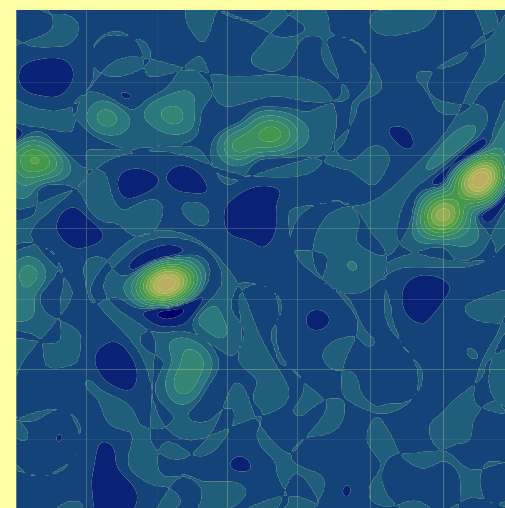
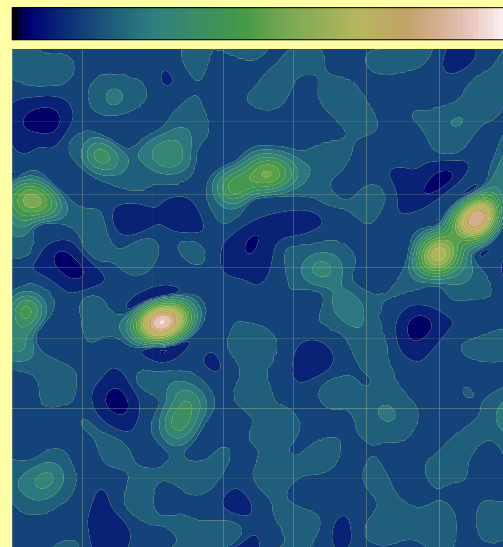
3.9'

-0.041 0.095 0.23

0 1.6e-04 6.2e-04



-3.2e-07 1.0e-06 6.9e-06



shear γ

(2.3', 2.3', 3.9')

(1.4', 1.4', 2.3')

(1.4', 2.3', 3.9')

$M_{\text{ap}}^3(\theta_1, \theta_2, \theta_3)$

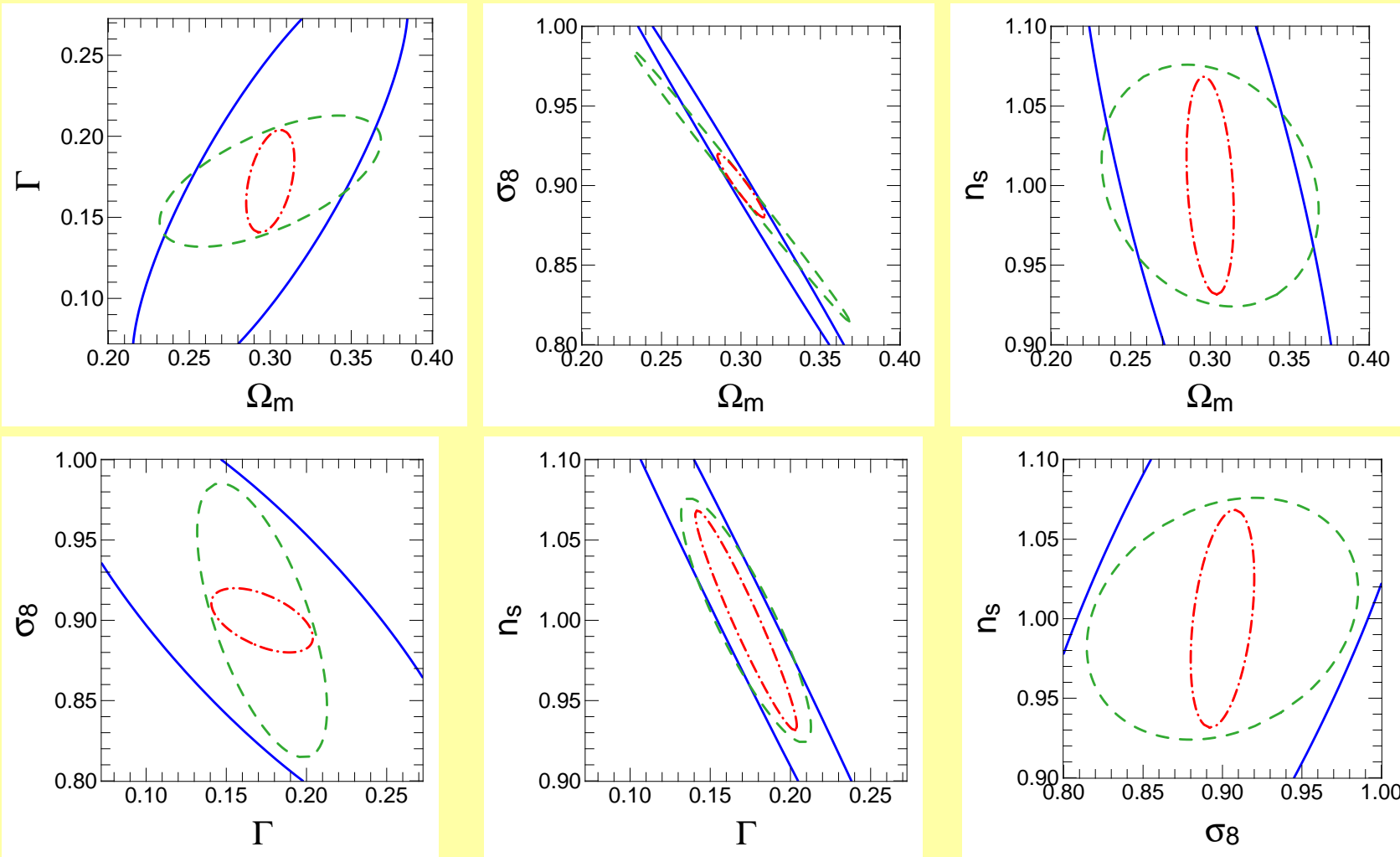
Why M_{ap} ?

- $\langle M_{\text{ap}}^3 \rangle$ is a scalar
- separates E- & B-modes
- one can obtain $\langle M_{\text{ap}}^2 \rangle$, $\langle M_{\text{ap}}^3 \rangle$ from 2pcf, 3pcf *or* directly from galaxy field [work in progress]
- localized filter \rightarrow simple relation to power spectrum/bispectrum
- no information loss if $\langle M_{\text{ap}}^3 \rangle$ is used instead of 3pcf [work in progress]

Predictions for CFHTLS wide

- Ray-tracing simulations [T. Hamana]
- All galaxies at redshift $z_0 = 0.977$, $\bar{n} = 25/\text{arcmin}^2$
- Shear signal for aperture radii θ between 1 and 15 arc minutes
- Error budget: shape noise, (non-Gaussian) cosmic variance, cross-correlation between $\langle M_{\text{ap}}^2 \rangle$ and $\langle M_{\text{ap}}^3 \rangle$ [MK & Schneider 2005]
- Fisher matrix for 1σ errors

Parameter constraints

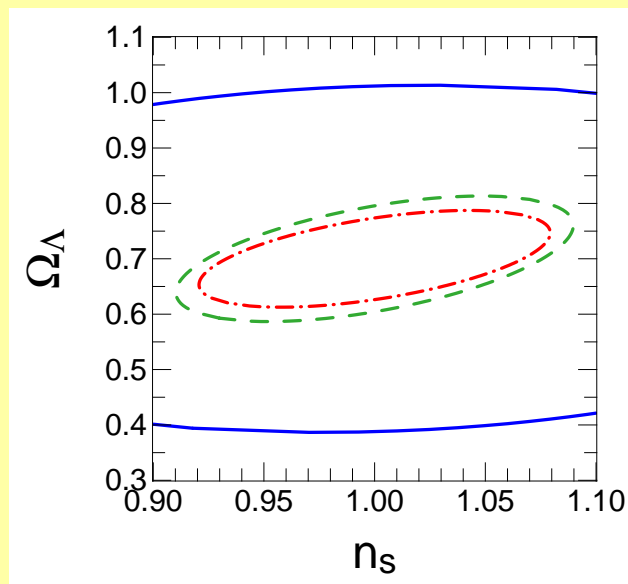
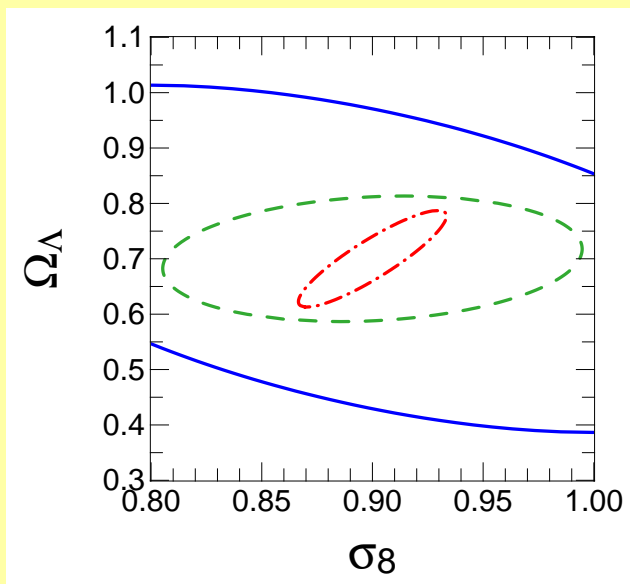
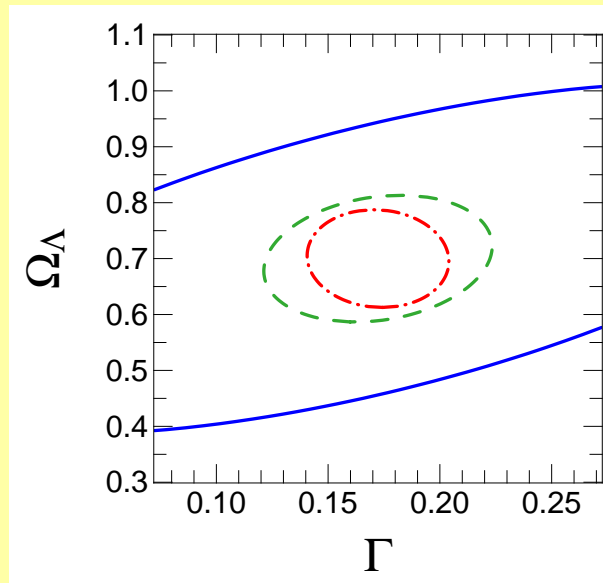
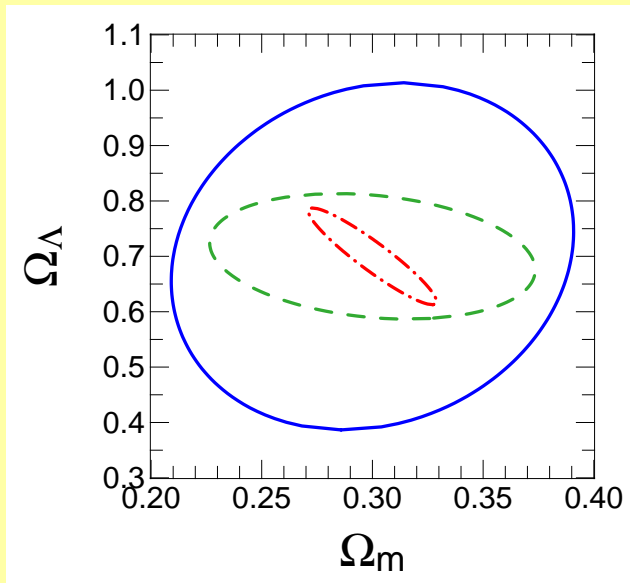


1σ -ellipses

$\langle M_{\text{ap}}^2 \rangle$ ————
 $\langle M_{\text{ap}}^3 \rangle$ - - - -
 combin. - · - · -

Fisher matrix with 5 parameters assuming flat Λ CDM cosmology
 $\sigma(z_0) = 0.01$

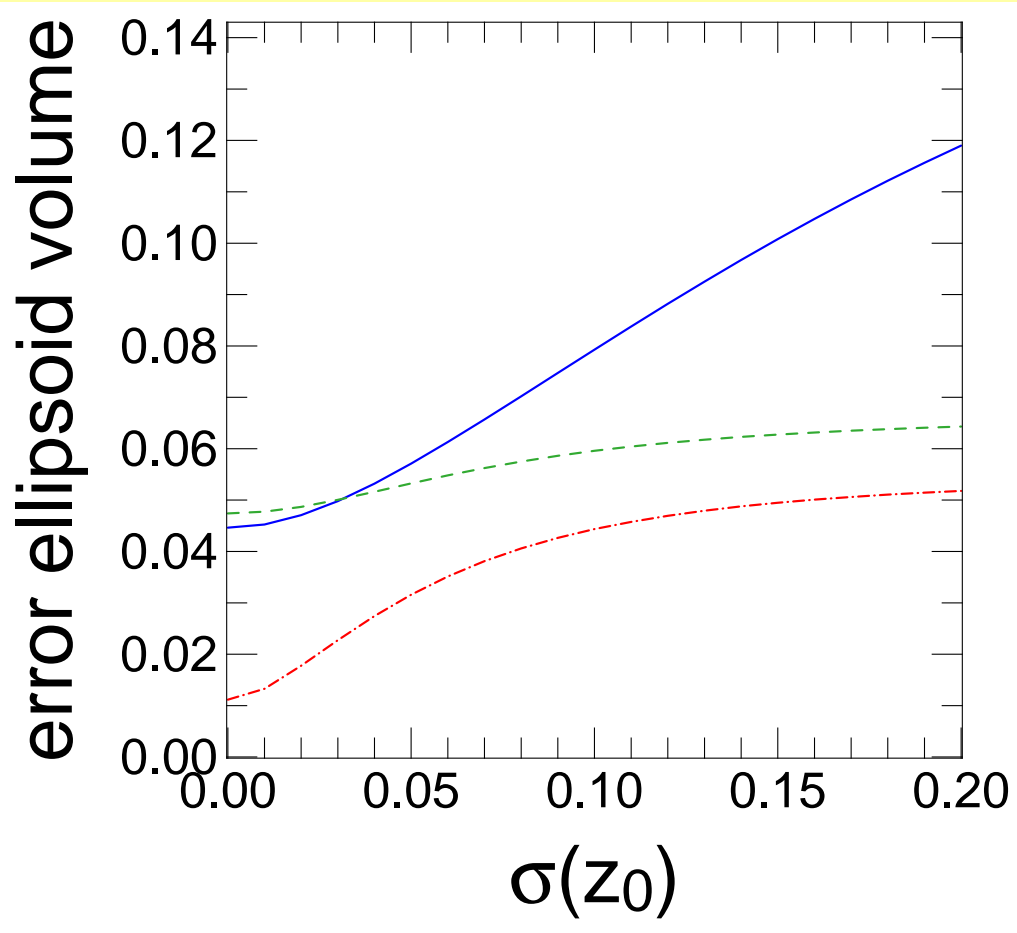
Dark energy constraints



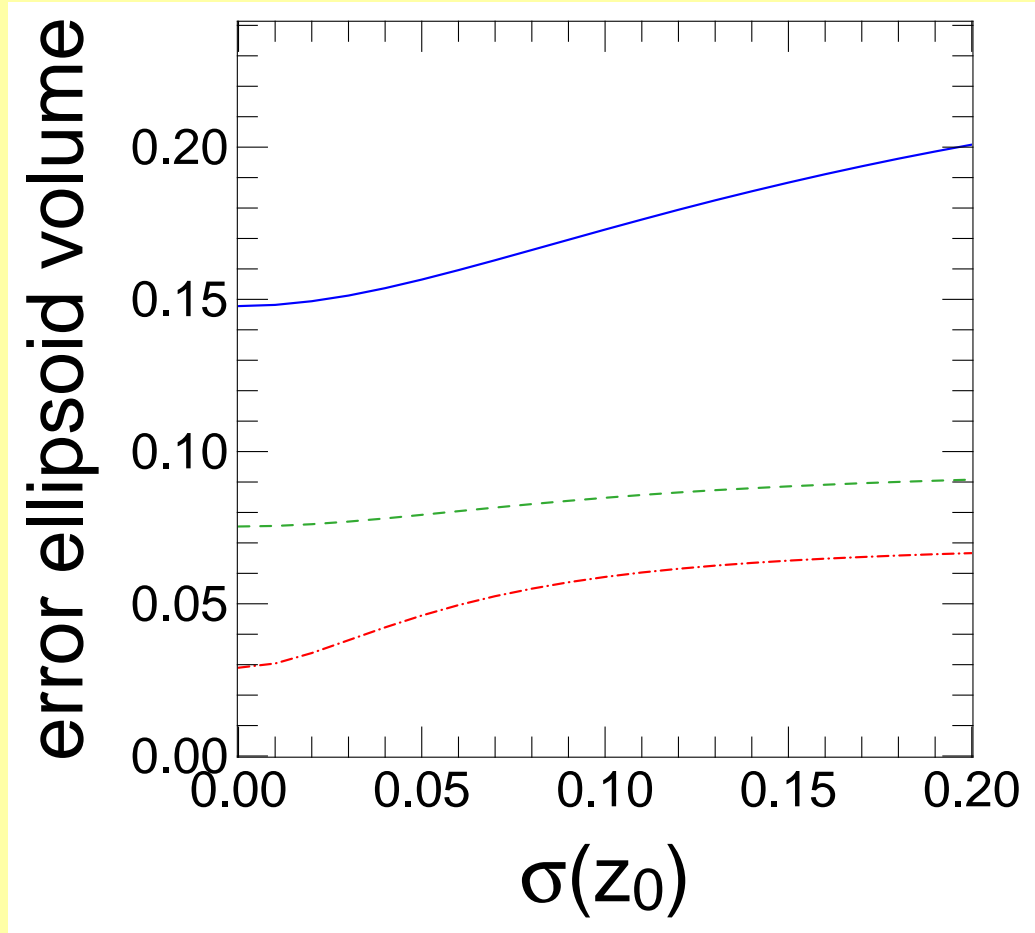
1σ -ellipses

$\langle M_{\text{ap}}^2 \rangle$ ————
 $\langle M_{\text{ap}}^3 \rangle$ - - - -
combin. - · - · -

Redshift uncertainty



$(\Omega_m, \Gamma, \sigma_8)$, flat



$(\Omega_m, \Omega_\Lambda, \Gamma, \sigma_8)$

Parameter constraints summary

Flat Λ CDM

	Ω_m	Γ	σ_8	n_s
$\langle M_{\text{ap}}^2 \rangle$	0.085	0.13	0.141	0.27
$\langle M_{\text{ap}}^3 \rangle$	0.068	0.04	0.085	0.076
combination	0.015	0.032	0.02	0.068

General Λ CDM

	Ω_m	Ω_Λ	Γ	σ_8	n_s
$\langle M_{\text{ap}}^2 \rangle$	0.091	0.31	0.21	0.2	0.27
$\langle M_{\text{ap}}^3 \rangle$	0.074	0.11	0.051	0.095	0.09
combination	0.029	0.087	0.032	0.033	0.079

Figurs are marginalized 1σ -errors, $\sigma(z_0) = 0.01$

Summary and Outlook

- Third-order cosmic shear statistic has been detected and is expected to be measured with high significance in CFHTLS
- Joint $\langle M_{\text{ap}}^2 \rangle$ and $\langle M_{\text{ap}}^3 \rangle$ observations can improve parameter constraints by a factor of 3 – 6 (with flat prior)
- For tight constraints on dark energy tomography needed
[\[Takada & Jain 2004\]](#)
- Source redshifts have to be known accurately
- Need better models of non-linear structure formation!

References

- Bernardeau, F., van Waerbeke, L. & Mellier, Y. 1997, A&A, 322, 1
- —. 2002, A&A, 389, L28
- —. 2003, A&A, 397, 405
- Hamana, T. & Mellier, Y. 2001, MNRAS, 327, 169
- Jarvis, M., Bernstein, G., & Jain, B. 2004, MNRAS, 352, 338
- Kilbinger, M. & Schneider, P. 2005, A&A, 442, 69
- Pen, U.-L., Zhang, T., van Waerbeke, L. et al. 2003, ApJ, 592, 664
- Schneider, P., Kilbinger, M. & Lombardi, M. 2005, A&A, 431, 9
- Schneider, P. & Lombardi, M. 2003, A&A, 307, 809
- Scoccimarro, R. & Couchman, H, 2001, MNRAS, 325, 1312
- Takada, M. & Jain, B. 2004, MNRAS, 348, 897